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needed. In both of them, shape changes of atomic nuclei can be described by a specific Hamil-tonian. From algebraic point of view, existence of the shape phase transitions is proposed by the З study of dynamical symmetries of some algebraic structures that different shapes coincide with з particular dynamical symmetries [11]. From geometrical aspect, critical point symmetries can be studied by solving the collective Hamiltonian with an appropriate potential [12-14]. Inves-tigation of eigenfunctions and eigenvalues of the Bohr Hamiltonian yields predicting the phase transitions in the shapes of the deformed nuclei [15]. Bohr Hamiltonian has been investigated in various situations by different physicist such as Fortunato [15], Bonatsos et al. [16,17], Capak et al. [18], Chabab et al. [19] and Hassanabadi et al. [20–24]. Davydov–Chaban is an approach to study collective motion of atomic nuclei. This is obtained from Bohr Hamiltonian by considering rigidity for γ in Bohr Hamiltonian. This approach also has been an important issue in nuclear physics since it covers some interesting results such as what is mentioned in articles published by Buganu and Budaca [25] or Bonatsos et al. [26]. Also, we published an article in which we discussed in detail that using Davydov-Chaban Hamiltonian how we can describe instability of unstable nuclei [27].

In this article, we are interested in analytical investigation of Bohr Hamiltonian for two well-known separable cases, the approximate and exact one, in which there is a Killingbeck interaction. So, in Sec 2, Bohr Hamiltonian for two cases and the potential are introduced. In Sec. 3, experimental data related to energy and B(E2) transition rates of some isotopes are re-produced.

2. Bohr Hamiltonian and Killingbeck interaction

The original Bohr Hamiltonian [28,40] is

$$H = -\frac{\hbar^2}{2B} \left[\frac{1}{\beta^4} \frac{\partial}{\partial \beta} \beta^4 \frac{\partial}{\partial \beta} + \frac{1}{\beta^2 \sin(3\gamma)} \frac{\partial}{\partial \gamma} \sin(3\gamma) \frac{\partial}{\partial \gamma} - \frac{1}{4\beta^2} \sum_{k=1}^3 \frac{Q_k^2}{\sin^2(\gamma - \frac{2}{3}\pi k)} \right] + V(\beta, \gamma)$$
(2.1)

in which, intrinsic deformation coordinates are denoted by β (deformation coordinate measuring departure from the spherical shape), γ (the angle measuring departure from axial symmetry) and Q_k (k = 1, 2, 3) the operators of the total angular momentum projections in the intrinsic reference, mass parameter and Euler angles are shown by B and θ respectively. Considering the wave function as $\Psi^{as(es)}(\beta, \gamma, \theta_i) = \xi_{L,\alpha}^{as(es)}(\beta) \Gamma^{as(es)}(\gamma) D_{M,\alpha}^L(\theta_i)$ where $D_{M,\alpha}^L(\theta_i)$ represents Wigner functions, M and α are the eigenvalues of the projections of angular momentum on the laboratory fixed \hat{z} axis and the body-fixed \hat{x}' axis respectively, es stands for exact separation and as means approximate separation case, we want to study the Bohr Hamiltonian for two cases.

2.1. Triaxial case with
$$\gamma \approx \frac{\pi}{6}$$

In this case in which the potential has a minimum around $\gamma \approx \frac{\pi}{6}$ one can write the last term of Eq. (2.1) in the form [40]

$$\sum_{k=1}^{45} \frac{Q_k^2}{\sin^2\left(\gamma - \frac{2\pi}{3}k\right)} \approx Q_1^2 + 4\left(Q_2^2 + Q_3^2\right) = 4\left(Q_1^2 + Q_2^2 + Q_3^2\right) - 3Q_1^2.$$
(2.2)

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