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2 *H. Sobhani et al. / Nuclear Physics A* ••• *(*••••*)* •••*–*•••

¹ needed. In both of them, shape changes of atomic nuclei can be described by a specific Hamil-² tonian. From algebraic point of view, existence of the shape phase transitions is proposed by the ² ³ study of dynamical symmetries of some algebraic structures that different shapes coincide with ³ ⁴ particular dynamical symmetries [\[11\]](#page--1-0). From geometrical aspect, critical point symmetries can ⁴ ⁵ be studied by solving the collective Hamiltonian with an appropriate potential [\[12–14\]](#page--1-0). Inves-⁶ tigation of eigenfunctions and eigenvalues of the Bohr Hamiltonian yields predicting the phase ⁶ ⁷ transitions in the shapes of the deformed nuclei [\[15\]](#page--1-0). Bohr Hamiltonian has been investigated in ⁷ ⁸ various situations by different physicist such as Fortunato [\[15\]](#page--1-0), Bonatsos et al. [\[16,17\]](#page--1-0), Capak⁸ ⁹ et al. [\[18\]](#page--1-0), Chabab et al. [\[19\]](#page--1-0) and Hassanabadi et al. [\[20–24\]](#page--1-0). Davydov–Chaban is an approach to 9 ¹⁰ study collective motion of atomic nuclei. This is obtained from Bohr Hamiltonian by considering ¹⁰ ¹¹ rigidity for γ in Bohr Hamiltonian. This approach also has been an important issue in nuclear ¹¹ ¹² physics since it covers some interesting results such as what is mentioned in articles published¹² ¹³ by Buganu and Budaca $[25]$ or Bonatsos et al. $[26]$. Also, we published an article in which we ¹³ ¹⁴ discussed in detail that using Davydov–Chaban Hamiltonian how we can describe instability of ¹⁴ ¹⁵ unstable nuclei [\[27\]](#page--1-0).

¹⁶ In this article, we are interested in analytical investigation of Bohr Hamiltonian for two ¹⁶ ¹⁷ well-known separable cases, the approximate and exact one, in which there is a Killingbeck¹⁷ ¹⁸ interaction. So, in Sec 2, Bohr Hamiltonian for two cases and the potential are introduced. In ¹⁸ 19 19 Sec. [3,](#page--1-0) experimental data related to energy and B*(E*2*)* transition rates of some isotopes are re-20 produced 20 and 20 produced.

21 21

23 23

22 22 **2. Bohr Hamiltonian and Killingbeck interaction**

 $\frac{24}{24}$ The original Bohr Hamiltonian [\[28,40\]](#page--1-0) is $\frac{24}{24}$ 25 25

$$
H = -\frac{\hbar^2}{2B} \left[\frac{1}{\beta^4} \frac{\partial}{\partial \beta} \beta^4 \frac{\partial}{\partial \beta} + \frac{1}{\beta^2 \sin(3\gamma)} \frac{\partial}{\partial \gamma} \sin(3\gamma) \frac{\partial}{\partial \gamma} - \frac{1}{4\beta^2} \sum_{k=1}^3 \frac{Q_k^2}{\sin^2(\gamma - \frac{2}{3}\pi k)} \right]_{27}^{28} + V(\beta, \gamma) \tag{2.1}
$$

 31 in which, intrinsic deformation coordinates are denoted by *β* (deformation coordinate measuring 32 departure from the spherical shape), *γ* (the angle measuring departure from axial symmetry) 33 and Q_k ($k = 1, 2, 3$) the operators of the total angular momentum projections in the intrinsic 33 34 reference, mass parameter and Euler angles are shown by *B* and *θ* respectively. Considering the 35 wave function as $\Psi^{as(es)}(\beta, \gamma, \theta_i) = \xi_{L,\alpha}^{as(es)}(\beta) \Gamma^{as(es)}(\gamma) D_{M,\alpha}^L(\theta_i)$ where $D_{M,\alpha}^L(\theta_i)$ represents 35 ³⁶ Wigner functions, *M* and *α* are the eigenvalues of the projections of angular momentum on the ³⁶ 37 laboratory fixed *z*ˆ axis and the body-fixed *x*ˆ axis respectively, *es* stands for exact separation and 38 *as* means approximate separation case, we want to study the Bohr Hamiltonian for two cases.

$$
40 \quad 2.1. \text{ Triaxial case with } \gamma \approx \frac{\pi}{6}
$$

⁴² In this case in which the potential has a minimum around $\gamma \approx \frac{\pi}{6}$ one can write the last term ⁴² ⁴³ of Eq. (2.1) in the form [\[40\]](#page--1-0) 44 44

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$$
\sum_{k=1}^{45} \frac{Q_k^2}{\sin^2(\gamma - \frac{2\pi}{3}k)} \approx Q_1^2 + 4\left(Q_2^2 + Q_3^2\right) = 4\left(Q_1^2 + Q_2^2 + Q_3^2\right) - 3Q_1^2. \tag{2.2}
$$

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