

S. Goudarzi et al. / Nuclear Physics A ••• (••••) •••-•••

¹ pure neutron matter (PNM), has not been established well despite of its crucial role in both ² nuclear physics and astrophysics. Motivated by this fact, studying the nuclear symmetry energy ³ E_{sym} which mostly governs the PNM EOS is currently an active field of research in nuclear ⁴ physics. ⁴

Since the symmetry energy can not be directly measured, it is of fundamental importance to identify observables which strongly correlate with E_{sym} and its density slope L to impose constraints on these quantities. Additional information on the value and density dependence of symmetry energy can be extracted from the astrophysical observations of compact objects, namely neutron stars. Among the terrestrial laboratory observables, the neutron skin thickness is widely studied [1-6] during last decades. Nowadays, there are many other observables such as isospin fractionation, neutron-proton differential flow, nuclear mass systematics, low-lying E1 mode, etc., which are believed to be sensitive to symmetry energy and particularly, its density slope. We refer to references [7-11] and references therein quoted for more details. Density de-pendence of this quantity also relates the heavy-ion reactions [12-18], stability of superheavy nuclei [19], cooling [20–24] and the mass-radius relations of neutron stars [25,26], and proper-ties of nuclei involved in r-process nucleosynthesis [27]. On the other hand, the symmetry energy or symmetry free energy of hot neutron-rich matter is important for understanding some physi-cal and astrophysical phenomena such as the liquid-gas phase transition of asymmetric nuclear matter, the dynamical evolution of massive stars and the supernova explosion mechanisms [28].

In order to theoretically predict the symmetry energy, asymmetric nuclear matter is stud-ied in the framework of various microscopic and phenomenological approaches by using a variety of microscopic and phenomenological two-nucleon forces and phenomenological three-body forces. Such many-body techniques include the Brueckner–Hartree–Fock (BHF) [7,29] and its relativistic counterpart Dirac-Brueckner-Hartree-Fock (DBHF) [30] frameworks, the self-consistent Green's function method [31], variational approaches [32,33], mean-field model [34], the Quantum Monte Carlo (QMC) technique [35], extended relativistic mean field (ERMF) model [36], and Skyrme-Hartree-Fock (SHF) method [37]. There are also some recent calcu-lations for the symmetry energy of finite nuclei (see e.g. [38] and references therein). Recent progress and challenges in both theoretically and experimentally measuring the symmetry en-ergy and its density dependence, specially at supra nuclear densities, are also reviewed in Refs. [39-43]. On the other hand, less theoretical and experimental attention is so far devoted to studying the temperature and density dependence of the symmetry energy or symmetry free energy of nuclear matter as well as finite nuclei [44–50].

Despite of numerous attempts that were made to determine the symmetry energy properties, especially the value of $E_{sym}(\rho_0)$ and its density slope L at saturation density ρ_0 , these quan-tities are still uncertain. Results are strongly dependent on different experimental techniques and many-body approaches as well as nuclear interaction models. Especially, the density de-pendence of symmetry energy is poorly known. The high density prediction of many-body models for the symmetry energy are extremely diverse and sometimes contradictory. On the other hand, although the semi-empirical value of the symmetry energy is known to be in the range of $\sim 31.6 \pm 2.66$ MeV [41], explorations on L show wide variations. The obtained theoretical lower limit is ~ 20 MeV [51] and the upper one is even higher than 170 MeV [52]. However, recent surveys of analyses of terrestrial nuclear laboratory experiments and astrophysical observations found that the mean semi-empirical value of the slope parameter at saturation density lies in the range of $58.9 \pm 16.0 \text{ MeV}$ [41,11].

In the present work, within the self-consistent lowest order constrained variational (LOCV)
approach using the Argonne V18 two-body potential [53] supplemented by an Urbana type three-

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