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## Structure of states in <sup>12</sup>Be via the <sup>11</sup>Be(d, p) reaction

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#### ABSTRACT

The s-wave neutron fraction of the 0<sup>+</sup> levels in <sup>12</sup>Be has been investigated for the first time through the <sup>11</sup>Be(d, p) transfer reaction using a 5 A MeV <sup>11</sup>Be beam at TRIUMF, Canada. The reaction populated all the known bound states of <sup>12</sup>Be. The ground state *s*-wave spectroscopic factor was determined to be  $0.28^{+0.03}_{-0.07}$  while that for the long-lived 0<sup>+</sup><sub>2</sub> excited state was  $0.73^{+0.27}_{-0.40}$ . This observation, together with the smaller effective separation energy indicates enhanced probability for an extended density tail beyond the <sup>10</sup>Be core for the 0<sup>+</sup><sub>2</sub> excited state compared to the ground state.

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Understanding the structure evolution from stable to halo nuclei requires systematic studies of changes in orbital occupancies. So far only Borromean nuclei, with unbound core + n clusters, have

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shown a two-neutron halo structure, such as in <sup>11</sup>Li [1]. The existence of such a structure in a non-Borromean nucleus is not yet established and investigating this in <sup>12</sup>Be is of particular interest. The disappearance of the conventional N = 8 magic number in <sup>12</sup>Be is suggested from the lowering of the  $2_1^+$  state [2]. This is associated with the  $2s_{1/2}$  and  $1d_{5/2}$  orbitals intruding into the *p*-shell. <sup>11</sup>Be is known to be a one-neutron halo with the neutron dominantly occupying the  $2s_{1/2}$  orbital with spectroscopic factor  $S \sim 0.74$  in agreement with the one-neutron removal cross-section [3]. This is also in agreement with the observations from transfer reactions [4–6].

The one-neutron separation energy  $(S_n)$  of <sup>12</sup>Be is 3169(16) keV, while that of <sup>11</sup>Be is 504(6) keV [7], and <sup>11</sup>Li has a two-neutron separation energy of 369.15(65) keV [8]. An important question is how the increase in neutron separation energy due to pairing influences the neutron distribution in the  $2s_{1/2}$  orbital.

Here we report the first measurement that determines the *s*-wave neutron occupancy of the 0<sup>+</sup> bound states in <sup>12</sup>Be through the <sup>11</sup>Be(*d*, *p*)<sup>12</sup>Be neutron-transfer reaction. Since the ground state of <sup>11</sup>Be is  $1/2^+$ , the selectivity of this reaction offers a clean way to disentangle the *s*-wave occupancies for the 0<sup>+</sup> levels. The results are therefore the first clean determination of the <sup>12</sup>Be<sub>gs</sub> spectroscopic factor without any influence of the isomeric 0<sup>+</sup> excited state. The observations also show the *s*-wave composition of the 0<sup>+</sup> excited state.

The first excited state  $(2^+)$  in <sup>12</sup>Be was observed at 2.11(2) MeV through inelastic scattering [2,9] and heavy ion transfer [10]. A state at 2.68(3) MeV, was interpreted as an l = 1 excitation [9, 12] which is considered to be the same state as observed in the <sup>10</sup>Be(t, p)<sup>12</sup>Be reaction at 2.73(3) MeV [11]. Recently, a long-lived 0<sup>+</sup> state at 2.24(2) MeV [13] was populated in the production of a <sup>12</sup>Be secondary beam. The mean lifetime of the state was determined to be 331(12) ns [14]. The presence of this state makes it important to investigate the *s*-wave configuration in the two bound 0<sup>+</sup> levels of <sup>12</sup>Be.

The spectroscopic factor (*S*) for the <sup>11</sup>Be<sub>gs</sub> +  $n(2s_{1/2})$  configuration in <sup>12</sup>Be<sub>gs</sub> was found to be 0.42 ± 0.06 from a one-neutron removal reaction [15]. This is smaller than the  $2s_{1/2}$  component in <sup>11</sup>Be as expected due to larger neutron separation energy. The <sup>12</sup>Be<sub>gs</sub> has *p*-wave and *d*-wave spectroscopic factors of 0.37 ± 0.06 [15] and 0.48 ± 0.06 [16], respectively. The small *s*-wave probability in <sup>12</sup>Be makes it necessary to explore whether any of the bound excited levels have a <sup>11</sup>Be<sub>gs</sub> +  $n(2s_{1/2})$  configuration. It should be mentioned here that in the neutron removal reactions the <sup>12</sup>Be beam cannot be distinguished as being in its ground state or its  $0_2^+$  long-lived state. It is therefore of utmost importance to find the *s*-wave strength through a different method that is free from this problem. In this work, the <sup>11</sup>Be(*d*, *p*)<sup>12</sup>Be reaction provides the first clean signature of the *s*-wave strength for the ground state and the  $0_2^+$  state.

Theoretically, the *s*-wave spectroscopic factor for the  $0_2^+$  level was predicted in [17] to be 1.34, but it is stated that these numbers should not be taken too seriously, since the matrix elements used are appropriate to a <sup>12</sup>C core. Ref. [18], with different wavefunctions, discusses an increased *s*-wave spectroscopic factor of 1.06 for the ground state while that for the excited  $0_2^+$  state is predicted to be 0.54.<sup>1</sup> Models of <sup>11,12</sup>Be based on a deformed potential [19] consider the ground state and the  $0_2^+$  excited state to be linear combinations of  $1/2[220]^2$  and  $1/2[101]^2$  two-neutron configurations with equal amplitudes. A study based on a <sup>10</sup>Be + *n* + *n* 



**Fig. 1.** (a) The kinematic loci of the protons identified in the upstream silicon detector in coincidence with <sup>12</sup>Be in the downstream silicon detector. (b) The Q-value spectrum integrated over the full angular range. The different states of <sup>12</sup>Be are labeled in the figure.

three-body structure model discussed the second  $0^+$  excited state to have two neutrons in the  $1p_{1/2}$  orbital [20]. No experimental information exists on the detailed configuration of the excited states in <sup>12</sup>Be.

The experiment was performed at the ISAC-II facility, TRIUMF, Canada. The radioactive <sup>11</sup>Be beam was extracted from the TRILIS laser ion source [21], and accelerated to 5 A MeV by the ISAC-I room-temperature RFQ and DTL [22] accelerators, followed by the new ISAC-II superconducting linear accelerator. The uncertainty in the average beam energy was < 10 A keV. An accelerated beam intensity of typically  $10^5$  ions/s was delivered to a 40 µg/cm<sup>2</sup> self supporting (CD<sub>2</sub>)<sub>n</sub> reaction target whose thickness was measured from energy-loss using a standard alpha source. The beam spot size at the target was < 3 mm in diameter.

A 140 µm thick annular S3 type double-sided silicon strip detector detected the protons in the backward direction covering laboratory angles of  $130^{\circ}$ – $160^{\circ}$ . The 24 ring segments of the detector defined the scattering angle. The reverse side of the detector is segmented azimuthally into 32 sectors. The scattered <sup>12</sup>Be along with <sup>11</sup>Be from elastic scattering were detected in another annular silicon detector placed 75 cm downstream of the reaction target with laboratory angular coverage of  $0.8^{\circ}$ – $2.7^{\circ}$ . This coverage ensured detection of the protons and <sup>12</sup>Be nuclei in coincidence. The coincident detection suppressed nearly all the background events making it possible to identify the kinematic loci for the <sup>11</sup>Be(*d*, *p*)<sup>12</sup>Be reaction through proton energy and angle correlations (Fig. 1a). The coincidence efficiency was determined from a Monte Carlo simulation and the region with efficiency  $\ge$  70% is shown in Fig. 3.

The Q-value spectrum constructed from the energy and angle of the detected protons is shown in Fig. 1b integrated over the

<sup>&</sup>lt;sup>1</sup> This value is considering the  $s^2$  wavefunction probability to be 0.27 instead of 0.17 which is likely a misprint in [18] because it does not lead to a total probability of unity.

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