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Dislocation nucleation from symmetric tilt grain boundaries in body-centered cubic vanadium

Shuozhi Xu [a,*], Yanqing Su [b]

[a] California NanoSystems Institute, University of California, Santa Barbara, Santa Barbara, CA 93106-6105, USA
[b] Department of Mechanical Engineering, University of California, Santa Barbara, Santa Barbara, CA 93106-5070, USA

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ABSTRACT

We perform molecular dynamics (MD) simulations with two interatomic potentials to study dislocation nucleation from six symmetric tilt grain boundaries (GB) using bicrystal models in body-centered cubic vanadium. The influences of the misorientation angle are explored in the context of activated slip systems, critical resolved shear stress (CRSS), and GB energy. It is found that for four GBs, the activated slip systems are not those with the highest Schmid factor, i.e., the Schmid law breaks down. For all misorientation angles, the bicrystal is associated with a lower CRSS than their single crystalline counterparts. Moreover, the GB energy decreases in compressive loading at the yield point with respect to the undeformed configuration, in contrast to tensile loading.

1. Introduction

Much insight has been provided into the inelastic deformation behavior of nanocrystalline (NC) materials following the landmark paper by Gleiter [1]. Experiments showed that artifact-free bulk NC materials with a narrow grain size (< 100 nm) distribution exhibit tensile yield strength several times higher than that of their coarse-grained (CG) counterparts [2]. Unlike CG metals which are considered "yield" when a large number of dislocations are simultaneously transmitted across the grain boundaries (GBs) [3,4], the strength of an NC metal is mainly controlled by dislocation nucleation from GBs as a result of its inability to accomondate long range dislocation pile-up in individual grains. Indeed, GBs are especially important in NC metals due to the latter's high GB area to material volume ratio. In the last three decades, numerous experimental studies [5,6], continuum-based modeling [7], and atomisitc simulations [8,9] have been devoted to exploring plastically deformed NC metals, pointing to the need to understand the role played by dislocation nucleation from GBs in governing mechanical properties of NC materials [10].

As an NC metal contains a large number of GBs of different types and complex triple junctions, direct simulations of the NC metals are not desirable if one were to focus on responses of individual GBs. To isolate a GB from other lattice defects, the idealized sample containing only one GB, i.e., a bicrystal, is frequently used in experiments [11], multiscale modeling [12], and atomistic simulations [13]. It is found that besides the five macroscopic degrees of freedom (DOFs) that dictate a GB at the macroscopic level, atomic-level parameters such as the three microscopic DOFs, inter-atom deletion, local lattice rotation, GB ledges, interfacial porosity, and free volume spatial distribution may also significantly influence GB structures and corresponding dislocation nucleation [14]. When a Cu bicrystal is subject to a tensile loading, the natural conformation of the interface porosity with respect to the primary dislocation slip systems is responsible for the easy emission of Shockley partial dislocations from GBs containing a certain structural unit (SU) [15]. Of note is that most atomistic simulations of bicrystals so far considered face-centered cubic systems [16,17], while body-centered cubic (BCC) lattices which have more complicated plastic deformation mechanisms [18–21] are much less explored. On the other hand, investigations of dislocation nucleation from GBs in BCC systems are necessitated by the need to understand the plasticity of relevant NC materials.

Therefore, in this work, we utilize molecular dynamics (MD) simulations to investigate dislocation nucleation from six $\langle 112 \rangle$ symmetric tilt GBs subject to compressive loading along the GB plane normal direction. BCC vanadium (V) is chosen as the model material because V alloys are among the primary candidate first wall structural materials in fusion reactors [22] yet few prior MD simulations were devoted to the GBs in V or V alloys. Six single crystalline counterparts are also examined for reference. The remainder of the paper is organized as follows. It starts in Section 2 which details the process to create desired atomistic bicrys-

* Corresponding author.
E-mail address: shuozhixu@ucsb.edu (S. Xu).

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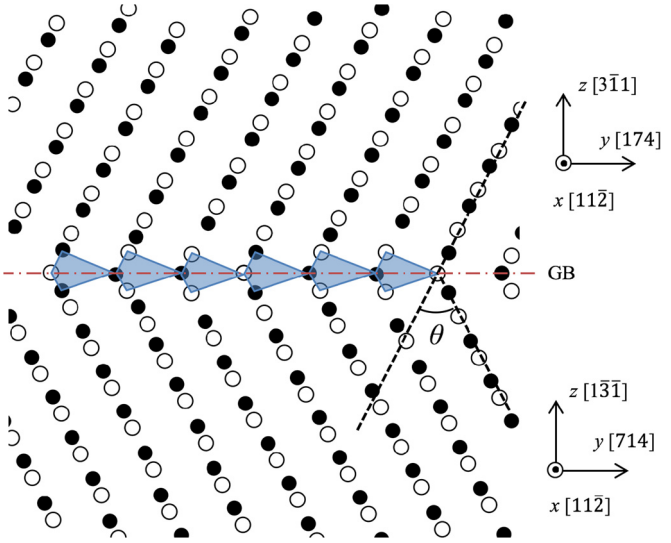

Fig. 1. A schematic of a bicrystal model with $\Sigma 11(1\bar{3}\bar{1})$ symmetric tilt GB, where the B SUs are highlighted in blue. Atoms on adjacent $(11\bar{2})$ planes are black and white, respectively. The misorientation angle θ is 62.96° in this case but needs to be varied for other GBs. (For interpretation of the colors in the figure(s), the reader is referred to the web version of this article.)

tal models. Then, in Section 3, the impacts of misorientation angle on slip system, yield stress, and GB energy are analyzed. Main findings are summarized in Section 4.

2. Methodology

A schematic of a bicrystal model is shown in Fig. 1, with the two grains rotated with respect to each other by a misorientation angle θ around a common $x[11\bar{2}]$ axis. The transformation matrix between the two crystallographic orientation systems can be written as a function of θ which is assumed positive for counterclockwise rotation, i.e.,

$$\begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\frac{\theta}{2} & \sin\frac{\theta}{2} \\ 0 & -\sin\frac{\theta}{2} & \cos\frac{\theta}{2} \end{bmatrix}$$

where θ varies from 34.04°, 62.96°, 85.14°, 101.54°, 122.88°, to 135.58° to produce six coincident site lattice (CSL) GBs with relatively low Σ, as listed in Table 1. Periodic boundary conditions are applied on all three axes, which effectively results in two identical, parallel, infinitely large GBs in one supercell [23]. The edge lengths of the simulation cell along each direction are set such that $L_x = 158.69$ Å, while L_y and L_z equal 32 times the corresponding lattice periodicity length. Since the atomistic simulation results significantly rely on the interatomic potential [24,25], two semi-empirical potentials — a modified embedded atom method (MEAM) potential [26] and a Finnis–Sinclair (FS) type potential [27] — are adopted for interactions between V atoms in a BCC lattice, with a lattice constant a_0 of 3.03686 Å and 3.03 Å, respectively.

As mentioned earlier, altering the five macroscopic DOFs that determine the CSL designation and GB plane does not result in unique GB structures [28]. To achieve the most probable equilibrium GB structures as in experiments, for each θ, 100 initial configurations considering a variety of in-plane (i.e., the x–y plane) rigid body translations and atom deletion criteria are created. For each configuration, energy minimization using a conjugate gradient algorithm is conducted [29,30], and the GB energy E_{GB} is calculated by

$$E_{GB} = \frac{E_{tot} - NE_{coh}}{2A_{GB}} \tag{1}$$

where E_{tot} is the total potential energy, N is the total number of atoms given in the last column of Table 1, $A_{GB} = L_x L_y$ is the GB area in one supercell, and E_{coh}, the cohesive energy of a V atom in a perfect lattice, is −5.3 eV for both potentials. It follows that, for each θ, the GB with the lowest E_{GB} among all 100 configurations is considered the equilibrium one [31] and is subject to homogeneous compressive loading in the plane stress condition along the z axis at a constant strain rate 10^9 s^{-1}. An NPT ensemble is adopted to maintain the temperature at 10 K and to zero the stress tensor components associated with the x and y directions. We remark that tensile simulation results of Cu bicrystals attained by high strain rates dynamic deformation at 10 K are similar to those by molecular statics at 0 K [13]; hence, the temperature we use is sufficiently low such that the thermal and mechanical effects do not intertwin. Compressive loading of six single crystals with the same crystal orientations as the lower grain in each case is also conducted for reference. All atomistic simulations are performed using LAMMPS [32] and atomic configurations are visualized in OVITO [33] with the lattice defects identified by the centrosymmetry parameter (CSP) [34].

3. Results and discussion

In all models, homogeneous dislocations are nucleated within the single crystals while dislocations are nucleated from GBs in the bicrystals. All dislocations are on {110} planes, with $(a_0/2)\langle 111\rangle$ Burgers vector. Snapshots of dislocation nucleation from $\Sigma 11(1\bar{3}\bar{1})$ and $\Sigma 35(\bar{1}5\bar{3})$ GBs based on the MEAM potential [26] are shown in Fig. 2, while the FS potential [27] predicts the same activated slip systems. We also calculate the Schmid factor (SF) and the normal factor (NF) for the activated slip systems because they both have a significant effect on the dislocation nucleation from GBs [14,35,36]; specifically,

$$\mathrm{SF} = \cos\phi\cos\lambda \tag{2}$$

$$\mathrm{NF} = \cos\phi\sin\lambda \tag{3}$$

where ϕ is the angle formed between the loading axis (i.e., the z axis) and the slip plane normal direction, and λ is the angle be-

Table 1
The misorientation angle θ, CSL designation Σ, GB plane normal in terms of lower grain in Fig. 1, simulation cell size, and number of atoms for the six bicrystals.

Misorientation angle θ (°)	CSL designation and GB plane normal	Simulation cell size (Å)			Number of atoms
		L_x	L_y	L_z	
34.04	$\Sigma 35(3\bar{5}\bar{1})$	158.69	160.95	147.84	258462
62.96	$\Sigma 11(1\bar{3}\bar{1})$	158.69	143.54	146.50	227779
85.14	$\Sigma 59(1\bar{7}\bar{3})$	158.69	134.29	139.17	202356
101.54	$\Sigma 5(0\bar{2}\bar{1})$	158.69	141.94	130.38	201402
122.88	$\Sigma 35(\bar{1}5\bar{3})$	158.69	147.53	147.84	236278
135.58	$\Sigma 7(\bar{1}3\bar{2})$	158.69	148.44	155.83	251509

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