



Contents lists available at ScienceDirect

Physics Letters A

www.elsevier.com/locate/pla



Sub- and super-diffusion on Cantor sets: Beyond the paradox

Alireza K. Golmankhaneh^{a,*}, Alexander S. Balankin^b

^a Department of Physics, Urmia Branch, Islamic Azad University, Urmia, Iran

^b Grupo "Mecánica Fractal", ESIME-Zacatenco, Instituto Politécnico Nacional, México D.F. 07738, Mexico

ARTICLE INFO

Article history:

Received 13 November 2017

Received in revised form 2 January 2018

Accepted 5 February 2018

Available online xxxx

Communicated by C.R. Doering

Keywords:

Anomalous diffusion

Random walk

Middle- ϵ Cantor set

F^α -measure

Spectral dimension

ABSTRACT

There is no way to build a nontrivial Markov process having continuous trajectories on a totally disconnected fractal embedded in the Euclidean space. Accordingly, in order to delineate the diffusion process on the totally disconnected fractal, one needs to relax the continuum requirement. Consequently, a diffusion process depends on how the continuum requirement is handled. This explains the emergence of different types of anomalous diffusion on the same totally disconnected set. In this regard, we argue that the number of effective spatial degrees of freedom of a random walker on the totally disconnected Cantor set is equal to $n_{sp} = [D] + 1$, where $[D]$ is the integer part of the Hausdorff dimension of the Cantor set. Conversely, the number of effective dynamical degrees of freedom (d_s) depends on the definition of a Markov process on the totally disconnected Cantor set embedded in the Euclidean space E^n ($n \geq n_{sp}$). This allows us to deduce the equation of diffusion by employing the local differential operators on the F^α -support. The exact solutions of this equation are obtained on the middle- ϵ Cantor sets for different kinds of the Markovian random processes. The relation of our findings to physical phenomena observed in complex systems is highlighted.

© 2018 Elsevier B.V. All rights reserved.

1. Introduction

Diffusion phenomena play an important role in physics, chemistry, technological processes, and biology [1–4]. Classical diffusion is associated with a Brownian motion. The laws of Brownian motion crucially rely on the hypothesis that the steps of a random walker are small (with finite variance) and uncorrelated. Whenever these assumptions are violated, the standard diffusion picture breaks down, leading to emergence of an anomalous diffusion [5]. In contrast to the Gaussian diffusion, which is characterized by the linear dependence of the mean square displacement on the time, the anomalous diffusion is non-universal in that it involves a parameter called the fractal dimension of a random walk [2]. Accordingly, for long enough times the mean square displacement of an ensemble of walkers behaves asymptotically as

$$\langle r^2 \rangle = K_\gamma t^\gamma, \quad (1)$$

where K_γ is the generalized diffusion coefficient and $\gamma = 2/D_w$ is the diffusion exponent. In the case of normal diffusion $\gamma = 1$ whereas the sub- and super-diffusion are characterized by $\gamma < 1$

and $\gamma > 1$ respectively [1–5]. The generalized Einstein law provides the relation between the fractal dimension of random walk D_w , electric resistance exponent ζ and spectral dimension of medium d_s , which reads as $D_w = 2\zeta/(2 - d_s)$ [6]. Furthermore, many types of fractals (but not all) obey the Alexander–Orbach relation

$$D_w = \frac{2D}{d_s}, \quad (2)$$

where D is the fractal (e.g. Hausdorff, box-counting, or self-similarity) dimension, such that $\zeta = D_w - D$ [5,6]. For path-connected fractals $d_s \leq D$ and so $\gamma \leq 1$ while $D_w \geq 2$ [7]. Accordingly, the random walk on path-connected fractals has become a paradigmatic model for the sub-diffusion observed in a great variety of physical systems (see Refs. [1–9] and references therein). On the other hand, the super-diffusion is commonly associated with Lévy flights [9,10]. At the same time, characteristic features of many physical systems can be modeled using totally disconnected fractals, e.g. Cantor sets embedded in E^n [11–14]. The ternary Cantor set of real numbers was introduced by Georg Cantor in order to illustrate the statement that a perfect set does not need to be everywhere dense. Geometrically, it is constructed by iterative deletion of open middle-third intervals from remaining intervals of the previous iteration, starting from the unit interval $[0, 1]$ ad infinitum. The ternary Cantor set is a self-similar, totally discon-

* Corresponding author.

E-mail addresses: alirezakhalili2002@yahoo.co.in, akhalilig1700@alumno.ipn.mx (A.K. Golmankhaneh), abalankin@ipn.mx (A.S. Balankin).

<https://doi.org/10.1016/j.physleta.2018.02.009>

0375-9601/© 2018 Elsevier B.V. All rights reserved.

1 nected, bounded, compact, perfect set which is nowhere dense on
 2 $[0, 1]$. Its Hausdorff dimension is equal to the similarity dimension
 3 $D = \ln 2 / \ln 3$. The middle- ϵ Cantor sets with the Hausdorff dimen-
 4 sion

$$5 \quad D = \frac{\ln 2}{\ln 2 - \ln(1 - \epsilon)} \quad (3)$$

6 can be constructed by the iterative deletion of open intervals of
 7 length equal to the ϵ fraction of the length of intervals remain-
 8 ing after the previous iteration [15]. The Cantor set in E^n (also
 9 called the Cantor dust) with the Hausdorff dimension $0 < D < n$
 10 can be constructed either by similar iterative method, starting from
 11 n -dimensional unit cube $[0, 1]^n$, or as the Cartesian product of
 12 n orthogonal Cantor sets of the same or different Hausdorff dimen-
 13 sions. Every Cantor set (dust) is homeomorphic to the ternary
 14 Cantor set [16]. Furthermore, it has been proved (see Ref. [17]) that
 15 any compact metric space is a continuous image of the ternary
 16 Cantor set. Due to these remarkable properties the totally discon-
 17 nected Cantor sets have found a celebrated place in mathematical
 18 analysis and its applications. In particular, stochastic processes on
 19 Cantor sets are an active topic of research [18–31]. In this regard,
 20 the ternary Cantor set is interesting because the diffusion on it ex-
 21 hibits some apparent paradoxes. Specifically, in some works, (see
 22 Refs. [18–24]) it was suggested that the Markov processes defined
 23 on the totally disconnected Cantor set lead to the super-diffusion
 24 allied with the set's self-similarity. Conversely, it was argued (see
 25 [28–31]) that the ultrametric properties of the totally disconnected
 26 Cantor sets stipulate the sub-diffusive nature of the Markov pro-
 27 cesses. It has been also recognized that local and average measure-
 28 ments can display different asymptotic behavior [12,33]. Accord-
 29 ingly, alternative models of the anomalous diffusion on the ternary
 30 Cantor set were suggested and disputed in the literature (see, for
 31 example, Refs. [32–36] and references therein). In this work, we
 32 argue that diffusion on the totally disconnected Cantor set can be
 33 either the super-, or the sub-diffusion, depending on how the ran-
 34 dom walk is constructed. In this regard, we stress that the number
 35 of effective dynamical degree of freedom of a random walker on
 36 the Cantor set depends on the Markov process definition, while
 37 the number of effective spatial degree of freedom is equal to the
 38 minimum dimension of the Euclidean space in which the Cantor
 39 set can be bi-Lipschitz embedded. This explains the emergence of
 40 different types of the anomalous diffusion on the ternary Cantor
 41 set. The diffusion equation describing different types of diffusion
 42 is deduced and solved for the middle- ϵ Cantor sets. The phys-
 43 ical implications of these findings are discussed. The rest of the
 44 paper is organized as follows. Sec. 2 is devoted to the analysis
 45 of different Markov processes on the totally disconnected Cantor
 46 set. The processes associated with the sub- and super diffusion on
 47 the ternary Cantor set are outlined. The predicted values for the
 48 diffusion exponents are compared with available data from numer-
 49 ical simulations of Lévy flights and walks on the Cantor sets. In
 50 Sec. 3 the equation of diffusion is deduced using the algorithmic
 51 F^α -calculus based on the Riemannian-like method. The relations
 52 between the orders of F^α -derivatives and the numbers of the ef-
 53 fective degrees of freedom are established. The probability density
 54 distributions functions are found for different types of diffusion
 55 on the middle- ϵ Cantor sets. Some physical phenomena associ-
 56 ated with different types of anomalous diffusion on the Cantor sets
 57 are highlighted. The main findings and conclusions are outlined in
 58 Sec. 4.

62 2. Markov processes on Cantor sets

63 It is a straightforward matter to understand that there is no
 64 way to build a nontrivial Markov process having continuous tra-
 65 jectories on a totally disconnected Cantor set embedded in the

66 Euclidean space E^n . Therefore, in order to delineate the diffu-
 67 sion process, one needs to relax the continuum requirement. In
 68 this way, different properties of the ternary Cantor set were ex-
 69 plored to define the Markov random processes on it. Specifically,
 70 the Brownian motion on the totally disconnected fractal set was
 71 introduced by Fujita [18]. This model was used to determine the
 72 spectral dimension of the Cantor set from the growth order of
 73 eigenvalues of the Brownian motion generator. Further, the tran-
 74 sition probability densities for generalized one-dimensional diffu-
 75 sion processes were studied in [19]. Later, Takahashi and Tamura
 76 [20] have defined diffusion processes on totally disconnected self-
 77 similar fractal sets as the limits of suitably scaled random walk.
 78 On the other hand, Evans [21] has introduced Markov processes
 79 with stationary independent increments taking values in a non-
 80 discrete, locally compact, metrizable, totally disconnected Abelian
 81 Cantor group. Further, Aldous and Evans [22] have used the Dirich-
 82 let form methods to construct and analyze a general class of re-
 83 versible Markov processes with totally disconnected state space.
 84 At the same time, Freiberg [23] has introduced the measure theo-
 85 retic Dirichlet forms on compact subsets of the real line. Using the
 86 technique of Dirichlet–Neumann-bracketing, he has obtained the
 87 estimates of the eigenvalue counting functions of the associated
 88 measure geometric Laplacians. Using the Dirichlet form technique,
 89 Karwowski [24] has developed a model of diffusion on the real
 90 line with jumps on the Cantor set which preserve the ultrametric
 91 feature of random process on unit ball of 2-adic numbers. In this
 92 regard, it is pertinent to note that the Markov processes studied in
 93 Refs. [18–24] lead to the super-diffusion on the ternary Cantor set.

94 Alternatively, Lobus [25] has constructed a strong Markov pro-
 95 cesses on Cantor sets from the Wiener processes by means of time
 96 change, killing, and space transformation. Bhamidi et al. [26] have
 97 proposed an analogue of Brownian motion that has as its state
 98 space an arbitrary closed subset of the line that is unbounded
 99 above and below. It was shown that there is a unique such process,
 100 which turns out to be automatically a reversible Feller–Dynkin
 101 Markov process. This process is martingale and has the identity
 102 function as its quadratic variation process. So, it is continuous
 103 in the sense that its sample paths don't skip over points. It was
 104 also recognized that the totally disconnected Cantor sets exhibit
 105 a natural ultrametric structure [27,28]. An analogue of the Rie-
 106 mannian geometry for an ultrametric Cantor set was suggested in
 107 Ref. [29]. Consequently, Pearson and Bellissard [29] have defined
 108 a probability measure, Dirichlet forms, and associated analogue of
 109 the Laplace–Beltrami operator. In this way, it was deduced that
 110 the diffusion on the ultrametric Cantor set exhibits nonstandard
 111 sub-diffusive behavior, as $t \rightarrow 0$. Later, Bakhtin [30] has defined
 112 an analog of Brownian motion on the triadic Cantor set by intro-
 113 ducing requirements on the Markov semigroup, most important of
 114 which are isometry invariance, and scale invariance. This allows to
 115 build up the jump statistics, while the generators of the symmet-
 116 ric self-similar Markov process play the role of the Laplacian on
 117 the Cantor set. More recently, Kigami [31] has shown that a tran-
 118 sient random walk on a tree induces a Dirichlet form on its Martin
 119 boundary, which is the Cantor set. This was used to define an in-
 120 trinsic metric on the Cantor set associated with the random walk.
 121 Accordingly, the harmonic measure and induced Dirichlet form on
 122 the Cantor set were explicitly expressed in terms of the effec-
 123 tive resistances. Consequently, the asymptotic behaviors of the heat
 124 kernel, the jump kernel, and moments of displacements were ob-
 125 tained. In this regard, it is pertinent to point out that, in contrast
 126 to Markov processes studied in Refs. [18–24], the Markov processes
 127 defined in Refs. [25–31] lead to the sub-diffusion on the ternary
 128 Cantor set.

129 In this background, we stress that, generally, the type of dif-
 130 fusion is determined by the numbers of effective dynamical and
 131 spatial degrees of freedom of a random walker (see Refs. [37,38]).
 132

Download English Version:

<https://daneshyari.com/en/article/8203687>

Download Persian Version:

<https://daneshyari.com/article/8203687>

[Daneshyari.com](https://daneshyari.com)