# Abundant lumps and their interaction solutions of (3+1)-dimensional linear PDEs 

Wen-Xiu Ma*<br>Department of Mathematics, Zhejiang Normal University, Jinhua 321004, Zhejiang, China<br>Department of Mathematics and Statistics, University of South Florida, Tampa, FL 33620, USA<br>College of Mathematics and Systems Science, Shandong University of Science and Technology, Qingdao 266590, Shandong, China<br>College of Mathematics and Physics, Shanghai University of Electric Power, Shanghai 200090, China<br>International Institute for Symmetry Analysis and Mathematical Modelling, Department of Mathematical Sciences, North-West University, Mafikeng Campus, Private Bag X2046, Mmabatho 2735, South Africa

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#### Abstract

The paper aims to explore the existence of diverse lump and interaction solutions to linear partial differential equations in (3+1)-dimensions. The remarkable richness of exact solutions to a class of linear partial differential equations in (3+1)-dimensions will be exhibited through Maple symbolic computations, which yields exact lump, lump-periodic and lump-soliton solutions. The results expand the understanding of lump, freak wave and breather solutions and their interaction solutions in soliton theory.


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## 1. Introduction

Lump solutions are a particular kind of exact solutions, which describe various important nonlinear phenomena in nature [1,2]. More specifically, such solutions can be generated from solitons by taking long wave limits [3]. There are also positons and complexitons to integrable equations, enriching the diversity of solitons [4,5]. Interaction solutions between two different kinds of exact solutions exhibit more diverse nonlinear phenomena [6].

Soliton solutions are exponentially localized in all directions in space and time, and lump solutions, rationally localized in all directions in space. Through a Hirota bilinear form:

$$
\begin{equation*}
P\left(D_{x}, D_{t}\right) f \cdot f=0 \tag{1.1}
\end{equation*}
$$

where $P$ is a polynomial and $D_{x}$ and $D_{t}$ are Hirota's bilinear derivatives, an $N$-soliton solution in (1+1)-dimensions can be defined by

$$
\begin{equation*}
f=\sum_{\mu=0,1} \exp \left(\sum_{i=1}^{N} \mu_{i} \xi_{i}+\sum_{i<j} \mu_{i} \mu_{j} a_{i j}\right) \tag{1.2}
\end{equation*}
$$

[^0]where
\[

\left\{$$
\begin{array}{l}
\xi_{i}=k_{i} x-\omega_{i} t+\xi_{i, 0}, 1 \leq i \leq N  \tag{1.3}\\
\mathrm{e}^{a_{i j}}=-\frac{P\left(k_{i}-k_{j}, \omega_{j}-\omega_{i}\right)}{P\left(k_{i}+k_{j}, \omega_{j}+\omega_{i}\right)}, 1 \leq i<j \leq N
\end{array}
$$\right.
\]

with $k_{i}$ and $\omega_{i}$ satisfying the dispersion relation and $\xi_{i, 0}$ being arbitrary shifts. The KPI equation

$$
\begin{equation*}
\left(u_{t}+6 u u_{x}+u_{x x x}\right)_{x}-u_{y y}=0 \tag{1.4}
\end{equation*}
$$

has a lump solution [7]:

$$
\begin{equation*}
u=2(\ln f)_{x x}, f=\left(a_{1} x+a_{2} y+a_{3} t+a_{4}\right)^{2}+\left(a_{5} x+a_{6} y+a_{7} t+a_{8}\right)^{2}+a_{9} \tag{1.5}
\end{equation*}
$$

where

$$
\begin{equation*}
a_{3}=\frac{a_{1} a_{2}^{2}-a_{1} a_{6}^{2}+2 a_{2} a_{5} a_{6}}{a_{1}^{2}+a_{5}^{2}}, a_{7}=\frac{2 a_{1} a_{2} a_{6}-a_{2}^{2} a_{5}+a_{5} a_{6}^{2}}{a_{1}^{2}+a_{5}^{2}}, a_{9}=\frac{3\left(a_{1}^{2}+a_{5}^{2}\right)^{3}}{\left(a_{1} a_{6}-a_{2} a_{5}\right)^{2}}, \tag{1.6}
\end{equation*}
$$

and the other parameters $a_{i}$ 's are arbitrary but need to satisfy $a_{1} a_{6}-a_{2} a_{5} \neq 0$, which guarantees rational localization in all directions in the ( $x, y$ )-plane. Other integrable equations, possessing lump solutions, include the three-dimensional threewave resonant interaction [8], the BKP equation [9,10], the Davey-Stewartson equation II [3], the Ishimori-I equation [11] and many others $[12,13]$.

It is recognized by making symbolic computations that many nonintegrable equations possess lump solutions as well, including ( $2+1$ )-dimensional generalized KP, BKP and Sawada-Kotera equations [14-16]. Moreover, various studies show the existence of interaction solutions between lumps and another kind of exact solutions to nonlinear integrable equation in (2+1)-dimensions, which contain lump-soliton interaction solutions (see, e.g., [17-20]) and lump-kink interaction solutions (see, e.g., [21-24]). Nevertheless, in the (3+1)-dimensional case, only lump-type solutions are presented for the integrable Jimbo-Miwa equations, which are rationally localized in almost all but not all directions in space. All presented analytical rational solutions to the (3+1)-dimensional Jimbo-Miwa equation in [25-27] and the (3+1)-dimensional Jimbo-Miwa like equation in [28] are not rationally localized in all directions in space. It is absolutely very interesting and important to explore lump and interaction solutions to partial differential equations in (3+1)-dimensions.

This paper aims at showing that there do exist abundant lump solutions and their interaction solutions to linear partial differential equations in (3+1)-dimensions. A class of particular examples in (3+1)-dimensions will be considered to exhibit such solution phenomena. We will explicitly generate lump solutions and mixed lump-periodic and lump-soliton solutions for a specially chosen class of (3+1)-dimensional linear partial differential equations. Based on Maple symbolic computations, sufficient conditions and examples of lump and interaction solutions will be provided, together with three-dimensional plots and contour plots of special examples of the presented solutions. Some concluding remarks will be given in the final section.

## 2. Abundant lump and interaction solutions

Let $u=u(x, y, z, t)$ be a real function of $x, y, z, t \in \mathbb{R}$. We consider a class of linear partial differential equations (PDEs) in (3+1)-dimensions:

$$
\begin{equation*}
\alpha_{1} u_{x y}+\alpha_{2} u_{x z}+\alpha_{3} u_{x t}+\alpha_{4} u_{y z}+\alpha_{5} u_{y t}+\alpha_{6} u_{z t}+\alpha_{7} u_{x x}+\alpha_{8} u_{y y}+\alpha_{9} u_{z z}+\alpha_{10} u_{t t}=0, \tag{2.1}
\end{equation*}
$$

where $\alpha_{i}, 1 \leq i \leq 10$, are real constants, and the subscripts denote partial differentiation.
We search for a kind of exact solutions

$$
\begin{equation*}
u=v\left(\xi_{1}, \xi_{2}, \xi_{3}, \xi_{4}\right) \tag{2.2}
\end{equation*}
$$

where $v$ is an arbitrary real function, and $\xi_{i}, 1 \leq i \leq 4$, are four wave variables:

$$
\begin{equation*}
\xi_{i}=a_{i} x+b_{i} y+c_{i} z+d_{i} t+e_{i}, \quad 1 \leq i \leq 4 \tag{2.3}
\end{equation*}
$$

in which $a_{i}, b_{i}, c_{i}, d_{i}$ and $e_{i}, 1 \leq i \leq 4$, are real constants to be determined. Then, the linear PDE (2.1) becomes

$$
\begin{equation*}
\sum_{i=1}^{4} \sum_{j=i}^{4} w_{i j} v_{\xi_{i} \xi_{j}}=0 \tag{2.4}
\end{equation*}
$$

where $w_{i j}, 1 \leq i \leq j \leq 4$, are quadratic functions of the parameters $a_{i}, b_{i}, c_{i}$ and $d_{i}, 1 \leq i \leq 4$. Upon setting all coefficients of the ten second partial derivatives of $v$ to be zero, we obtain a system of equations on the parameters:

$$
\left\{\begin{array}{l}
\alpha_{1} a_{i} b_{i}+\alpha_{2} a_{i} c_{i}+\alpha_{3} a_{i} d_{i}+\alpha_{4} b_{i} c_{i}+\alpha_{5} b_{i} d_{i}  \tag{2.5}\\
\quad+\alpha_{6} c_{i} d_{i}+\alpha_{7} a_{i}^{2}+\alpha_{8} b_{i}^{2}+\alpha_{9} c_{i}^{2}+\alpha_{10} d_{i}^{2}=0,1 \leq i \leq 4, \\
\alpha_{1}\left(a_{i} b_{j}+a_{j} b_{i}\right)+\alpha_{2}\left(a_{i} c_{j}+a_{j} c_{i}\right)+\alpha_{3}\left(a_{i} d_{j}+a_{j} d_{i}\right)+\alpha_{4}\left(b_{i} c_{j}+b_{j} c_{i}\right)+\alpha_{5}\left(b_{i} d_{j}+b_{j} d_{i}\right) \\
\quad+\alpha_{6}\left(c_{i} d_{j}+c_{j} d_{i}\right)+2 \alpha_{7} a_{i} a_{j}+2 \alpha_{8} b_{i} b_{j}+2 \alpha_{9} c_{i} c_{j}+2 \alpha_{10} d_{i} d_{j}=0,1 \leq i<j \leq 4 .
\end{array}\right.
$$

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[^0]:    * Correspondence to: Department of Mathematics and Statistics, University of South Florida, Tampa, FL 33620, USA.

    E-mail address: mawx@cas.usf.edu.

