



On the asymptotically Poincaré-Einstein 4-manifolds with harmonic curvature

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HIGHLIGHTS

- We investigate the AH mass of an APE 4-manifold with harmonic curvature.
- We prove a rigidity of complete non-compact 4-manifold with harmonic curvature and non-positive scalar curvature.
- We prove a rigidity of complete non-compact Einstein 4-manifold with non-positive scalar curvature.
- We prove a rigidity of an APE 4-manifold with harmonic curvature. The bound of the L^2 -norm of the curvature is a highlight.

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ABSTRACT

In this paper, we discuss the mass aspect tensor and the rigidity of an asymptotically Poincaré-Einstein (APE) 4-manifold with harmonic curvature. We prove that the trace-free part of the mass aspect tensor of an APE 4-manifold with harmonic curvature and normalized Einstein conformal infinity is zero. As to the rigidity, we first show that a complete noncompact Riemannian 4-manifold with harmonic curvature and positive Yamabe constant as well as a L^2 -pinching condition is Einstein. As an application, we then obtain that an APE 4-manifold with harmonic curvature and positive Yamabe constant is isometric to the hyperbolic space provided that the L^2 -norm of the traceless Ricci tensor or the Weyl tensor is small enough and the conformal infinity is a standard round 3-sphere.

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1. Introduction

Poincaré-Einstein (PE) manifolds have been deeply investigated recently because of the so called AdS/CFT correspondence proposed in the theory of quantum gravity in theoretical physics. Since an asymptotically Poincaré-Einstein (APE) manifold and a Riemannian manifold with harmonic curvature are both a kind of generalization of a PE manifold, we wonder whether an APE manifold with harmonic curvature is Einstein. In this paper, we mainly answer this question and investigate the mass aspect tensor and the rigidity of an APE 4-manifold with harmonic curvature.

Before considering the APE Riemannian manifold, we first give some basic materials about conformally compact manifold. Suppose that M^n can be realized as the interior of a smooth compact manifold \bar{M}^n with boundary $\partial\bar{M}$. A defining function τ for the boundary $\partial\bar{M}$ is a smooth function on \bar{M}^n such that $\tau > 0$ in M^n ; $\tau = 0$ on $\partial\bar{M}$; $d\tau \neq 0$ on $\partial\bar{M}$. We refer to $\partial\bar{M}$ as the boundary-at-infinity of M^n and denote it by $\partial_\infty M$.

A complete noncompact Riemannian metric g on M^n is said to be $C^{k,\mu}$ conformally compact if the compactified metric $\bar{g} = \tau^2 g$ extends to be a $C^{k,\mu}$ Riemannian metric on \bar{M}^n . If in addition, $|d\tau|_{\tau^2 g}^2 = 1$ on $\partial_\infty M$, then (M^n, g) is called $C^{k,\mu}$ asymptotically hyperbolic (AH).

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The metric $\bar{g} = \tau^2 g$ induces a metric \hat{g} on the boundary $\partial_\infty M$, and the metric g induces a conformal class of metric $[\hat{g}]$ on the boundary $\partial_\infty M$ when defining functions vary. The conformal boundary manifold $(\partial_\infty M, [\hat{g}])$ is called the conformal infinity of the conformally compact manifold (M^n, g) .

Given a C^1 AH metric g and a representative \hat{g} in $[\hat{g}]$ on the conformal boundary $\partial_\infty M$, there is a uniquely determined defining function r such that $|dr|_{r^2 g}^2 \equiv 1$ in a neighborhood of the boundary. With this choice of r , in a collar neighborhood $[0, \epsilon) \times \partial_\infty M \subset M^n$ for some $\epsilon > 0$, g takes the geodesic normal form

$$g = r^{-2} (dr^2 + g_r), \tag{1}$$

where g_r is a curve of metrics on $\partial_\infty M$ with $g_r|_{r=0} = \hat{g}$. Such defining function r is called the geodesic defining function associated with \hat{g} .

Since the sectional curvatures of an AH metric approach -1 as $r \rightarrow 0$, the Ricci curvature will approach $-(n-1)g$ at infinity. Following [1,2], we now turn our attention to the AH manifolds that have Ricci curvature sufficiently pinched near infinity.

Definition 1.1. An AH metric g on M^n is called Poincaré-Einstein (PE) if $E(g) := Ric(g) + (n-1)g$ vanishes identically. It is called asymptotically Poincaré-Einstein (APE) if $|E|_g = O(r^n)$.

Let (M^n, g) be an APE manifold and $g = r^{-2} (dr^2 + g_r)$ with $g_r|_{r=0} = \hat{g}$. Assume that g_r is sufficiently regular that its asymptotical expansion may be calculated by the APE condition. Let E_{00} denote $E(\frac{\partial}{\partial r}, \frac{\partial}{\partial r})$ and E^\perp denote the projection of E orthogonal to $\frac{\partial}{\partial r}$. A straightforward calculation shows that

$$E_{00} = -\frac{1}{2} tr_{g_r} g_r'' + \frac{1}{2r} tr_{g_r} g_r' + \frac{1}{4} |g_r'|_{g_r}^2 \tag{2}$$

and

$$E^\perp = Ric(g_r) - \frac{1}{2} g_r'' + \frac{n-2}{2r} g_r' + \frac{1}{2r} g_r tr_{g_r} g_r' + \frac{1}{2} g_r' \cdot g_r^{-1} \cdot g_r' - \frac{1}{4} g_r' tr_{g_r} g_r', \tag{3}$$

where $'$ denotes $\frac{\partial}{\partial r}$ and $g_r' \cdot g_r^{-1} \cdot g_r'$ denotes the tensor with components $(g_r)'_{ik} (g_r)^{kl} (g_r)'_{lj}$.

By similar computation as that in [3], Bahuaud, Mazzeo and Woolgar showed in [1, Proposition 2.2] that for an APE metric g , the expansion of g_r has the following form:

$$g_r = \hat{g} + g^{(2)} r^2 + (\text{even powers of } r) + g^{(n-1)} r^{n-1} + O(r^n) \tag{4}$$

when n is even; and

$$g_r = \hat{g} + g^{(2)} r^2 + (\text{even powers of } r) + (g^{(n-1)} + \tilde{g}^{(n-1)} \log r) r^{n-1} + O(r^n \log r) \tag{5}$$

when n is odd, where

- (1) $g^{(2k)}$ are determined by \hat{g} for $2k < n-1$ and $-g^{(2)}$ is the Schouten tensor $P(\hat{g})$ of \hat{g} for $n \geq 3$, i.e.

$$g^{(2)} = -P(\hat{g}) = -\frac{1}{n-2} \left(Ric(\hat{g}) - \frac{R(\hat{g})}{2(n-1)} \hat{g} \right);$$

- (2) when n is even, $g^{(n-1)}$ is undetermined but $tr_{\hat{g}} g^{(n-1)} = 0$;
- (3) when n is odd, the ambient obstruction tensor $\tilde{g}^{(n-1)}$ is determined by \hat{g} and $tr_{\hat{g}} \tilde{g}^{(n-1)} = 0$;
- (4) when n is odd, $tr_{\hat{g}} g^{(n-1)}$ is determined by the prior coefficients, but the trace-free part of $g^{(n-1)}$ is undetermined.

Notice that when n is odd, g_r will always have an expansion involving $\log r$. But as the ambient obstruction tensor $\tilde{g}^{(n-1)}$ is determined by \hat{g} , if we choose \hat{g} carefully, for example, we choose \hat{g} to be Einstein, then we will have $\tilde{g}^{(n-1)} = 0$. We consider in this paper a class of APE manifolds which have special conformal infinities, that is, the representative metric \hat{g} is Einstein, normalized such that $Ric(\hat{g}) = \lambda(n-2)\hat{g}$ with $\lambda \in \{-1, 0, 1\}$. We call such $(\partial_\infty M, [\hat{g}])$ the normalized Einstein conformal infinity.

We will also need the following notion of AH mass for an APE Riemannian manifold developed in [2].

Lemma 1.2 ([2, Lemma 2.2 and Definition 2.5]). *Let (M^n, g) be an APE Riemannian manifold with normalized Einstein conformal infinity $(\partial_\infty M, [\hat{g}])$, then g has the form*

$$g = r^{-2} \left(dr^2 + \left(1 - \frac{\lambda}{4} r^2 \right)^2 \hat{g} + \frac{1}{n-1} r^{n-1} \theta + \frac{1}{n} r^n \kappa + O(r^{n+1}) \right) \tag{6}$$

where θ is the Neumann data for g with $tr_{\hat{g}} \theta = 0$ and κ is the mass aspect tensor for g . Moreover, $tr_{\hat{g}} \kappa$ is the mass aspect function and $\int_{\partial_\infty M} tr_{\hat{g}} \kappa dv_{\hat{g}}$ is the mass for g .

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