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# SYMMETRIES AND CONSERVATION LAWS OF A NONLINEAR SIGMA MODEL WITH GRAVITINO

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ABSTRACT. We study the symmetries and invariances of a version of the action functional of the nonlinear sigma model with gravitino, as considered in [12]. The action is invariant under rescaled conformal transformations, super Weyl transformations and diffeomorphisms. In particular cases the functional possesses a degenerate supersymmetry. The corresponding conservation laws lead to a geometric interpretation of the energy-momentum tensor and supercurrent as holomorphic sections of appropriate bundles.

## 1. INTRODUCTION

The main motivation for the introduction of the two-dimensional supersymmetric nonlinear sigma model in quantum field theory, or more specifically supergravity and superstring theory, are its symmetries, see for instance [2, 5, 6, 10]. Furthermore, as argued in [14], the functional is determined by its symmetries together with suitable bounds on the order of its Euler–Lagrange equations. While supersymmetric models are usually formulated using anticommuting variables, in [12] an analogue of the two-dimensional nonlinear supersymmetric sigma model using only commuting variables was introduced. Here we would like to give a detailed geometric account of the symmetries of this purely commutative model.

We briefly recall the two-dimensional nonlinear sigma model constructed in [12]. Let  $M$  be a Riemann surface and let  $(N, h)$  be a Riemannian manifold. In the classical nonlinear sigma model, the action functional is given by the Dirichlet energy functional which is defined for a map  $\phi: M \rightarrow N$  and a Riemannian metric  $g$  on  $M$ . In our model we need to take also their superpartners into consideration. These superpartners are geometrically formulated via suitable spinor fields. To be more precise, given the Riemannian metric  $g$ , we fix a spin structure  $\xi: P_{\text{Spin}}(M, g) \rightarrow P_{\text{SO}}(M, g)$ . An irreducible representation of the Clifford algebra  $\text{Cl}_{0,2}$  induces the real spin representation  $\mu: \text{Spin}(2) \rightarrow \text{GL}(V)$ , where  $V$  is a representation space of real dimension four. The associated spinor bundle  $S_g := P_{\text{Spin}}(M, g) \times_{\mu} V$  is equipped with the canonical spinor metric  $g_s$  and the spin connection  $\nabla^s$  induced by the Levi-Civita connection on  $TM$ . Choosing an isomorphism between the vector spaces  $V$  and  $\text{Cl}_{0,2}$  we get a Clifford map  $\gamma: TM \rightarrow \text{End}(S_g)$  which satisfies the Clifford relation

$$\gamma(X)\gamma(Y) + \gamma(Y)\gamma(X) = -2g(X, Y), \quad \forall X, Y \in \Gamma(TM).$$

Sections of  $S_g$  will be referred to as (pure) spinors, which describe matter fields with half-integer spins in physics. The spin Dirac operator  $\not{D}_g := \gamma(e_\alpha)\nabla_{e_\alpha}^s: \Gamma(S_g) \rightarrow \Gamma(S_g)$  is a first-order elliptic differential operator, which is self-adjoint if  $M$  is closed.

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