



# Fringing field can prevent infinite time quenching



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## ARTICLE INFO

### Article history:

Received 12 March 2015

Received in revised form 25 October 2015

Accepted 18 December 2015

Available online 11 January 2016

### Keywords:

Quenching

Singular nonlinearity

Fringing field

Gradient quadratic growth

MEMS

Global convergence

Convex nonlinearity

## ABSTRACT

Consider the following problem arising in Micro-Electro-Mechanical Systems (MEMS)

$$\begin{cases} u_t - \Delta u = \frac{\lambda(1 + \delta |\nabla u|^2)}{(1 - u)^p}, & (x, t) \in \Omega \times (0, T), \\ u = 0, & (x, t) \in \partial\Omega \times (0, T), \\ u(x, 0) = u_0(x), \quad 0 \leq u_0(x) < 1, & x \in \Omega, \end{cases}$$

where  $\delta > 0$ ,  $p > 1$  and  $\Omega$  is a bounded smooth domain in  $\mathbb{R}^N (N \geq 1)$ . We prove that infinite time quenching is impossible for any  $\lambda > 0$  in this problem. It provides a remarkable contrast to the case of  $\delta = 0$ , in which infinite time quenching must happen for some  $\lambda$  when  $\Omega$  is a ball in  $\mathbb{R}^N (N \geq 8)$ . This means that the presence of the fringing field  $\delta|\nabla u|^2$  dramatically changes the quenching behavior of the solution. We also obtain some new results about global convergence and quenching in finite time.

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## 1. Introduction

Consider the following nonlinear parabolic problem

$$\begin{cases} u_t - \Delta u = \frac{\lambda(1 + \delta |\nabla u|^2)}{(1 - u)^p}, & (x, t) \in \Omega \times (0, T), \\ u = 0, & (x, t) \in \partial\Omega \times (0, T), \\ u(x, 0) = u_0(x), \quad 0 \leq u_0(x) < 1, & x \in \Omega, \end{cases} \quad (1.1)$$

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where  $\lambda > 0, \delta > 0, p > 1, \Omega$  is a bounded smooth domain in  $\mathbb{R}^N (N \geq 1)$  and  $u_0 \in H_0^1(\Omega)$ . The corresponding stationary problem is

$$\begin{cases} -\Delta u &= \frac{\lambda(1 + \delta|\nabla u|^2)}{(1-u)^p}, & x \in \Omega, \\ u &= 0, & x \in \partial\Omega, \end{cases} \tag{1.2}$$

modeling an electrostatically actuated Micro-Electro-Mechanical System (MEMS) device. In this model,  $\lambda$  represents the applied voltage, the term  $\delta|\nabla u|^2$  describes the effect of a fringing electrostatic field, and  $\delta$  is the aspect ratio of the device. About details, we refer to [1,2]. Notice that  $p = 2$  (the inverse square relation in Coulomb’s Law) and the physically relevant dimensions –  $N = 1, 2$  – are the main concerns of MEMS researchers.

In this paper, we are concerned with quenching phenomena for problem (1.1). The solution of (1.1) is said to *quench* at time  $T$  if  $\max_{\bar{\Omega}} u(x, t) \rightarrow 1$  as  $t \rightarrow T$ . If  $T = +\infty$ , then  $u$  quenches in *infinite* time; if  $T < +\infty$ , then  $u$  quenches in *finite* time (see e.g. [3–7]).

We recall some related results about (1.1) and (1.2) with  $\delta = 0$ . When  $\Omega$  is a ball of  $\mathbb{R}^N (N \geq 2)$ , (1.2) has been well investigated by Joseph and Lundgren [8] (also see Fila et al. [9] and Zheng [10]); for  $N = 1$ , see Laetsch [11] or Levine [7]. A striking feature is the dependence of bifurcation diagrams on the dimension of space (see e.g. Fig. 1). The critical value  $\lambda^*$  of  $\lambda$ , also referred to as the *extremal parameter*, depends on  $p$  and  $\Omega$ .  $\lambda^*$  is also intimately associated with quenching phenomena of (1.1). Acker and Walter [3] proved, under our notations, that for  $u_0(x) \equiv 0$  and  $\Omega$  is a ball, if  $\lambda < \lambda^*$ , the solution of (1.1) globally exists, while if  $\lambda > \lambda^*$ , the solution must quench in finite time; these conclusions are still true for  $p = 2$  and general bounded domains (see [12]). For  $\lambda = \lambda^*$ , the situation is more complicated. Fila et al. [9] gave a complete answer for  $u_0(x) \equiv 0$  and  $\Omega$  being a ball (also see [10] for  $0 \leq u_0(x) < 1$ ). They proved, under our notations, that infinite-time quenching occurs at  $\lambda^*$  if and only if  $3 \leq N \leq 9$  and  $0 < p \leq \frac{N-2\sqrt{N-1}}{4-N+2\sqrt{N-1}}$  or  $N > 9$ . For example, when  $p = 2$ , the critical value of  $N$  is  $\frac{14+6\sqrt{6}}{3} \approx 7.933$ , this means that when  $\lambda = \lambda^*$ , infinite-time quenching occurs just for all dimensions  $N \geq 8$ ; see Fig. 1 for a comparison. In [4], Fila and Kawohl showed that when  $\lambda = \lambda^*$  and  $\Omega$  is a bounded convex domain, quenching in infinite time cannot occur if  $N = 2, p > 1$  or  $N = 3, p > 3$ . For  $p = 2$  and general bounded domains  $\Omega$ , Esposito et al. [12] proved that when  $\lambda = \lambda^*$ , infinite time quenching cannot occur for  $N \leq 7$ ; on the other hand, if  $N \geq 8$  and the domain  $\Omega$  is not a ball, the situation is still unclear. Besides, when  $2 \leq N \leq 7$  and  $\lambda = \lambda_*$  (see Fig. 1), the possibility of infinite time quenching is still not excluded for general initial data.

We now come back to the case  $\delta > 0$ . For  $N = 1$ , bifurcation diagrams of stationary problem (1.2) are similar to those of  $\delta = 0$ ; see Fila et al. [13] for  $p = 1$  and the authors [14] for general  $p > 0$  (see Fig. 2). For  $N \geq 2$ , Wei and Ye [15] proved an interesting result that when  $p > 1$ , there exists a number  $\lambda^* > 0$  such that (1.2) has at least two solutions for  $\lambda \in (0, \lambda^*)$ , exactly one solution for  $\lambda = \lambda^*$ , and no solutions for  $\lambda \in (\lambda^*, +\infty)$ ; also see [1,2] for more related results of numerical simulations and [16] for another proof. For  $N = 1$  and  $p = 2$ , (1.1) was investigated by Liu and Wang [17] and the results were similar to those when  $\delta = 0$ . Recently, Wang [18] considered a slightly more general problem than (1.1) and proved that for  $u_0(x) \equiv 0$ , if  $\lambda < \lambda^*$ , the solution converges pointwise to the minimal stationary solution; if  $\lambda > \lambda^*$ , the solution quenches in finite time. In a very recent work, Luo and Yau [22] gave another proof for Wang’s result on Eq. (1.1), covering the case  $\lambda = \lambda^*$ .

Notice that Wei and Ye’s result implies that the extremal solution of (1.2) at  $\lambda^*$  is always classical and the regularity is independent of the dimension  $N$ . It is natural to conjecture that there exists a counterpart in the corresponding parabolic problem (1.1). Namely, whether the quadratic gradient term  $\delta|\nabla u|^2$  (i.e. the fringing field) can remove the effect of the dimension on quenching at  $\lambda^*$ ?

The purpose of this paper is to give an affirmative answer for this problem. We prove that for any  $\lambda > 0$  and  $0 \leq u_0 < 1$ , infinite time quenching cannot occur for problem (1.1). In other words, the solution  $u$

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