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Large deviation for the nonlocal Kuramoto–Sivashinsky SPDE

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a r t i c l e i n f o

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1. Introduction

It is known that the deterministic 1-dimensional Kuramoto–Sivashinsky (K–S) equation,

$$
du(t) + (u_{xxxx} + u_{xx} + uu_x) dt = 0
$$
\n(1)

arises in the modelling of the flow of a thin film of viscous liquid falling down an inclined plane, subject to an applied electric field. With an impact of a nonlocal term, Duan and Vincent [\[1\]](#page--1-0) studied the dynamics concerning the deterministic nonlocal K–S equation. In a successive paper [\[2\]](#page--1-1), the authors discussed a stochastic version of Eq. [\(1\)](#page-0-3) with an additive white noise. They proved that a unique weak solution exists in $L^4(0,T;L^4(G))$, P-a.s. for the equation with homogeneous Dirichlet boundary conditions. In [\[3\]](#page--1-2), Yang discussed an analogous subject as in [\[1\]](#page--1-0) for the equation driven by an additive white noise with the impact of the nonlocal term, which is described in the following form:

$$
\begin{cases} du(t) + (u_{xxxx}(t) + u_{xx}(t) + u(t)u_x(t)) dt + \alpha \mathcal{H} \mathcal{I} (u_{xxx}(t)) dt = \sigma dW(t), & \text{in } G, \\ u \text{ is periodic on } G := (-\ell, \ell), \quad \text{i.e., } u(x + \ell) = u(x - \ell) \text{ for } x \in G, \end{cases}
$$
\n
$$
(2)
$$

where the constants $\ell > 0$, $\alpha > 0$, $\sigma \in \mathbb{R}$ and the nonlocal term $\mathcal{H}(\mu)$ is the Hilbert transform given by

$$
\mathcal{H} \mathcal{I}(u)(x) := -\frac{1}{2\ell} \int_{-\ell}^{\ell} \cot \frac{\pi (x - y)}{2\ell} u(y) dy, \quad x \in G.
$$
 (3)

The noise term in [\(2\)](#page-0-4) is an additive noise σdW , where $W = (W(t); 0 \le t \le 1)$ is a Q-Wiener process on a complete probability space $(\Omega, \mathcal{F}, (\mathcal{F}_t; t \geq 0), \mathbb{P})$, where the filtration $(\mathcal{F}_t; t \geq 0)$ satisfies the usual conditions.

a b s t r a c t

In this paper, we establish a large deviation principle for the (weak) solution to a nonlocal Kuramoto–Sivashinsky stochastic partial differential equation with small noise perturbation. The key technique is an application of the contraction principle. © 2013 Elsevier Ltd. All rights reserved.

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In this paper, we are concerned with a large deviation principle (LDP) for the (weak) solution $u^\varepsilon = (u^\varepsilon(t); 0 \le t \le 1)$ to Eq. [\(2\)](#page-0-4) with small noise perturbations. Namely, for any $\varepsilon > 0$ and $t \in [0, 1]$, the $(\mathcal{F}_t; t \ge 0)$ -adapted process u^ε is governed by

$$
\begin{cases} du^{\varepsilon}(t) + \left(u^{\varepsilon}_{xxxx}(t) + u^{\varepsilon}_{xx}(t) + u^{\varepsilon}(t)u^{\varepsilon}_{x}(t) \right)dt + \alpha \mathcal{H} \mathcal{I} \left(u^{\varepsilon}_{xxx}(t) \right)dt = \varepsilon \sigma dW(t), & \text{in } G; \\ u^{\varepsilon} \text{ is periodic on } G. \end{cases}
$$
 (4)

The LDP for various stochastic partial differential equations (SPDEs) driven by white noise have been studied in the literature (see, e.g., [\[4–6\]](#page--1-3) and references therein). In [\[4\]](#page--1-3), Cardon-Weber obtained the LDP for the 1-dimensional stochastic Burgers type equation by proving the uniform Freidlin–Wentzell estimates. Further, Carreras and Sarrà studied the LDP for a *d*-dimensional stochastic heat equation with spatially correlated noise in [\[5\]](#page--1-4). Recently, Röckner et al. [\[6\]](#page--1-5) established the LDP for stochastic generalized porous media equations using the generalized contraction principle (see Theorem 3.2 therein). Motivated by the idea employed in [\[6\]](#page--1-5), we prove the LDP of the (weak) solution to Eq. [\(2\)](#page-0-4) with small noise perturbations (namely, the solution to Eq. [\(4\)](#page-1-0) for $\varepsilon > 0$) by adopting a version of the contraction principle (see Theorem 4.2.23 in [\[7\]](#page--1-6)). Note that this technique has been applied to derive the corresponding LDP for diffusion processes or delay SDEs (see, e.g. [\[7,](#page--1-6)[8\]](#page--1-7)).

The rest of this paper is organized as follows. In Section [2,](#page-1-1) some preliminaries are given. Section [3](#page--1-8) is devoted to establishing probability properties of the (weak) solution to Eq. [\(4\).](#page-1-0) In Section [4,](#page--1-8) we explore the skeleton equation corresponding to Eq. [\(4\).](#page-1-0) Finally a LDP of the (weak) solution to the nonlocal K–S SPDE with small noise perturbation is established in Section [5.](#page--1-9)

2. Preliminaries

This section would introduce some basic notation, function spaces and functional inequalities used frequently in the paper.

First, we recall a basic fact on the solution to Eq. [\(2\)](#page-0-4) with the initial value u_0 when the diffusive coefficient $\sigma = 0$, namely, the spatial average \bar{u}_0 of *u*,

$$
\bar{u}_0 := \frac{1}{2\ell} \int_{-\ell}^{\ell} u(t,x) dx = \frac{1}{2\ell} \int_{-\ell}^{\ell} u_0(x) dx, \quad \forall t \geq 0.
$$

W.L.G., suppose that $\bar{u}_0 = 0$ throughout the paper. Thus, we can define the following function spaces,

$$
\begin{cases}\nH := \left\{ u \in L^2(G); \ u \text{ is periodic on } G, \int_{-\ell}^{\ell} u(x) dx = 0 \right\}, \\
H_{\text{per}}^p := \left\{ u \in W^{2,p}(G); \ u \text{ is periodic on } G \right\}, \quad \text{for } p \in \mathbb{N}, \\
\dot{H}_{\text{per}}^p := H_{\text{per}}^p \cap H, \quad \text{for } p \in \mathbb{N}.\n\end{cases}
$$

For $i \in \mathbb{N} \cup \{0\}$, let $D_i \coloneqq \frac{\partial^i}{\partial x^i}$ $\frac{\partial^i}{\partial x^i}$ and *D*₀ = *I* (the identity operator on *H*). Then *A* = −*D*₂ is a positive self-adjoint unbounded linear operator with domain $D(A)$. Let $(\lambda_k)_{k\in\mathbb{N}}$ and $(e_k)_{k\in\mathbb{N}} := (\phi_k(x), \psi_k(x))_{k\in\mathbb{N}}$ be the eigenvalues and corresponding eigenfunctions of $A: D(A) \rightarrow H$. Then, it holds that

$$
\begin{cases}\n\lambda_k = \frac{\pi^2 k^2}{\ell^2}, \\
\phi_k(x) = \frac{1}{\sqrt{\ell}} \sin\left(\frac{k\pi x}{\ell}\right), \\
\psi_k(x) = \frac{1}{\sqrt{\ell}} \cos\left(\frac{k\pi x}{\ell}\right),\n\end{cases}
$$

and $(e_k)_{k \in \mathbb{N}}$ forms a complete orthonormal basis of *H*. By the properties of the operator *A*, for $s \in \mathbb{R}$, the spectral theory allows us to define the powers *A ^s* of *A* by (see [\[9\]](#page--1-10))

$$
\begin{cases}\nD(A^s) = \left\{ u \in H; \sum_{k=1}^{\infty} \lambda_k^{2s} (u, e_k)^2 < \infty \right\}, \\
A^s u = \sum_{k=1}^{\infty} \lambda_k^s (u, e_k) e_k, \quad \text{for } u \in D(A^s),\n\end{cases}
$$

where (\cdot, \cdot) and $|\cdot|$ denote the inner product and the corresponding norm of *H*. We endow the domain $D(A^s)$ of $A^s : D(A^s) \to$ *H* with the following inner product and the norm

$$
\begin{cases}\n(u, v)_{2s} = (A^s u, A^s v), \\
|u|_{2s} = |A^s u|,\n\end{cases}
$$

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