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## **Nonlinear Analysis**





# Blow-up of rough solutions to the fourth-order nonlinear Schrödinger equation\*

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#### ABSTRACT

This paper deals with the formation of singularities of rough blow-up solutions to the fourth-order nonlinear Schrödinger equation. The limiting profile and  $L^2$ -concentration of the rough blow-up solutions are obtained in  $H^s(\mathbb{R}^4)$  with  $s>s_0$ , where  $s_0 \leq \frac{9+\sqrt{721}}{20} \approx 1.793$ . The new ingredient relies on the refined compactness result developed by Zhu et al. [S.H. Zhu, J. Zhang, H. Yang, Limiting profile of the blow-up solutions for the fourth-order nonlinear Schrödinger equation, Dyn. Partial Differ. Equ. 7 (2010) 187–205].

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#### 1. Introduction

In this paper, we study the Cauchy problem of the following fourth-order nonlinear Schrödinger equation

$$iu_t - \Delta^2 u + |u|^2 u = 0, \quad t > 0, \ x \in \mathbb{R}^4,$$
 (1.1)

$$u(0,x) = u_0,$$
 (1.2)

where  $i = \sqrt{-1}$ ;  $\Delta^2 = \Delta \Delta$  is the biharmonic operator defined in  $\mathbb{R}^4$  and  $\Delta = \sum_{j=1}^4 \frac{\partial^2}{\partial x_j^2}$  is the Laplace operator in  $\mathbb{R}^4$ ;

u=u(t,x):  $[0,T^*)\times\mathbb{R}^4\to\mathbb{C}$  is the complex valued function and  $0< T^* \leq +\infty$ . Fourth-order Schrödinger equations are introduced by Karpman [1], Karpman and Shagalov [2] to take into account the role of small fourth-order dispersion terms in the propagation of intense laser beams in a bulk medium with Kerr nonlinearity, and such fourth-order Schrödinger equations are written as

$$i\phi_t + \varepsilon \Delta^2 \phi + \mu \Delta \phi + |\phi|^{p-1} \phi = 0, \qquad \phi = \phi(t, x) : I \times \mathbb{R}^d \to \mathbb{C},$$
 (1.3)

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where  $\varepsilon$  and  $\mu$  are two parameters, d is the space dimension. Note that Eq. (1.1) is a special case of Eq. (1.3) by taking  $\varepsilon = -1$ ,  $\mu = 0$  and  $p = 1 + \frac{8}{d} = 3$ . Eq. (1.1) is called the mass-critical due to the mass  $M(u) = \int_{\mathbb{R}^4} |u(t,x)|^2 dx$  and the equation itself are invariant under the rescaling symmetry  $u \mapsto \lambda^2 u(\lambda^4 t, \lambda x)$ .

We recall some known results for the classical focusing mass-critical nonlinear Schrödinger equation

$$iv_t + \Delta v + |v|^{\frac{4}{d}}v = 0, \quad v(0, x) = \varphi,$$
 (1.4)

where  $v=v(t,x):I\times\mathbb{R}^d\to\mathbb{C}$ . Ginibre and Velo [3] showed the local well-posedness in  $H^1(\mathbb{R}^d)$ . In this space energy arguments apply, and a blow-up theory has been developed in the last two decades (see [4–6] and the references therein). On the other hand, Cazenave [4] established the local well posedness in  $H^s(\mathbb{R}^d)$  with  $0\le s<1$ . In this space, the energy arguments fail and studying the rough blow-up solutions is more difficult and interesting. Recently, many researchers are attracted to study the qualitative properties of blow-up solutions in lower regular space  $H^s(\mathbb{R}^d)$  with  $0\le s<1$  (see [7–12]). Combining the harmonic techniques with variational characteristic of the ground state, Colliander et al. [8] firstly obtained the mass concentration properties of the radially symmetric blow-up solutions in  $H^s(\mathbb{R}^2)$  with  $1>s>\frac{1+\sqrt{11}}{5}$ . Hmidi and Keraani [10] extended this result to the general blow-up solutions. Tzirakis [13] obtained the analogue results in  $H^s(\mathbb{R}^1)$  with  $1>s>\frac{10}{11}$ . Visan and Zhang [12] extended Colliander et al.'s results [8] to the general blow-up solutions in dimension  $d\ge 3$  for some  $1>s>s_0$ . For  $\varphi\in L^2(\mathbb{R}^2)$ , Bourgain [14] showed that some small amount of mass must concentrate in parabolic windows (at least along a subsequence). Keraani [15], Bégout and Vargas [16], Chae et al. [17] extended Bourgain's results [14] to d=1 and  $d\ge 3$ .

In Eq. (1.1), if one replaces the nonlinearity  $|u|^2u$  with  $|u|^{p-1}u$ , it is a class of semilinear fourth-order Schrödinger equations similar to Eq. (1.1), which has been widely investigated. For  $1 (we use the convention: <math>\frac{2d}{(d-4)^+} = +\infty$  when  $d \le 4$  and  $\frac{2d}{(d-4)^+} = \frac{2d}{d-4}$  when d > 4), Ben-Artzi et al. [18] established the local well-posedness in  $H^2(\mathbb{R}^d)$ . Fibich et al. [19] obtained the general results of global well-posedness in  $H^2(\mathbb{R}^d)$ . Pausader [20] and Segata [21] studied the global well-posedness and scattering of the fourth-order nonlinear Schrödinger equation with cubic nonlinearity. For  $p = \frac{2d}{d-4}$ , Miao et al. [22], Pausader [23] studied the global existence and scattering of the focusing fourth-order nonlinear Schrödinger equation; Miao et al. [24], Pausader [25] studied the global existence and scattering of the defocusing fourth-order nonlinear Schrödinger equation. The above studies focused on global solutions.

In the present paper, we study the limiting profile and  $L^2$ -concentration of blow-up solutions to the Cauchy problem (1.1)–(1.2) in  $H^s := H^s(\mathbb{R}^4)$  with 0 < s < 2. We first recall some known results for the Cauchy problem (1.1)–(1.2) in  $H^2 := H^2(\mathbb{R}^4)$ . Kenig et al. [26], Ben-Artzi et al. [18], Pausader [25] established the local well-posedness in  $H^2$ . Moreover, for the solution u(t, x) of the Cauchy problem (1.1)–(1.2), there are two conservation laws in  $H^2$ :

#### (i) Conservation of mass

$$\int_{\mathbb{R}^4} |u(t,x)|^2 dx = \int_{\mathbb{R}^4} |u_0|^2 dx.$$
 (1.5)

#### (ii) Conservation of energy

$$E(u(t)) := \frac{1}{2} \int_{\mathbb{R}^4} |\Delta u(t, x)|^2 dx - \frac{1}{4} \int_{\mathbb{R}^4} |u(t, x)|^4 dx = E(u_0).$$
 (1.6)

On the other hand, we mean a special periodic solution of Eq. (1.1) in the form  $u(t, x) = Q(x)e^{-it}$  by the standing wave. It is easy to check that Q(x) satisfies

$$\Delta^2 Q + Q - |Q|^2 Q = 0, \quad Q \in H^2, \tag{1.7}$$

and Q is called the ground state of Eq. (1.7) (see [19]). Zhu et al. [27] showed the existence of the ground state of Eq. (1.7). Levandosky [28] discussed the stability of the standing waves. Fibich et al. [19] showed some numerical observations of the solution to the Cauchy problem (1.1)–(1.2), which implies that if the initial data  $\|u_0\|_{L^2} < \|Q\|_{L^2}$ , then the solution u(t,x) exists globally; if the initial data  $\|u_0\|_{L^2} \ge \|Q\|_{L^2}$ , then the solution u(t,x) may blow up in finite time. Because of the effect of fourth-order dispersion  $\Delta^2 u$ , whether the variance identity arguments can be extended to show the existence of blow-up solutions for the fourth-order nonlinear Schrödinger equation is still unknown (see [29,19]). However, the numerical observations in [19] showed the existence of blow-up solutions. Baruch et al. [29], Zhu et al. [27] showed some dynamical properties of the blow-up solutions in the energy space  $H^2$ , such as blow-up rate,  $L^2$ -concentration and limiting profile. Meanwhile, Zhu et al. [27] obtained a refined compactness result, which is of particular significance to study the dynamical properties of blow-up solutions to fourth-order nonlinear Schrödinger equations.

**Theorem 1.1.** Let d=4 and  $\{v_n\}_{n=1}^{+\infty}$  be a bounded sequence in  $H^2$  such that

$$\lim_{n \to +\infty} \sup \|\Delta v_n\|_{L^2} \le M \quad and \quad \lim_{n \to +\infty} \sup \|v_n\|_{L^4} \ge m. \tag{1.8}$$

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