



# Asymptotically self-similar solutions of the damped wave equation

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## ABSTRACT

We consider the damped hyperbolic equation

$$\varepsilon u_{\tau\tau} + u_{\tau} = (a(\xi)u_{\xi})_{\xi} - |u|^{p-1}u, \quad (\xi, \tau) \in \mathbb{R} \times \mathbb{R}_+, \quad (1)$$

where  $\varepsilon > 0$ ,  $a(\xi) \rightarrow 1$  as  $|\xi| \rightarrow +\infty$  and  $1 < p < 3$ .

We prove in this article that the exact self-similar solutions of the semi-linear parabolic equation obtained by setting  $\varepsilon = 0$  and  $a(\xi) \equiv 1$  in (1) are also asymptotically stable self-similar solutions of the Eq. (1). The proof of our result relies on various energy estimates rewritten in the variables  $\xi/\sqrt{\tau}$ ,  $\ln \tau$ .

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## 1. Introduction

In this paper, we study the asymptotic stability of the solution of the damped hyperbolic equation given as follows:

$$\varepsilon u_{\tau\tau} + u_{\tau} = (a(\xi)u_{\xi})_{\xi} - |u|^{p-1}u, \quad (\xi, \tau) \in \mathbb{R} \times \mathbb{R}_+, \quad (1.1)$$

where  $\varepsilon$  is a positive, (not necessarily small) parameter, and  $p \in ]1, 3[$ . We assume that the diffusion coefficient  $a(\xi)$  is positive and satisfies  $\lim_{\xi \rightarrow \pm\infty} a(\xi) = 1$ .

Equations of type (1.1) are commonly used as mathematical models in several fields dealing with spreading and interacting particles: especially in genetics and population dynamics; see for example [1,2]. In the case of an inhomogeneous medium, the diffusion coefficients in such equations may depend on the space variable  $\xi$ .

In the case  $p > 3$ , Gally and Raugel [3] have studied the stability of small solutions of Eq. (1.1) under the hypothesis  $\lim_{\xi \rightarrow \pm\infty} a(\xi) = a_{\pm} > 0$ . In fact, they introduced scaling variables and used energy estimates to show that the asymptotic behavior of a small solution  $u(\xi, \tau)$  of Eq. (1.1) is entirely determined, up to the second order, by a linear parabolic equation depending only on  $a_{\pm}$ .

In the case  $\varepsilon = 0$  and  $a(\xi) \equiv 1$ , Eq. (1.1) reduces to the following semi-linear parabolic equation:

$$u_{\tau} = u_{\xi\xi} - |u|^{p-1}u. \quad (1.2)$$

The asymptotic behavior of Eq. (1.2) has been rather extensively studied. We here recall some known results. First, we remark that if  $u(\xi, \tau)$  is a solution of Eq. (1.2), then for all  $\lambda > 0$ ,  $u_{\lambda}(\xi, \tau) = \lambda^{\frac{2}{p-1}} u(\lambda^2 \tau, \lambda \xi)$  is also a solution. A solution  $u \neq 0$  is said to be self-similar, when  $u_{\lambda} \equiv u$ , for all  $\lambda > 0$ .

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In the  $N$  multidimensional case, with  $1 < p < 1 + \frac{2}{N}$ , we already know that if Eq. (1.2) has a self-similar solution with initial data  $u_0$ , then  $u_0$  is homogenous of degree  $\frac{2}{1-p}$ . On the other hand, Cazenave, Dickstein, Escobedo and Weissler [4] have shown, among many other results, that there exists a unique self-similar solution  $\tilde{u}$  of Eq. (1.2) such that

$$\lim_{|\xi| \rightarrow +\infty} |\xi|^{\frac{2}{p-1}} (\tilde{u}(\xi, 1) - u_0(\xi)) = 0. \quad (1.3)$$

These self-similar solutions are related to the asymptotic behavior of  $u$  solution of Eq. (1.2) in the following sense: suppose that there exist  $\omega$ , a homogenous function of degree 0, and  $u_0 \in \mathcal{C}(\mathbb{R})$  an asymptotically homogenous function in space in the following sense:

$$\lim_{|\xi| \rightarrow +\infty} |\xi|^{\frac{2}{p-1}} u_0(\xi) - \omega\left(\frac{\xi}{|\xi|}\right) = 0. \quad (1.4)$$

Assume in addition that  $u_0 \geq 0$ ,  $u_0 \not\equiv 0$ . Then, the solution  $u$  of Eq. (1.2) is asymptotically self-similar in the sense that

$$\lim_{\tau \rightarrow +\infty} \sup_{\xi \in \mathbb{R}} \left( (\tau + |\xi|^2)^{\frac{1}{p-1}} (u(\xi, \tau) - \tilde{u}(\xi, \tau)) \right) = 0. \quad (1.5)$$

Recently, in [5], the authors proved that if the initial data does not behave asymptotically like the profile, as in (1.4), in a certain precise way, then the solution is not asymptotically self-similar. In the same way, an instability result has been established in [6,7].

Returning now to the one-dimensional case, a self-similar solution of Eq. (1.2) has the form  $u(\xi, \tau) = \tau^{-\frac{1}{p-1}} f\left(\frac{\xi}{\sqrt{\tau}}\right)$ , where  $f$  satisfies an ordinary differential equation

$$f''(x) + \frac{x}{2} f'(x) + \frac{1}{p-1} f(x) - |f(x)|^{p-1} f(x) = 0, \quad \text{in } \mathbb{R}_+. \quad (1.6)$$

Here we restrict our study to positive self-similar solutions satisfying  $f'(0) = 0$ . As shown in [8–10], for  $p \in ]1, 3[$  and  $\gamma \in [\gamma_p, \gamma^*]$ , where  $0 < \gamma_p < \gamma^* = (p-1)^{-\frac{1}{p-1}}$ , there exists a smooth, even, positive solution  $f_\gamma$  of Eq. (1.6) such that  $f_\gamma(0) = \gamma$ . Moreover, the function

$$c : [\gamma_p, \gamma^*] \longrightarrow [0, +\infty[ \\ \gamma \longmapsto \lim_{x \rightarrow +\infty} x^{\frac{2}{p-1}} f_\gamma(x), \quad (1.7)$$

is a bijection. On the other hand, the decay at infinity of these solutions  $f_\gamma$  is as follows (see [8]):

$$\begin{aligned} \text{for } \gamma > \gamma_p, \quad f_\gamma(x) &= c(\gamma) |x|^{\frac{-2}{p-1}} + \mathcal{O}\left(|x|^{\frac{-2}{p-1}-2}\right), & \text{as } |x| \rightarrow +\infty, \\ \text{for } \gamma = \gamma_p, \quad f_\gamma(x) &= c(\gamma_p) |x|^{\frac{2}{p-1}-1} e^{-\frac{x^2}{4}} + \mathcal{O}\left(|x|^{\frac{2}{p-1}-3} e^{-\frac{x^2}{4}}\right), & \text{as } |x| \rightarrow +\infty, \end{aligned} \quad (1.8)$$

where  $c(\gamma)$  is defined by (1.7) and  $c(\gamma_p)$  is a positive constant. Moreover,

$$\text{for } \gamma > \gamma_p, \quad \frac{f'_\gamma(x)}{f_\gamma(x)} = \frac{-2}{p-1} |x|^{-1} + \mathcal{O}(|x|^{-3}), \quad \text{as } |x| \rightarrow +\infty. \quad (1.9)$$

The stability of these self-similar solutions has been widely studied (see for example [11,12,10,13,14]).

For instance, using the properties (1.8), Escobedo, Kavian and Matano [9] have shown that, for  $\gamma \in [\gamma_p, \gamma^*]$ , for initial data  $u(\xi, 1)$ , at  $\tau = 1$  near  $f_\gamma(\xi)$  in the space  $L^\infty\left(\mathbb{R}, \left(1 + |x|^{\frac{2}{p-1}}\right)\right)$  and  $\lim_{|\xi| \rightarrow +\infty} |\xi|^{\frac{2}{p-1}} u(\xi, 1) = \lim_{|\xi| \rightarrow +\infty} |\xi|^{\frac{2}{p-1}} f_\gamma(\xi)$ , there exists a unique global classical solution  $u(\xi, \tau)$  of (1.2)  $u(\xi, \tau)$  in the space  $L^\infty\left([0, +\infty[, L^\infty\left(\mathbb{R}, \left(1 + |x|^{\frac{2}{p-1}}\right)\right)\right)$ , which moreover satisfies

$$\lim_{\tau \rightarrow +\infty} \left\| \tau^{\frac{1}{p-1}} u(\xi \sqrt{\tau}, \tau) - f_\gamma(\xi) \right\|_{L^\infty} = 0. \quad (1.10)$$

In a more recent paper, Brimont and Kupiainen [15] have shown, that if  $\gamma \in [\gamma_p, \gamma^*]$  and if the initial data  $u(\xi, 1)$  are near  $f_\gamma(\xi)$  in the weighted space  $L^\infty(\mathbb{R}, (1 + |x|^q))$ , with  $q > \frac{2}{p-1}$ , Eq. (1.2) has a unique global classical solution  $u(\xi, \tau)$  in the space  $L^\infty([0, +\infty[, L^\infty(\mathbb{R}, (1 + |x|^q)))$ . Moreover, there exist  $\mu > 0$  and  $C > 0$  such that, for all  $\tau \geq 1$ ,

$$\left\| \left[ \tau^{\frac{1}{p-1}} u(\xi \sqrt{\tau}, \tau) - f_\gamma(\xi) \right] (1 + |\xi|^q) \right\|_{L^\infty} \leq C \tau^{-\mu} \|u(\xi, 1) - f_\gamma(\xi)\|_{L^\infty}. \quad (1.11)$$

In the whole paper, we fix a positive number  $\gamma$  in  $[\gamma_p, \gamma^*]$ , and we denote by  $g_\gamma$  the self-similar positive solution  $g_\gamma(\xi, \tau) = \tau^{-\frac{1}{p-1}} f_\gamma\left(\xi \tau^{-\frac{1}{2}}\right)$  of Eq. (1.2).

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