



Research on the spin-Hamiltonian parameters for the approximately trigonal Mn^{5+} tetrahedral center in $\text{Sr}_5(\text{PO}_4)_3\text{Cl}$ crystal



Wu Xiao-Xuan^a, Zheng Wen-Chen^{b,*}

^a Department of Physics, Civil Aviation Flight University of China, Guanghan, People's Republic of China

^b Department of Material Science, Sichuan University, Chengdu, People's Republic of China

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ABSTRACT

The spin-Hamiltonian parameters (g factors g_{\parallel} , g_{\perp} and zero-field splitting D) of the approximately trigonal Mn^{5+} tetrahedral cluster in $\text{Sr}_5(\text{PO}_4)_3\text{Cl}$ crystal are computed from the complete high-order perturbation formulas established on the two-mechanism model, in which in addition to the contributions due to the frequently-applied crystal-field (CF) mechanism in the CF theory, those due to the commonly-ignored charge-transfer (CT) mechanism are contained. The calculated results are in rational accordance with the experimented values. The calculations demonstrate that for the high valence state d^n clusters in crystals, the contributions to spin-Hamiltonian parameters from the CT mechanism should also be considered. The negative sign of zero-field splitting D and the defect structure of the approximately trigonal Mn^{5+} center in $\text{Sr}_5(\text{PO}_4)_3\text{Cl}$ are proposed from the calculations.

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1. Introduction

Since the infrared laser action was discovered for Cr^{4+} ion at the tetrahedral Si^{4+} site of forsterite Mg_2SiO_4 crystal [1,2], the tetrahedrally coordinated $3d^2$ (e.g., Cr^{4+} , Mn^{5+} and Fe^{6+}) ions doped in crystals and their spectroscopic properties have attracted significant attention because these materials have the potential applications in infrared tunable lasers [3–10] and the vivo fluorescence imaging [11]. Misra et al. [10] reported the EPR spectra of Mn^{5+} -doped $\text{Sr}_5(\text{PO}_4)_3\text{Cl}$ crystal. From the EPR experiment, an orthorhombic Mn^{5+} tetrahedral center attributed to Mn^{5+} occupying the tetrahedral P^{5+} site was found and its spin-Hamiltonian parameters (g factors g_{\parallel} , g_{\perp} and zero-field splittings D , E) were given [10]. Since the orthorhombic component E of zero-field splitting is very small and may be neglected, the Mn^{5+} center can be approximated as C_{3v} symmetry [10]. To date no theoretical explanations or calculations for the spin-Hamiltonian parameters g_{\parallel} , g_{\perp} and D of the approximately trigonal Mn^{5+} center in $\text{Sr}_5(\text{PO}_4)_3\text{Cl}$ crystal have been found in the literature. It is known that the reasonable calculations rely upon the suitable theoretical method or model. Mn^{5+} ion has high valence state which leads the charge-transfer (CT) energy levels to become lower, so, besides the contributions to spin-Hamiltonian parameters due to crystal-field (CF) mechanism related to the interactions of CF excited states with the ground state in the vastly-applied CF theory, the contributions due to CT mechanism concerning the interactions of CT excited states with the ground state are also considered. Thus, a two (CF and CT)-mechanism model [12,14] should be applied here. In the present paper, the high-order perturbation formulas established on the two-mechanism model are adopted to calculate the spin-Hamiltonian parameters g_{\parallel} , g_{\perp} and D for the approximately trigonal Mn^{5+} tetrahedral center in $\text{Sr}_5(\text{PO}_4)_3\text{Cl}$ crystal. The results are discussed.

2. Calculations

By means of the one-electron basis functions (which consist of anti-bonding and bonding orbitals, corresponding to CF and CT mechanisms, respectively) and the perturbation method in the two-mechanism model [14], the high-order perturbation formulas of spin-Hamiltonian parameters g_{\parallel} , g_{\perp} and D for a trigonal d^2 tetrahedral cluster are deduced as [14]

* Corresponding author. Tel.: +86 2885412371; fax: +86 28 85416050.

E-mail address: zhengwc1@163.com (Z. Wen-Chen).

$$\begin{aligned}
g_{//} &= g_s + \Delta g_{//}^{CF} + \Delta g_{//}^{CT} \\
\Delta g_{//}^{CF} &= -4k'_{CF}\zeta'_{CF}/E_1 - [g_e(\zeta'_{CF})^2 + k'_{CF}\zeta_{CF}\zeta'_{CF} - k_{CF}(\zeta'_{CF})^2/2]/E_1^2 \\
&\quad - (g_e - k_{CF}/2)(\zeta'_{CF})^2/E_2^2 - k'_{CF}\zeta_{CF}\zeta'_{CF}/(E_1E_2) - 6Bk'_{CF}(\zeta'_{CF})^2 \\
&\quad \times [2/(E_1E_2E_5) + 1/(E_2^2E_5)] + 4k'_{CF}\zeta'_{CF}\nu/(3E_1^2) - 4\sqrt{2}\zeta'\nu' \\
&\quad \times [k_{CF}/(E_1E_3) + 12Bk'_{CF}/(E_1E_3E_4)] \\
\Delta g_{//}^{CT} &= 4(k'_{CT}\zeta'_{CT}/E_n + k_{CT}\zeta_{CT}/E_a) - 8k_{CT}\zeta_{CT}\nu/(9E_a^2) \\
g_{\perp} &= g_s + \Delta g_{\perp}^{CF} + \Delta g_{\perp}^{CT} \\
\Delta g_{\perp}^{CF} &= \Delta g_{//}^{CF} - 2k'_{CF}\zeta'_{CF}\nu/E_1^2 + 6\sqrt{2}\zeta'_{CF}\nu'[k_{CF}/(E_1E_3) + 12Bk'_{CF}/(E_1E_3E_4)] \\
\Delta g_{\perp}^{CT} &= \Delta g_{//}^{CT} + 4k_{CT}\zeta_{CT}\nu/(3E_a^2) \\
D &= D^{CF} + D^{CT} \\
D^{CF} &= \frac{\nu}{2}(\zeta'_{CF})^2(1/E_1^2 - 1/E_2^2) + \frac{3\nu'}{\sqrt{2}}\{\zeta_{CF}\zeta'_{CF}[1/(E_2E_3) - 1/(E_1E_3)] + \\
&\quad 4B(\zeta'_{CF})^2[1/(E_2E_3E_5) + 1/(E_2^2E_5) - 3/(E_1E_3E_4) - 3/(E_2E_3E_4)]\} \\
D^{CT} &= -\zeta_{CT}^2\nu/(3E_a^2)
\end{aligned} \tag{1}$$

with

$$\begin{aligned}
\zeta_{CF} &= (N_t^a)^2 \left\{ \zeta_d^0 + [\sqrt{2}\lambda_{\pi}^a\lambda_{\sigma}^a - (\lambda_{\pi}^a)^2/2] \zeta_p^0 \right\}, \\
\zeta'_{CF} &= N_t^a \cdot N_e^a \left\{ \zeta_d^0 + [\lambda_{\pi}^a\lambda_{\sigma}^a/\sqrt{2} + (\lambda_{\pi}^a)^2/2] \zeta_p^0 \right\}, \\
\zeta_{CT} &= N_t^a \cdot N_t^b \left\{ \zeta_d^0 + \left[\frac{\lambda_{\pi}^a\lambda_{\sigma}^b + \lambda_{\pi}^b\lambda_{\sigma}^a}{\sqrt{2}} - \frac{\lambda_{\pi}^a\lambda_{\pi}^b}{2} \right] \zeta_p^0 \right\}, \\
\zeta'_{CT} &= N_t^b \cdot N_e^a \left\{ \zeta_d^0 + \left[\frac{\lambda_{\pi}^a\lambda_{\sigma}^b}{\sqrt{2}} + \frac{\lambda_{\pi}^a\lambda_{\pi}^b}{2} \right] \zeta_p^0 \right\}, \\
k_{CF} &= (N_t^a)^2 \left[1 - (\lambda_{\pi}^a)^2/2 + \sqrt{2}\lambda_{\pi}^a\lambda_{\sigma}^a + 2\lambda_{\sigma}^a S_{dp}(\sigma) + 2\lambda_{\pi}^a S_{dp}(\pi) \right], \\
k'_{CF} &= N_t^a \cdot N_e^a \left[1 - (\lambda_{\pi}^a)^2/2 + \lambda_{\pi}^a\lambda_{\sigma}^a/\sqrt{2} + 4\lambda_{\pi}^a S_{dp}(\pi) + \lambda_{\sigma}^a S_{dp}(\sigma) \right], \\
k_{CT} &= N_t^a N_t^b \left\{ 1 + \left[\frac{\lambda_{\pi}^a\lambda_{\sigma}^b + \lambda_{\pi}^b\lambda_{\sigma}^a}{\sqrt{2}} - \frac{\lambda_{\pi}^a\lambda_{\pi}^b}{2} \right] + (\lambda_{\sigma}^a + \lambda_{\sigma}^b) S_{dp}(\sigma) + (\lambda_{\pi}^a + \lambda_{\pi}^b) S_{dp}(\pi) \right\}, \\
k'_{CT} &= N_t^b N_e^a \left\{ 1 + \left[\frac{\lambda_{\pi}^a\lambda_{\sigma}^b}{\sqrt{2}} - \frac{\lambda_{\pi}^a\lambda_{\pi}^b}{2} \right] + \lambda_{\sigma}^b S_{dp}(\sigma) + (3\lambda_{\pi}^a + \lambda_{\pi}^b) S_{dp}(\pi) \right\},
\end{aligned} \tag{2}$$

where the superscripts and subscripts *CF* and *CT* stand for the parameters in *CF* and *CT* mechanisms. $g_s \approx 2.0023$, the g value of free electron. E_j ($j = 1-5$) are the zero-order energy denominations between the *CF* excited states and the ground state. They are the functions of Racah parameters B , C and the cubic field parameter Dq of the studied system and their expressions are given in Ref. [14]. E_a and E_n are the *CT* energy levels. For $(\text{MnO}_4)^{3-}$ clusters considered, $E_n \approx 31,600 \text{ cm}^{-1}$ and $E_a \approx 33,600 \text{ cm}^{-1}$ [15]. ν and ν' are the trigonal field parameters. ζ_d^0 and ζ_p^0 are the spin-orbit coupling parameters of free d^n ion and free ligand ion. For $(\text{MnO}_4)^{3-}$ cluster, we have $\zeta_d^0 (\text{Mn}^{5+}) \approx 503 \text{ cm}^{-1}$ [16] and $\zeta_p^0 (\text{O}^{2-}) \approx 150 \text{ cm}^{-1}$ [17]. The group overlap integrals $S_{dp}(\beta)$ ($\beta = \sigma, \pi$) rely on the average metal-ligand distance \bar{R} of the studied cluster. The metal-ligand distance R in an impurity center may be unlike the corresponding distance R_h in the host crystal because of the difference between the ionic radii r_i of impurity and r_h of the replaced host ion. An approximate formula $R \approx R_h + (r_i - r_h)/2$ [18] is often applied to estimate the distance R . From $R_{h1} \approx 1.540 \text{ \AA}$, $R_{h2} \approx 1.539 \text{ \AA}$ [19], $r_i (\text{Mn}^{5+}) \approx 0.47 \text{ \AA}$ and $r_h (\text{P}^{5+}) \approx 0.31 \text{ \AA}$ [20], we have $R_1 \approx 1.62 \text{ \AA}$, $R_2 \approx 1.619 \text{ \AA}$ and $\bar{R} \approx 1.639 \text{ \AA}$. Thus, we obtain $S_{dp}(\sigma) \approx -0.1151$ and $S_{dp}(\pi) \approx 0.0393$ from the Slater-type self-consistent field (SCF) functions [21,22] for the $(\text{MnO}_4)^{3-}$ cluster in $\text{Sr}_5(\text{PO}_4)_3\text{Cl}:\text{Mn}^{5+}$ crystal.

The molecular orbital (MO) coefficients N_{γ}^X ($X = a, b, \gamma = t, e$) and λ_{β}^X needed for calculations of the parameters in Eq. (2) follow the normalization correlations [14]

$$\begin{aligned}
N_e^X &= \left[1 + 3(\lambda_{\sigma}^X)^2 + 6\lambda_{\pi}^X S_{dp}(\pi) \right]^{-1/2} \\
N_t^X &= \left[1 + (\lambda_{\sigma}^X)^2 + (\lambda_{\pi}^X)^2 + 2\lambda_{\sigma}^X S_{dp}(\sigma) + 2\lambda_{\pi}^X S_{dp}(\pi) \right]^{-1/2},
\end{aligned} \tag{3}$$

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