Optik 126 (2015) 1209–1212

Contents lists available at ScienceDirect

Optik

journal homepage: www.elsevier.de/ijleo

Study on impact of spatial filter on a hot image through medium with gain

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ARTICLE INFO

Article history:
Received 18 February 2014
Accepted 3 March 2015

Keywords:
Hot image
Spatial filter
Self-focusing
Gain

ABSTRACT

The evolution of hot image formation through medium with gain in consideration of the effect of spatial filter is theoretically and numerically investigated. Based on the linear diffraction theory and small-scale self-focusing theory of Bespalov and Talanov, intensity distribution of hot image in conjugate plane is derived analytically. Then, the peak intensity of hot image for different medium gain and different pinhole sizes are discussed in detail, the results show theoretical analysis is mostly approximate to the numerical simulations, furthermore, it is found that suppressing effect on peak intensity ratio with small gain coefficient is larger than that with bigger ones for determined pinhole size.

1. Introduction

In high power laser systems, hot image resulting from the nonlinear holography is a paramount process that draws considerable attention; it is one of predominant factors that limit maximum output power available from solid-state laser. Once plaguing the safe operation of high power system, damage caused by hot image has not clearly understood for several decades because of its queer characteristic of hard tracks. The formation of hot image originates from the scatter embedded in a strong background beam. After propagating a distance in free space, the scatter wave interferes with the strong background beam and produces interference pattern, as they enter a second order nonlinear medium, an intensity-dependent term is imposed onto the phase front of the incident beam owing to the functionality of the nonlinear medium, then the phase modulated beam proceed to propagate in free space, subsequently a conjugate wave to the initial scatter wave is generated and converges to produce an intensified holographic image of the scatter downstream in corresponding position. So hot image phenomena is also nicknamed nonlinear holographic imaging. The peak intensity of hot image may be several times larger than the initial background beam, What's worse, costly optics may be damaged if intensities of hot image exceed damage threshold of materials even the anticipated average fluences should have been at the safe operation point. The physical mechanism exposing the formation of the hot image [1] is first demonstrated by Hunt et al. afterwards many researches on the character of hot image is constantly pursued. Widmayer et al. [2,3] successively provided the nonlinear formation of images of obscuration and phase errors experimentally, and the computer model is proved to be in good agreement with experimental results. Moreover, he revealed that phase scatter pose a larger damage threat to optical components than the amplitudes ones. Xie et al. [4–6] developed an analytical method for hot image. Peng et al. [7] numerically investigated the evolution of hot image in high power systems with a single thick medium. In addition, some researchers extended studies of the hot image in complicated systems comprising the cascaded nonlinear medium and in special conditions in which the multiple obscurations and arrayed mechanical defects [8–10] are introduced. Peng et al. [11] analyzed the restraining effect of spatial filters on hot image with nonlinear medium of no gain, as we know, in high power systems spatial filters are most commonly used between amplifying chains for filtering out high frequency, so it is required to discuss the impacts of spatial filter on a hot image through medium with gain, in this paper, we present theoretical and numerical treatment for hot image formation considering the spatial filter effect. It may be helpful for designing the high power systems and minimizing the damage risks coming from the hot image.

2. Model and theoretical analysis

The principle suppressing effect of spatial filter on a hot image is sketched in Fig. 1. Briefly, the scatter illuminated by an intense background beam is located in plane A, after propagating a distance d_0, The scatter wave along with the background beam passes through the spatial filter, which is composed of two lens with the

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http://dx.doi.org/10.1016/j.ijleo.2015.03.014

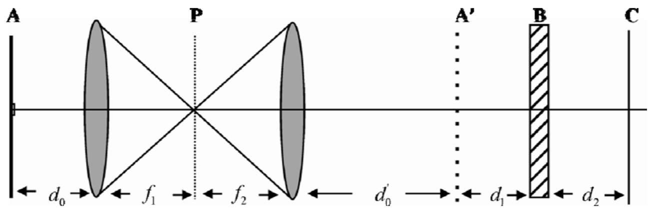

Fig. 1. Sketch of the spatial filter and formation of hot image.

focal length f_1 and f_2, respectively, a pinhole is positioned on the common focus plane of two lens. Then the beams go on propagating until they fall onto a nonlinear medium with gain coefficient g_0, the thickness of the medium is L, as the beams propagate through certain distance in free space, a hot image with large peak intensity is generated.

Assuming $\tau_0(x,y)$ as the transmission function of the scatter and taking $A(x,y)\exp(jkz)$ as optical field of the continuous background beam, the beam modulated by the scatter in plane A is given by

$$E_A(x,y,0)=A_0(x,y,0)[1+\tau(x,y)]=E_{A0}(x,y,0)+E_{A1}(x,y,0) \quad (1)$$

where $\tau(x,y)$ is written as

$$\tau(x,y)=\tau_0(x,y)-1=\begin{cases} a_0\exp(j\phi)-1 & \text{inside the scatter area} \\ 0 & \text{outside the scatter area} \end{cases} \quad (2)$$

where $a_0(0\le a_0\le 1)$ and $\phi(0\le\phi\le 2\pi)$ denote the amplitude and the phase modulation coefficient of the scatter. Supposing

$$E_{A0}(x,y,0)=A_0(x,y,0),\quad E_{A1}(x,y,0)=A_0(x,y,0)\tau(x,y) \quad (3)$$

According to Collins formula theory, the field in the pinhole plane P is obtained as

$$E_p(x_p,y_p)=\frac{\exp\left[jk(x_p^2+y_p^2)/(2f_0)\right]}{j\lambda f_1}F(x_p,y_p) \quad (4)$$

where k is wave number, λ is wavelength of the incident beam, supposed $F_A(f_x,f_y)$ is Fourier transform spectrum function of optical field $E_A(x,y,0)$, f_x, f_y are components of spatial frequency in x and y directions, respectively by substituting f_x with $x_p/(\lambda f_1)$, $F(x_p,y_p)$ in Eq. (4) can be written as

$$F(x_p,y_p)=F_A\left(\frac{x_p}{\lambda f_1},\frac{y_p}{\lambda f_1}\right)\exp\left[-j\pi\lambda d_0\left(\left(\frac{x_p}{\lambda f_1}\right)^2+\left(\frac{y_p}{\lambda f_1}\right)^2\right)\right] \quad (5)$$

The transmission of the pinhole is described as

$$T_p(x_p,y_p)=\begin{cases} 1 & |r_p|=\sqrt{x_p^2+y_p^2}\le a_c/2 \\ 0 & |r_p|=\sqrt{x_p^2+y_p^2}\ge a_c/2 \end{cases} \quad (6)$$

where a_c is diameter of the pinhole, then after passing through the pinhole, optical field is expressed as

$$E_p'(x_p,y_p)=E_p(x_p,y_p)\times T_p(x_p,y_p) \quad (7)$$

Due to basic properties of image relaying of spatial filter incident beams are reimaged subsequently in plane A′, let setting $f=f_1=f_2$, the relation between project distance d_0 and image distance d_0' is deduced as $d_0'=2f-d_0$ by transfer matrix method. Optical field in plane A′ can be given by

$$E_{A'}(x_{A'},y_{A'})=\frac{\exp(jkS)}{(j\lambda f)}\iint F_A\left(\frac{x_p}{\lambda f},\frac{y_p}{\lambda f}\right)T_p(x_p,y_p)\times\exp\left[-\frac{jk}{f}(x_p x_{A'}+y_p y_{A'})\right]dx_p dy_p \quad (8)$$

From Eq. (8), the Fourier transform spectrum of $E_{A'}$ is obtained.

$$F_{A'}(f_x,f_y)=F_A(f_x,f_y)\times T_p(\lambda f f_x,\lambda f f_y) = \left[F_{A0}(f_x,f_y)+F_{A1}(f_x,f_y)\right]T_p(\lambda f f_x,\lambda f f_y) \quad (9)$$

where $F_{A0}(f_x,f_y)$ and $F_{A1}(f_x,f_y)$ are Fourier transform spectrum of optical field $E_{A0}(x,y)$ and $E_{A1}(x,y)$, respectively. It is clearly shown from Eq. (9) that effect of spatial filter on incident beam is cutting off high frequency components with shutting frequency $\sqrt{f_{xc}^2+f_{yc}^2}=a_c/(\lambda f)$.

Starting from plane A′, the beams go through a distance d_1, then they are injected into the nonlinear medium with gain. As described in Ref. [4], the optical field of the scattered wave away from the rear surface of the nonlinear medium d_2 can be derived based on Bespalov and Talanov (BT) theory and linear transfer diffraction formula, namely that is

$$\begin{aligned} E_{C1}(x,y,z) &= \Im^{-1}\left\{F_{C1\_r}+jF_{C1\_i}\right\} \\ &= \exp\left[\frac{\beta-\alpha}{2}L\right]\Im^{-1}\left\{\begin{aligned} & j\left[\left(\frac{k_0 g}{q_\perp^2}+\frac{q_\perp^2}{4k_0 g}\right)\sinh(gL)\right]\times(F_{A1\_r}T_P-jF_{A1\_i}T_P)e^{-j(\theta_2-\theta_1)} \\ & +\left[\cosh(gL)+j\left(\frac{k_0 g}{q_\perp^2}-\frac{q_\perp^2}{4k_0 g}\right)\sinh(gL)\right]\times(F_{A1\_r}T_P+jF_{A1\_i}T_P)e^{-j(\theta_2+\theta_1)} \end{aligned}\right\} \end{aligned} \quad (10)$$

where F_{A1_r} and F_{A1_i} are expressed as the real part and image part of $F_{A1}(f_x,f_y)$ respectively, $\Im^{-1}$ is defined as reverse Fourier transform function, $q_x=2\pi f_x$, $q_y=2\pi f_y$, and $q_\perp=\sqrt{q_x^2+q_y^2}$ is transverse spatial frequency, $g=\sqrt{q_\perp^2(q_c^2-q_\perp^2)/(4k_0^2)}$ is denoted as the growing rate corresponding to $q_\perp$, in which $q_c=\sqrt{4k_0^2\gamma\bar{I}/n_0}$ is critical point of growing spatial frequency, k_0 is wave number in medium with refractive index n_0, $\theta_2=(q_\perp)^2 d_2/(2k)$, $\theta_1=(q_\perp)^2 d_1/(2k)$. From Eq. (10), it is easy to see that the term on the right hand containing $(F_{A1\_r}T_P-jF_{A1\_i}T_P)$ is related to the hot image, and the other term containing $(F_{A1\_r}T_P+jF_{A1\_i}T_P)$ can provide the minimum intensity in some place. Therefore, intensity distribution in hot image plane satisfying the relation $d_2=d_1$ is given by

$$I=I_0\exp\left[\frac{\beta-\alpha}{2}L\right]\times\left|1+\Im^{-1}\left\{j\left[\left(\frac{k_0 g}{q_\perp^2}+\frac{q_\perp^2}{4k_0 g}\right)\sinh(gL)\right]\times(F_{A1\_r}T_P-jF_{A1\_i}T_P)e^{-j(\theta_2-\theta_2)}\right\}\right|^2 \quad (11)$$

3. Simulations and comparison

In this section, the effect of the spatial filter on hot image is numerically simulated by using split-step Fourier method in nonlinear medium, besides, we also give the analytical results according to Eq. (11), the parameters of the nonlinear medium are taken as follows: the thickness of the medium $L=15$, refractive index $n_0=1.528$, the nonlinear index coefficient $\gamma=3\times10^{-16}\,\mathrm{cm^2/W}$; the incident background beam with