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Optimal convergence rate and regularity of nonrelativistic limit for the nonlinear pseudo-relativistic equations

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ABSTRACT

In this paper, we are concerned with the nonrelativistic limit of the following pseudo-relativistic equation with Hartree nonlinearity or power type nonlinearity

$$\left(\sqrt{-\hbar^2 c^2 \Delta + m^2 c^4} - mc^2\right) u + \mu u = \mathcal{N}(u),$$

where c denotes the speed of light. We prove that as $c \rightarrow \infty$ the ground states of this equation converges to the ground state of its nonrelativistic counterpart

$$-\frac{\hbar^2}{2m} \Delta u + \mu u = \mathcal{N}(u)$$

with convergence rate $1/c^2$ in every H^s norms. Moreover, we show that this rate is optimal.

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1. introduction

We consider the pseudo-relativistic equation

$$i\hbar\partial_t\psi = (\sqrt{-\hbar^2c^2\Delta + m^2c^4} - mc^2)\psi - \mathcal{N}(\psi), \quad (1.1)$$

where

$$\psi = \psi(t, x) : \mathbb{R} \times \mathbb{R}^n \rightarrow \mathbb{C}$$

is the wave function, $n \geq 1$ is the space dimension, \hbar is the reduced Planck constant, $c > 0$ denotes the speed of light and $m > 0$ represents the particle mass. The operator $\sqrt{-\hbar^2c^2\Delta + m^2c^4}$ is defined as the Fourier multiplier with symbol $\sqrt{\hbar^2c^2|\xi|^2 + m^2c^4}$. The nonlinear term $\mathcal{N}(\psi)$ is assumed to be either of the power-type

$$\mathcal{N}(\psi) = |\psi|^{p-2}\psi$$

or of the Hartree-type

$$\mathcal{N}(\psi) = (|x|^{-1} * |\psi|^2)\psi.$$

The equation (1.1) is referred as the pseudo-relativistic nonlinear Schrödinger equation (NLS) (or nonlinear Hartree equation (NLH), respectively) when the nonlinearity is of the power-type (or the Hartree-type, respectively). Throughout the paper, we always assume that

$$n \geq 1 \text{ and } 2 < p < \frac{2n}{n-1}$$

(we set $\frac{2n}{n-1} = \infty$ when $n = 1$) for the power-type nonlinearity, which makes $\mathcal{N}(u)$ be $H^{1/2}(\mathbb{R}^n)$ subcritical, while we assume that $n = 3$ when we refer $\mathcal{N}(u)$ to the Hartree nonlinearity.

This model is considered as a relativistic correction to the nonrelativistic counterpart, because by the Taylor series expansion of the symbol

$$\sqrt{\hbar^2c^2|\xi|^2 + m^2c^4} - mc^2 = mc^2 \left(\sqrt{1 + \frac{\hbar^2|\xi|^2}{m^2c^2}} - 1 \right) = \frac{\hbar^2|\xi|^2}{2m} - \frac{\hbar^4|\xi|^4}{8m^3c^2} + \dots, \quad (1.2)$$

the nonrelativistic equation

$$i\hbar\partial_t\psi = -\frac{\hbar^2}{2m}\Delta\psi - \mathcal{N}(\psi) \quad (1.3)$$

formally approximates the equation (1.1) in the nonrelativistic regime

$$|\mathbf{p}| = \hbar|\xi| \ll mc.$$

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