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# Normalized solutions and asymptotical behavior of minimizer for the coupled Hartree equations

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Received 4 June 2017; revised 1 March 2018

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## Abstract

We prove the existence and nonexistence of  $L^2(\mathbb{R}^N)$ -normalized solutions of coupled Hartree equations, which arise from the studies of the nonlinear optics and multi-component Bose–Einstein condensates. Under certain type trapping potentials, by proving some delicate energy estimates, we give a precise description on the concentration behavior of minimizer solutions of the system. Furthermore, an optimal blowing up rate for the minimizer solutions of the system is also proved.

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MSC: 35J61; 35J20; 35Q55; 49J40

Keywords: Coupled Hartree equations; Nonlocal interaction; Variational methods

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## 1. Introduction and main results

In the present paper we study the coupled nonlinear Hartree equations with nonlocal interaction in the following form:

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<https://doi.org/10.1016/j.jde.2018.03.003>

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$$\begin{cases} -\Delta u + V_1(x)u = \lambda_1 u + \mu_1 \left( \int_{\mathbb{R}^N} \frac{u^2(y)}{|x-y|^2} dy \right) u + \beta \left( \int_{\mathbb{R}^N} \frac{v^2(y)}{|x-y|^2} dy \right) u, & x \in \mathbb{R}^N, \\ -\Delta v + V_2(x)v = \lambda_2 v + \mu_2 \left( \int_{\mathbb{R}^N} \frac{v^2(y)}{|x-y|^2} dy \right) v + \beta \left( \int_{\mathbb{R}^N} \frac{u^2(y)}{|x-y|^2} dy \right) v, & x \in \mathbb{R}^N, \end{cases} \quad (1.1)$$

where  $u, v \in H^1(\mathbb{R}^N)$ ,  $\mu_1, \mu_2 > 0$ ,  $\lambda_1, \lambda_2 \in \mathbb{R}$  are the Lagrange multipliers,  $V_1(x)$  and  $V_2(x)$  are trapping potentials, and  $\beta \in \mathbb{R}$  is a coupling constant describing attractive or repulsive interactions.

The consideration of (1.1) is mainly motivated by recent studies on the nonlinear Hartree equation (NHE)

$$i \frac{\partial \psi}{\partial t} = -\frac{1}{2} \Delta \psi + V \psi - \chi \left( C(x) * |\psi|^2 \right) \psi, \quad x \in \mathbb{R}^3. \quad (1.2)$$

For identical and nonrelativistic basic particles (such as bosons or electrons) under the influence of an external potential and also two-body attractive interaction between two particles, the condensate in the mean field regime is governed by the NHE (see [9,10,12,16]). In (1.2),  $\psi$  is a radially symmetric two-body potential function defined in  $\mathbb{R}^3$  and  $*$  denotes the convolution in  $\mathbb{R}^3$ . The most typical external potential is the Coulomb function  $C(x) = |x|^{-1}$ . The equation (1.2) is also used in the description of the Bose–Einstein condensates, in which  $V$  is the trapping potential and the nonlocal interaction also describes the interaction between the bosons in the condensate [7,33,34]. When  $V = 1$ , the equation (1.2) is also known as the nonlinear Choquard equation (see [19,22,25]), and the equation (1.2) also arises from the model of wave propagation in a media with a large response length [1,17]. Recently, the papers [11,25,29,31] considered the stationary solution of (1.2) with the following generalized Choquard equation

$$-\Delta u + u = (J_\alpha * |u|^p) |u|^{p-2} u, \quad u \in H^1(\mathbb{R}^N), \quad (1.3)$$

where  $J_\alpha : \mathbb{R}^N \rightarrow \mathbb{R}$  is the Riesz potential defined at each point  $x \in \mathbb{R}^N \setminus \{0\}$  by

$$J_\alpha(x) = \frac{A_\alpha}{|x|^{N-\alpha}} \quad \text{and} \quad A_\alpha = \frac{\Gamma\left(\frac{N-\alpha}{2}\right)}{\Gamma\left(\frac{\alpha}{2}\right) \pi^{N/2} 2^\alpha}. \quad (1.4)$$

By using the moving plane method for the nonlocal problem (see [5,6]), Ma and Zhao [25] considered the existence and the uniqueness of the positive solutions of (1.3). By using the variational methods, Moroz and Van Schaftingen [29] proved the existence of the ground state solution of (1.3). Later on, Moroz and Van Schaftingen studied the existence of nontrivial solution of (1.3) with the lower critical case  $p = \frac{\alpha}{N} + 1$  in [31], subsequently Ghimenti and Van Schaftingen proved the existence of nodal and sign-changing solutions of (1.3) in [11]. One can see [30] for general nonlinearity case.

Systems of coupled nonlinear Schrödinger equations or Hartree equations have been the focus of many recent theoretical studies. The two-component nonlinear Schrödinger equations system with nonlocal Hartree type interaction can be written in the following form:

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