



# Enhanced thermoelectric performance of bulk tin telluride: Synergistic effect of calcium and indium co-doping

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## ABSTRACT

SnTe based materials are considered recently as a lead-free replacement of the well-known PbTe based thermoelectric (TE) materials in addressing the energy crisis worldwide. Herein we report both experimental and theoretical study on the effect of co-doping of calcium and indium on electronic structure and TE properties of SnTe. We show that the resonant levels introduced by indium and band gap opening caused by calcium, valence band convergence induced by both calcium and indium, synergistically increases the Seebeck coefficient for a wide range of temperatures. The co-doped SnTe with a high  $ZT$  of  $\sim 1.65$  at 840 K and record high power factor of  $\sim 47 \mu\text{Wcm}^{-1}\text{K}^{-2}$  for SnTe based materials make it a promising material for TE applications.

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## Introduction

Exhaustion of non-renewable fossil fuel resources caused due to its abuse has resulted in ever increasing energy crisis [1]. Materials for sustainable power generation is considered as a solution to this problem [2]. Thermoelectric materials which can directly and reversibly convert heat into electricity have attracted considerable attention in the scientific community as a pollution free technology [3,4]. The thermoelectric efficiency of a material is governed by its figure of merit  $ZT$  which is given by  $ZT = \sigma S^2 T / \kappa$ , where  $\sigma$  is electrical conductivity,  $S$  is Seebeck coefficient and  $\kappa$  is total thermal conductivity. Although at present all the materials are known have a nonzero thermoelectric effect, in most materials the  $ZT$  is too low to be useful [1]. Hence, there is a need to identify materials with higher efficiency than available at present [5–8]. Due to well-known interdependence of  $S$ ,  $\sigma$  and  $\kappa$ , developing strategies to improve material's average  $ZT$  is complicated [3]. Due to environmental toxicity of lead, PbTe based materials are not advised for large scale usage, despite it being the best known TE material [9,10]. Hence there is need to develop thermoelectric materials containing non-toxic elements [11–16]. SnTe, being a homologue of PbTe has

similar crystal and electronic structure. Though safer than PbTe, SnTe suffers a drawback of low  $ZT$  due to high carrier concentration resulting out of inherent Sn vacancies, smaller band gap ( $\sim 0.18$  eV) and higher separation between the valence sub-bands ( $\sim 0.3$  eV) than PbTe [17,18]. In order to improve the  $ZT$  of SnTe, electronic structure has been engineered using dopants such as Bi, Ca, Cd, Hg, Ga, In, Mg, Mn, Sr and co-dopants such as Ag-In, Bi-In, Cd-In, Hg-In, Mg-In, Mn-In [14,19–36]. These work involve improving the power factor by achieving carrier concentration optimization, valence band convergence, introduction of resonance level or all of them together [14,19–36]. Further the thermal conductivity is reduced by employing various nanostructuring [1,21,22,24,27–30]. While doping of Bi causes modulation of carrier concentration and an increase in density of state effective mass, In introduces resonance states around Fermi level, enhancing the TE properties of SnTe [19,32]. Dopants like Ca, Cd, Hg, Ga, Mg, Mn, Sr are known to bring about opening of band gap of SnTe and cause convergence of valence sub-bands to increase the Seebeck coefficient which in turn enhances the power factors of SnTe. Since fabricating the device requires higher  $ZT$  throughout the temperature range, recently co-doping was employed to get beneficial effects of dopants active at lower (In due to introduction of resonance states) as well as higher temperature (elements like Ag, Cd, Hg, Mg, Mn due to their ability to modify the electronic structure of SnTe) [14,31,33–36]. In spite of these efforts, the maximum  $ZT$  achieved is in the range of 1.3–1.6 with most of the work involving harmful materials like Cd and Hg

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[15,21,22,33–35]. Recently, we reported enhancement of thermoelectric performance of bulk SnTe by co-doping Mg and In [14]. The synergistic action of the two dopants led to a  $ZT$  maximum of  $\sim 1.5$  at 840 K. The study showed that the site occupied by the dopants play an important role in modification of its electronic structure. We know that for singly doped materials by far Ca doped SnTe holds the maximum  $ZT$  value of  $\sim 1.35$  at 873 K [20]. It would be of interest to know if Ca also goes into zinc blende site like Mg and whether the site occupied by Ca in SnTe has any effect on its electronic structure [14,26]. Since in some of the co-doping cases such as Bi-In and Mg-In reported by Zhang et al. we see that the  $ZT$  maximum achieved is lower than the singly doped Bi and Mg in SnTe, we wanted to see how Ca and In would interact, whether the co-doping would decrease the  $ZT$  than singly doped Ca in SnTe or would it increase the  $ZT$  due to synergistic action as in the case of Cd-In co-doped SnTe reported by Tan et al. and Mg-In co-doped SnTe previously reported by us [14,19,26,32,33]. Hence, we thought it is worthwhile to investigate the effect of co-doping of Ca and In in SnTe both experimentally through synthesis using a combination of self propagating high temperature synthesis (SHS) and direct current sintering (DCS) and theoretically through simulations using first principles density functional theory (DFT) approach to see if we can achieve further improvement in  $ZT$ . We show that the resonant levels introduced by In and band gap opening caused by Ca, valence band convergence caused by both In and Ca, synergistically increases the Seebeck coefficient for a wide range of temperatures without involving toxic elements like Cd and Hg. We show through DFT simulations that unlike Mg, Ca prefers rock salt site in SnTe [14]. Thus, co-doping of Ca and In results in enhancement of figure of merit  $ZT$  due to complimentary role of Ca and In in modifying the electronic structure of SnTe producing a record high power factor and  $ZT$  in SnTe based materials. The synergism shown here is much higher than any co-doped SnTe based materials including the recently reported Mg-In co-doped SnTe [14].

## Results and discussions

A series of samples  $\text{Sn}_{1.04-3x}\text{Ca}_{2x}\text{In}_x\text{Te}$  ( $x = 0$  to 0.04) was prepared by modified SHS and DCS process. The current strategies used to optimize the thermoelectric properties of SnTe largely depend on the solubility of the dopant in it [1,25]. Equilibrium phase diagram limits the solubility of the dopant in conventional solid solution alloying. Hence increasing the solubility is difficult. Since SHS is a non-equilibrium process higher solubility can be attained at room temperature [14,25]. SHS process is simple, quick and self-propagating due to the release of intense heat during the exothermic reaction. The combination of SHS-DCS process is cost effective and scalable for large scale production. Pristine SnTe is not considered as a good TE material due to large number of Sn vacancies. Self-compensation, addition of I or Bi have been proved to be a good way to reduce the carrier concentration in pristine SnTe [14,21,22]. Recently we reported an optimized concentration of  $\text{Sn}_{1.04}\text{Te}$  and the same is implemented here [14]. We then co-dope it with Ca and In maintaining 2:1 ratio unlike in the case of Mg-In co-doped SnTe where the ratio of Mg:In was 1:1 [14]. 2:1 ratio was chosen to take advantage of higher solubility of Ca and to avoid rapid decrease in electrical conductivity caused due to increase in In dopant concentration as explained in later sections [20,21,24].

The powder X-ray diffraction pattern of the synthesized compound revealed in Fig. 1a could be indexed to cubic NaCl structure with space group  $Fm\bar{3}m$ . Impurity peaks were absent within the detection limits confirming the formation of single phase. We see that the calculated lattice parameter value 'a' gradually increases with increase in concentration of Ca and In during co-doping due to larger ionic radius of  $\text{Ca}^{2+}$  in comparison with  $\text{Sn}^{2+}$  [20]. In

substitution does not have much effect on the lattice parameter of SnTe as InTe and SnTe have approximately similar lattice parameter ( $\sim 6.32 \text{ \AA}$ ) [14]. The linear increase in the lattice parameter with dopant concentration follows Vegard's law till 8 mol % of Ca (Fig. 1b).

Thermoelectric transport properties as a function of temperature were studied in the range 300 K–840 K. We observe that co-doping of Ca and In decreases the electrical conductivity values with increase in the concentration of the dopants and also with respect to increase in the temperature showing the degenerate nature of semiconducting materials (Fig. 2a) [22]. The values of electrical conductivity of the Ca-In co-doped samples are intermediate compared to the previously reported SnTe samples singly doped with Ca and In [20,24,25]. The room temperature electrical conductivity decreases from  $6600 \text{ Scm}^{-1}$  to  $1970 \text{ Scm}^{-1}$  as 'x' increases from 0 to 4 mol % in  $\text{Sn}_{1.04-3x}\text{Ca}_{2x}\text{In}_x\text{Te}$ . The carrier concentration ' $N_p$ ' of SnTe reduces from  $4.5 \times 10^{20} \text{ cm}^{-3}$  to  $2.1 \times 10^{20} \text{ cm}^{-3}$  on self compensation of Sn vacancies as established in the previous report [14,21]. The carrier concentration decreases to  $1.9 \times 10^{20} \text{ cm}^{-3}$  as the 'x' increases to 4 mol % in  $\text{Sn}_{1.04-3x}\text{Ca}_{2x}\text{In}_x\text{Te}$ . The details of carrier concentration and carrier mobility as a function of doping concentration are given in Fig. 2b. We observe that the carrier mobility ' $\mu$ ' decreases with increase in 'x' due to increase in scattering between the carriers, enhanced point defect scattering and increase in the effective mass of the carrier as explained in the later sections [20,33,37].

The Seebeck coefficient values show increasing trend with increase in temperature and dopant concentration (Fig. 3a). The positive value of Seebeck coefficient reveals that holes are the majority carriers leading to p-type conduction in the material [22]. The pristine and self compensated SnTe show a maximum value of Seebeck coefficient of  $110 \mu\text{VK}^{-1}$  and  $145 \mu\text{VK}^{-1}$  respectively. The co-doped sample shows room temperature Seebeck coefficient of  $95 \mu\text{VK}^{-1}$  and a maximum value of  $235 \mu\text{VK}^{-1}$  at 840 K (when  $x = 0.04$  in  $\text{Sn}_{1.04-3x}\text{Ca}_{2x}\text{In}_x\text{Te}$ ). To study the effect of co-dopants, we compared room temperature Seebeck coefficient versus  $N_p$  with the theoretical Pisarenko plot reported by Zhang et al. (Fig. 3b) [24]. The Pisarenko line is derived using two valence band model with light hole valence band having an effective mass of  $0.168 m_e$  and heavy hole valence band having an effective mass of  $1.92 m_e$  with the energy gap of 0.35 eV between the two sub-bands [24]. The Seebeck coefficient value for undoped SnTe falls exactly on top of the line indicating the validity of the theory. At room temperature the co-doped samples have significantly high Seebeck values than predicted by the Pisarenko line. The introduction of resonance levels by In increases the Seebeck coefficient values at lower temperatures by increasing the band effective mass  $m_b^*$  while, at higher temperatures the Seebeck value increases due to the decrease in the energy separation between the light hole and heavy hole valence sub-bands which is a consequence of Ca and In doping which is further studied making use of DFT calculations [1,14,24,25,31–36].

We carried out the simulations of pristine, Ca doped SnTe, In doped SnTe and Ca-In co-doped SnTe using a 32 atom supercell to achieve  $\sim 6.25$  mol % doping concentration. SnTe with a primitive rock salt crystal structure is known to possess an electronic structure with principal band gap at L point and it also has a heavy hole valence band at  $\Sigma$  point. Since our calculations involve a  $2 \times 2 \times 1$  supercell, the Brillouin zone folds accordingly and hence the states at L point in the primitive cell folds onto the Z point in the supercell while the states at  $\Sigma$  point corresponds to the states at  $M + \delta$  point along  $M \rightarrow \Gamma$  direction in the tetragonal supercell [14]. The electronic structure of SnTe which we obtained in our present calculation reveals a band gap of  $\sim 0.056 \text{ eV}$  at Z point, slightly higher than the  $0.049 \text{ eV}$  reported by Banik et al. [26] The theoretically

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