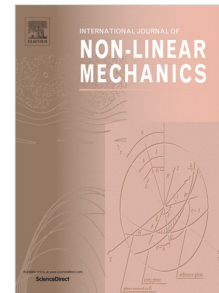


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Bounded and unbounded solutions of a discontinuous oscillator at resonance

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1 **BOUNDED AND UNBOUNDED SOLUTIONS OF A DISCONTINUOUS**
 2 **OSCILLATOR AT RESONANCE**

3 HEBAI CHEN^{1,2}, JINQIAO DUAN³

ABSTRACT. First, we have studied a specific discontinuous differential equation for a smooth and discontinuous (SD) oscillator $\ddot{x} + x(1 - 1/\sqrt{a^2 + x^2}) = p(t)$, where $p(t)$ is a given 2π -periodic forcing function and a is a real parameter. When the forcing $p(t) \in C^r(\mathbb{R}/2\pi\mathbb{Z})$ and $|\int_0^{2\pi} p(t)e^{-it} dt| < 4$, all solutions of the oscillator are shown to be bounded, i.e.,

$$\lim_{|t| \rightarrow +\infty} \{x(t)^2 + \dot{x}(t)^2\} < +\infty.$$

Moreover, the oscillator has at least one harmonic solution. When the forcing $p(t) \in C^0(\mathbb{R}/2\pi\mathbb{Z})$ and $|\int_0^{2\pi} p(t)e^{-it} dt| \geq 4$, all solutions of the oscillator are unbounded, i.e.,

$$\lim_{|t| \rightarrow +\infty} \{x(t)^2 + \dot{x}(t)^2\} = +\infty.$$

The two conditions are sharp because they complement each other. Moreover, a bifurcation that connects the bounded solutions with unbounded ones occurs, and 4 is a bifurcation value. Inspired by this special discontinuous oscillator, we have found the conditions for the existence of bounded and unbounded solutions for a more general discontinuous oscillator. Finally, we present interesting physical models to illustrate our results.

4 1. INTRODUCTION

5 Littlewood in [18, 19, 20] suggested to study the bounded and unbounded solutions of
 6 the following Liénard system for an oscillator $\ddot{x} + g(x) = p(t)$, with a nonlinear interaction
 7 g and periodic forcing p . A solution of this type of differential equation is called “bounded”
 8 if

$$\lim_{|t| \rightarrow +\infty} \{x(t)^2 + \dot{x}(t)^2\} < +\infty.$$

9 For example, the periodic solutions and quasi-periodic solutions are bounded solutions,
 10 where the definition of quasi-periodic solutions is shown in Section 2. A solution is called
 11 “unbounded” if

$$\lim_{|t| \rightarrow +\infty} \{x(t)^2 + \dot{x}(t)^2\} = +\infty.$$

12 In [15], Lazer and Leach studied the nonlinear perturbations of a forced harmonic oscillator
 13 at resonance and provided a necessary and sufficient condition for the existence of harmonic
 14 solutions (defined in the next section). Specifically, they considered the following equation

$$(1) \quad \ddot{x} + n^2x + \phi(x) = p(t),$$

15 where $n \in \mathbb{N}$, ϕ is a given continuous and bounded nonlinear function, and $p(t)$ is a given
 16 2π -periodic forcing. When

$$\left| \int_0^{2\pi} p(t)e^{-int} dt \right| < 2 \left(\liminf_{x \rightarrow +\infty} \phi(x) - \limsup_{x \rightarrow -\infty} \phi(x) \right),$$

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