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An optimized Runge–Kutta method for the solution of orbital problems

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Abstract

We present a new explicit Runge–Kutta method with algebraic order four, minimum error of the fifth algebraic order (the limit of the error is zero, when the step-size tends to zero), infinite order of dispersion and eighth order of dissipation. The efficiency of the newly constructed method is shown through the numerical results of a wide range of methods when these are applied to well-known periodic orbital problems. © 2004 Elsevier B.V. All rights reserved.

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1. Introduction

Many orbital problems in astronomy, astrophysics, celestial mechanics, etc. are expressed by the second-order differential equation of the form

$$y''(x) = f(x, y(x)), \quad y(x_0) = y_0, y'(x_0) = y'_0,$$
(1)

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that is, differential equations where f is independent from the first derivative of y. In order to use first-order numerical methods to solve these problems we set

$$y_1(x) = y(x),$$

 $y_2(x) = y'(x).$

In this way, (1) can be expressed by a system of two first-order ODEs

$$y'_{1}(x) = y_{2}(x),$$

$$y'_{2}(x) = f(x, y_{1}(x)).$$
(2)

2. Basic theory

2.1. General form of explicit Runge-Kutta methods

An *s*-stage explicit Runge–Kutta method used for the computation of the approximation of $y_{n+1}(x)$ in problem (2), when $y_n(x)$ is known, can be expressed by the following relations:

$$y_{n+1} = y_n + \sum_{i=1}^{s} b_i k_i,$$

$$k_i = hf\left(x_n + c_i h, y_n + h \sum_{j=1}^{i-1} a_{ij} k_j\right), \quad i = 1, \dots, s.$$
(3)

The method mentioned previously can also be presented using the Butcher table below:

The following equations must always hold:

c

$$c_i = \sum_{j=1}^{s} a_{ij}, \quad i = 2, \dots, s.$$
 (5)

Definition 1 (*Hairer et al.* [4]). A Runge–Kutta method has algebraic order p when the method's series expansion agrees with the Taylor series expansion in the p first terms

$$y^{(n)}(x) = y_n^{(n)}(x), \quad n = 1, 2, \dots, p.$$

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