

# $\Lambda$ -Doubling investigation of the $5^1\Pi_g$ Rydberg state of $\text{Na}_2$ using optical–optical double resonance spectroscopy

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## Abstract

The splitting of  $\Lambda$ -doubling in the  $5^1\Pi_g$  Rydberg state of  $\text{Na}_2$ , which dissociates to  $\text{Na}(3s) + \text{Na}(4d)$ , has been measured using the high-resolution cw optical–optical double resonance technique. The observed data are in the range of  $0 \leq v \leq 22$  and  $11 \leq J \leq 83$  with  $\Lambda$ -doubling revealed. A set of Dunham coefficients with  $\Lambda$ -doubling constants has been obtained from the experimental results. The splitting of  $\Lambda$ -doubling increases quadratically with the rotational quantum number  $J$  and weakly depends on the vibrational quantum number  $v$ . These splitting constants are much larger than those in the  $\text{Na}_2$   $B^1\Pi_u$  state, which dissociates to  $\text{Na}(3s) + \text{Na}(3p)$ . This indicates that the splitting of  $\Lambda$ -doubling in the  $5^1\Pi_g$  state is affected by both the perturbations by adjacent  $\Sigma$  states and the  $L$ -uncoupling.

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## 1. Introduction

In a diatomic molecule, the precession of electronic angular momentum  $L$  takes place along its internuclear axis with constant component  $M_L$  where  $M_L$  can take the values  $M_L = L, L-1, \dots, -L$  [1]. However, reversing the direction of precession does not change the energy of the state but does change  $M_L$  to  $-M_L$ . States with different  $M_L$  have in general a larger energy separation since the electric field which causes the separation is very strong. It is common to classify the electronic states of diatomic molecules according to the value of  $|M_L|$  and let  $\Lambda = |M_L|$ . The corresponding angular momentum vector  $\Lambda$  represents the component of the electronic orbital angular momentum  $L$  along the internuclear axis. For a given  $L$ , the quantum number  $\Lambda$  can take the values  $\Lambda = 0, 1, 2, \dots, L$ , corresponding to the designation of molecular states  $\Sigma, \Pi, \Delta, \dots$ , respectively. States  $\Pi, \Delta, \dots$  are doubly degen-

erate since  $M_L$  can have two values  $+\Lambda$  and  $-\Lambda$ , while states  $\Sigma$  are non-degenerate. When interaction with other electronic states is also considered, these doubly degenerate states are split by so-called  $\Lambda$ -doubling. Mathematically, the wavefunctions of these doubly degenerate states can be written as a phase factor  $e^{\pm i\Lambda\phi}$ , where  $\phi$  is the azimuthal angle rotating around the internuclear axis. The different angular momenta in the molecule—electron spin  $S$ , electronic angular momentum  $L$ , and angular momentum of nuclear rotation  $N^+$ —form a resultant that is always designated as  $J$ , the total angular momentum. The mutual interactions of these angular momenta are coupled in different ways according to their interaction strength and have been classified by Hund [2]. However, Hund's coupling cases represent only idealized limiting cases. In molecule, the rotation of nuclei  $N^+$  disturbs the energy of the states. For a larger speed of rotation, each total angular momentum  $J$  in the state of  $\Lambda \neq 0$  will split into two components. This splitting is called  $\Lambda$ -doubling and increases quadratically with  $J$ . This phenomenon of diatomic molecular spectra was experimentally and theoretically studied in the early years by Van Vleck [3] and Mulliken and Christy [4].

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For a molecule in a Rydberg state correlating with an atomic asymptote with smaller principal quantum number  $n$ , the electronic states with different  $A$  value are far apart in energy since the coupling of angular momentum to its internuclear axis is large and  $L$  precesses very rapidly. Each electronic state has its own set of rotational and vibrational energy levels. If one considers a series of orbits with increasing principal quantum number  $n$ , the influence of the internuclear axis on its electronic energy rapidly decreases when the orbit gets larger. The rate of precession of  $L$  and the energy interval between different  $A$  states diminishes quickly as the size of the orbit increases. Finally, the precession of  $L$  around the internuclear axis breaks down and  $A$  is no longer a good quantum number. In the Hund's coupling cases, transition occurs from case (a) or case (b) to case (d) while the coupling of  $L$  breaks down. This transition between Hund's coupling cases is shown in Fig. 1. Therefore, the degeneracy will be removed and the  $A$ -doubling splitting becomes larger as the size of the orbit increases. The transition between the Hund's coupling cases of the molecule NO has been reported by Klisch et al. [5]. To analyze the  $A$ -doubling splitting, the regular perturbation is applied by coupling the state of interest to the nearby electronic states whose  $A$  value differs by one [6]. The

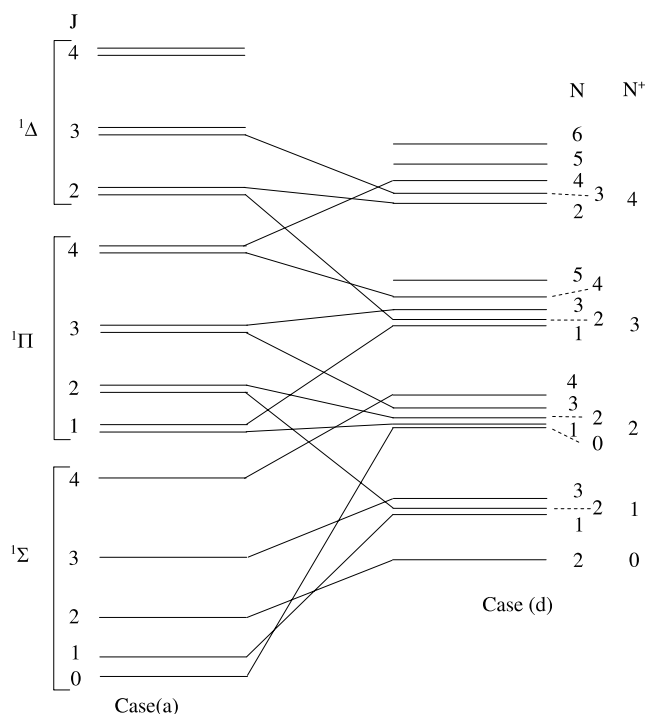


Fig. 1. Correlation diagram between Hund's coupling case (a) and case (d) for the  $s + d$  atomic asymptote. In case (d), coupling of the total electronic spin to the remaining angular momentum is not significant. So, only singlet states, i.e.,  $S = 0$ , are considered in this diagram (if  $S = 0$ , then  $J = N$ ). In case (a),  $\Sigma$ ,  $\Pi$ , and  $\Delta$  represent the possible states and  $J$  is the rotational angular momentum. For states with  $A \neq 0$ , each of the rotational levels  $J$  splits into two components of  $e/f$  parities known as  $A$ -doubling.

observations of  $A$ -doubling have been reported for molecules such as, NO,  $H_2$ ,  $Li_2$ , and NaK [7,8,9]. In the sodium dimer,  $A$ -doubling has been reported only in the  $B^1\Pi_u$  state with the first-order splitting constant  $q_0 = 0.115 \times 10^{-4} \text{ cm}^{-1}$ . In this report, the separation between the  $e/f$  levels of  $A$ -doubling in the  $5^1\Pi_g$  state has been measured by using high-resolution cw optical-optical double resonance (OODR) spectroscopy. The first-order splitting constant is 15 times larger than that in the  $B^1\Pi_u$  state. The  $A$ -doubling splitting increases as a quadratic function of rotational quantum number  $J$  and the dependence on the vibrational quantum number  $v$  is weak. The large  $A$ -doubling splitting constants indicate that the effects are from the perturbations between the adjacent electronic states as well as the uncoupling of orbital angular momentum  $L$ .

## 2. Experiment

The experimental setup has been described in more detail in [10]; hence we only dwell on it briefly. As shown in Fig. 2, sodium metal is loaded into a five-arm stainless steel heat-pipe oven with about 1 Torr of argon buffer gas and heats to  $350^\circ\text{C}$  at the center of the heat-pipe oven to generate sodium vapor. A single line  $\text{Ar}^+$  laser (Coherent I-90, total of nine lines) is used to excite the  $B^1\Pi_u(v', J') \leftarrow X^1\Sigma_g^+(v'', J'')$  transitions of  $\text{Na}_2$ . Several intermediate levels of  $B^1\Pi_u$  state can be excited simultaneously by one of the lines from the  $\text{Ar}^+$  laser. Table 1 lists the  $\text{Ar}^+$  laser wavelength (in air), transitions, and the term values of  $B^1\Pi_u$  state used in this study. The term values of the intermediate  $B^1\Pi_u$  state are calculated from the set III molecular constants in Table VII of [11]. We did not add the photon energy of the pump  $\text{Ar}^+$  laser to the term value of the ground state to form the term values of  $B^1\Pi_u$  state since the  $\text{Ar}^+$  laser frequency

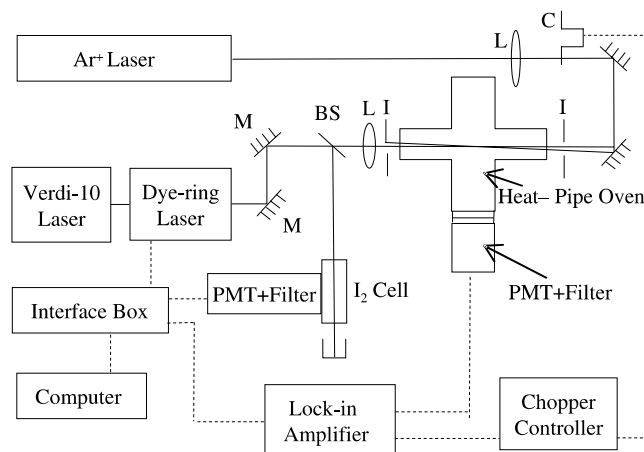


Fig. 2. The block diagram of experimental setup for OODR spectroscopy. The  $\text{Ar}^+$  laser beam is tilted after passing through the center of the heat-pipe to avoid interference with the dye laser. (M, mirror, BS, beam splitter, L, lens, I, iris, and C, chopper.)

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