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# The ${}^8\text{Li}(d, p){}^9\text{Li}$ reaction and astrophysical ${}^8\text{B}(p, \gamma){}^9\text{C}$ reaction rate

B. Guo, Z.H. Li\*, W.P. Liu, X.X. Bai, G. Lian, S.Q. Yan, B.X. Wang,  
S. Zeng, J. Su, Y. Lu

*China Institute of Atomic Energy, P.O. Box 275(46), Beijing 102413, PR China*

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## Abstract

Angular distribution of the  ${}^8\text{Li}(d, p){}^9\text{Li}_{\text{g.s.}}$  reaction at  $E_{\text{cm}} = 7.8$  MeV was measured in inverse kinematics. The square of asymptotic normalization coefficient (ANC) for the virtual decay  ${}^9\text{Li} \rightarrow {}^8\text{Li} + n$  was derived to be  $1.33 \pm 0.33 \text{ fm}^{-1}$  through distorted wave Born approximation (DWBA) analysis, for the first time. According to charge symmetry,  $(\text{ANC})^2$  for  ${}^9\text{C} \rightarrow {}^8\text{B} + p$  was then extracted to be  $1.14 \pm 0.29 \text{ fm}^{-1}$ . We have deduced the astrophysical S-factors and reaction rates for direct capture in  ${}^8\text{B}(p, \gamma){}^9\text{C}$  at energies of astrophysical relevance using the ANC for  ${}^9\text{C} \rightarrow {}^8\text{B} + p$  extracted from the mirror system.

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Keywords: NUCLEAR REACTIONS  ${}^2\text{H}({}^8\text{Li}, {}^9\text{Li})$ ,  $E_{\text{cm}} = 7.8$  MeV; measured  $\sigma(\theta)$ , DWBA analysis; deduced asymptotic normalization coefficient; mirror systems;  ${}^8\text{B}(p, \gamma)$ ,  $E = \text{low}$ ; deduced astrophysical S-factor

## 1. Introduction

Nucleosynthesis of light nuclei is impeded by the gap at mass number  $A = 8$ , where no stable nuclei exist. In some astrophysical environments, however, this gap

\* Corresponding author.

E-mail address: zhli@iris.ciae.ac.cn (Z.H. Li).

can be bypassed via the reactions involving unstable nuclei  ${}^8\text{B}$  and  ${}^8\text{Li}$ , such as  ${}^8\text{B}(p, \gamma){}^9\text{C}$ ,  ${}^8\text{Li}(\alpha, n){}^{11}\text{B}$ ,  ${}^8\text{Li}(n, \gamma){}^9\text{Li}$  and  ${}^8\text{Li}(d, p){}^9\text{Li}$ , to synthesize  $A > 8$  nuclides. The  ${}^7\text{Be}(p, \gamma){}^8\text{B}(p, \gamma){}^9\text{C}(\alpha, p){}^{12}\text{N}(p, \gamma){}^{13}\text{O}$  reaction chain is considered as one of the possible alternative paths to the  $3\alpha$  process for transforming the nuclei in the  $pp$  chains to the CNO nuclei in the peculiar astrophysical sites where the densities ( $\geq 2 \times 10^7 \text{ g/cm}^3$ ) and temperatures ( $T_9 \sim 0.1\text{--}0.4$ ) are so high that the proton- and  $\alpha$ -capture reactions become faster than the competing  $\beta$ -decays [1]. The  ${}^8\text{B}(p, \gamma){}^9\text{C}$  reaction may play an important role in the evolution of massive stars with very low metallicities [1,2], and thus has increasingly attracted both theoretical and experimental studies [1,3–8]. There are several microscopic and systematic calculations, and their results are in large discrepancy [1,3,4]. As for the experiments, it is very difficult to directly measure this reaction at energies of astrophysical relevance because of very small cross section and low intensity of the available  ${}^8\text{B}$  beam at present. Some indirect approaches have been applied to the study of this reaction [5–8]. Beaumel et al. [5] measured the  ${}^8\text{B}(d, n){}^9\text{C}$  angular distribution in inverse kinematics with a 14.4 A MeV  ${}^8\text{B}$  beam, and then derived the ANC for the virtual decay  ${}^9\text{C} \rightarrow {}^8\text{B} + p$  and the astrophysical  $S_{18}(0)$  factor for the  ${}^8\text{B}(p, \gamma){}^9\text{C}$  reaction. Trache et al. [6] analyzed the cross section data for one-proton-removal reaction of  ${}^9\text{C}$  on four different targets (C, Al, Sn and Pb) [9], and employed the Glauber model [10] to deduce the ANC and  $S_{18}(0)$  factor. Recently, Motobayashi [7] extracted the  $S_{18}$  factors in energy range 0.2–0.6 MeV by Coulomb dissociation approach ( $S_{18}(0)$  factor can be then obtained through an extrapolation by the slope of theoretical S-factor curve). Most recently, Enders et al. [8] studied the proton-removal from  ${}^9\text{C}$  on a carbon target at  $E = 78.3$  A MeV and derived the ANC and astrophysical  $S_{18}(0)$  factor. The  $S_{18}(0)$  obtained from Ref. [7] is significantly larger than other three ones.

The  ${}^8\text{Li}(d, p){}^9\text{Li}$  reaction not only leads to the production of  ${}^9\text{Be}$  (via the  ${}^9\text{Li}$   $\beta$ -decay) which acts as a precursor to heavier nuclides, but also can serve as a surrogate reaction to extract the  ${}^8\text{B}(p, \gamma){}^9\text{C}$  and  ${}^8\text{Li}(n, \gamma){}^9\text{Li}$  reaction rates for the direct capture. To date, only a few experiments for the  ${}^8\text{Li}(d, p){}^9\text{Li}$  reaction have been carried out by using the secondary  ${}^8\text{Li}$  beam. An earlier measurement, performed at  $E_{\text{cm}} = 1.5\text{--}2.8$  MeV [11], presented an upper limit of the cross section, though no  ${}^9\text{Li}$  event was detected. Very recently, the angular distributions for different states in  ${}^9\text{Li}$  were measured at  $E({}^8\text{Li}) = 76$  MeV to obtain information on the spins, parities and single-neutron spectroscopic factors [12]. In the present work, we measured the  ${}^8\text{Li}(d, p){}^9\text{Li}_{\text{g.s.}}$  angular distribution at  $E({}^8\text{Li}) = 39$  MeV through the coincidence detection of  ${}^9\text{Li}$  and recoil proton, and derived the ANC for the virtual decay  ${}^9\text{Li} \rightarrow {}^8\text{Li} + n$  based on DWBA analysis, and then deduced the ANC for  ${}^9\text{C} \rightarrow {}^8\text{B} + p$  based on charge symmetry. We have also calculated the direct capture S-factors and reaction rates for  ${}^8\text{B}(p, \gamma){}^9\text{C}$  at astrophysically relevant energies. Most recently, a short paper concerning the  ${}^8\text{Li}(d, p){}^9\text{Li}_{\text{g.s.}}$  angular distribution and indirect determination of the astrophysical  ${}^8\text{Li}(n, \gamma){}^9\text{Li}_{\text{g.s.}}$  reaction rates has been published elsewhere [13].

## 2. Experimental procedure and results

The measurement of  ${}^8\text{Li}(d, p){}^9\text{Li}$  angular distribution was performed using the secondary beam facility GIRAFFE [14,15] built at the HI-13 tandem accelerator of China

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