



# Fundamental understanding of wave generation and reception using $d_{36}$ type piezoelectric transducers



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## ABSTRACT

A new piezoelectric wafer made from a PMN–PT single crystal with dominant piezoelectric coefficient  $d_{36}$  is proposed to generate and detect guided waves on isotropic plates. The in-plane shear coupled with electric field arising from the piezoelectric coefficient is not usually present for conventional piezoelectric wafers, such as lead zirconate titanate (PZT). The direct piezoelectric effect of coefficient  $d_{36}$  indicates that under external in-plane shear stress the charge is induced on a face perpendicular to the poled  $z$ -direction. On thin plates, this type of piezoelectric wafer will generate shear horizontal (SH) waves in two orthogonal wave propagation directions as well as two Lamb wave modes in other wave propagation directions. Finite element analyses are employed to explore the wave disturbance in terms of time-varying displacements excited by the  $d_{36}$  wafer in different directions of wave propagation to understand all the guided wave modes accurately. Experiments are conducted to examine the voltage responses received by this type of wafer, and also investigate results of tuning frequency and effects of  $d_{31}$  piezoelectric coefficient, which is intentionally ignored in the finite element analysis. All results demonstrate the main features and utility of proposed  $d_{36}$  piezoelectric wafer for guided wave generation and detection in structural health monitoring.

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## 1. Introduction

Ultrasonic guided waves are being used extensively for effectively detecting the damage in structural health monitoring (SHM) [1,2]. The guided waves reflect, transmit or scatter, often with additional converted wave modes when they encounter discontinuous, such as geometric boundaries, flaws, or damages, therefore methods for damage detection are mainly based on analysis of waveform received by transducers at fairly high ultrasonic frequency range [3,4].

Conventional guided waves include Lamb waves with symmetric and antisymmetric modes, and shear horizontal (SH) waves in thin plate-like structures, while in elongated cylindrical structures like rods, pipes, cable strands or fibers, longitudinal modes wave, flexural modes wave and torsional modes wave propagate along

their long axes [5]. For most of wave modes, one of the fundamental characteristics is the existence of dispersion, which gives the nonlinear relationship between wavenumber and frequency. As a consequence, signals with a wide bandwidth will become significantly distorted as the propagation distance or time duration increases and further attenuated in amplitude due to geometric spreading. Therefore, dispersion results in complex signal processing in the application of guided waves based SHM methods.

A narrowband signal in small compact time duration is usually used to generate the guided wave to minimize influence of the dispersion, yet it still prevails for long propagation distances. One of the solutions is to utilize non-dispersive guided wave, such as fundamental shear horizontal ( $SH_0$ ) waves in thin plates or the first mode torsional waves in cylindrical structures. They are usually generated by a special type of transducers, such as magnetostrictive transducers e.g., [6,7] and thickness shear ( $d_{15}$ ) mode piezoelectric transducers [8], which can readily generate non-dispersive guided waves. Seung et al. proposed an omni-directional SH wave magnetostrictive patch transducer that consists of an annular magnetostrictive patch [9], a toroidal coil and a permanent magnet. However these techniques require additional confinement

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in order to generate such waves, resulting in a larger footprint than miniature piezoelectric wafers, which even can be embedded into the structures.

Another reason leading to complicated waveform is multi-mode of guided waves. For a given frequency, a finite number propagating wave modes will be generated and propagate simultaneously. This issue can be overcome by frequency tuning technique, which is to seek for a specific frequency range to amplify or decrease the amplitude of response with respect to various frequencies for different modes wave respectively [10].

A new piezoelectric wafer has been proposed to generate and detect the SH waves [11]. This preliminary investigation also demonstrated its ability to detect the damage. This special behavior owes to the piezoelectric coefficient  $d_{36}$  in its piezoelectric matrix. The direct piezoelectric effect of piezoelectric coefficient  $d_{36}$  indicates that under external in-plane shear stress  $\sigma_{12}$ , the charge is induced on a face perpendicular to the poled  $z$ -direction. The corresponding converse piezoelectric effect is when the external electric field is applied in the  $z$ -direction, the response is the in-plane shear deformation experienced by the material on the face. Toward the practical application, this paper investigates further the directionality and the tuning frequency characteristic of the proposed piezoelectric wafer by simulation and experiments.

## 2. Brief characteristics of guided waves in plates

For a plate bounded by the surfaces  $x_3 = \pm h/2$  and is of infinite extent in the  $x_1$  and  $x_3$  directions, the harmonic wave motion can be divided into two classes of wave motions: plane strain and antiplane shear (or shear horizontal) motions. For harmonic wave motion in plane strain in the  $(x_1, x_3)$ -plane of an elastic plate, the guided wave field can be represented by a standing wave in the  $x_2$  direction and a propagating wave in the  $x_1$  direction, which is known as Lamb waves. Lamb waves can be separated into symmetric and antisymmetric modes. One of the most discussed issues is their dispersion property given by:

$$\frac{\omega^4}{c_T^2} = 4k^2 q^2 \left[ 1 - \frac{p \tan(ph/2 + \gamma)}{q \tan(qh/2 + \gamma)} \right] \quad (1)$$

where  $\omega$  is angular frequency,  $c_T$  is transverse wave velocity,  $k$  is wave number,  $q = \sqrt{\omega^2/c_T^2 - k^2}$ ,  $p = \sqrt{\omega^2/c_L^2 - k^2}$ ,  $h$  is the thickness of the plate,  $\gamma = 0$  and  $\pi/2$  represent symmetric and antisymmetric modes, respectively. The dispersion relation results in an infinite number of wave modes defined by  $S_0, S_1, S_2, \dots, A_0, A_1, A_2, \dots$  for symmetric and antisymmetric modes, respectively. The corresponding displacement components for wave modes can be expressed as:

$$\begin{aligned} u &= q \left[ \cos(qx_3 + \gamma) - \frac{2k^2 \cos(qh/2 + \gamma)}{k^2 - q^2 \cos(ph/2 + \gamma)} \cos(px_3 + \gamma) \right] B \\ w &= -ik \left[ \sin(qx_3 + \gamma) + \frac{2pq \cos(qh/2 + \gamma)}{k^2 - q^2 \cos(ph/2 + \gamma)} \sin(px_3 + \gamma) \right] B \end{aligned} \quad (2)$$

where  $u$  and  $w$  are displacement components in  $x_1$  and  $x_3$ , respectively.  $B$  is a constant. Therefore, for Lamb waves, it only has displacement components  $u$  and  $w$ , and the displacement components  $v$  does not exist.

Another type of guided wave in the plate, SH wave, involves antiplane shear motion. In a plate with surfaces normal to  $x_3$ , the wave polarized in the  $(x_1, x_3)$ -plane with components  $u$  and  $w$  is decoupled from the SH wave polarized along  $x_2$  with displacement  $v$  alone. Its dispersion relation can be written in an analytical form as:

$$\left( \frac{\omega h}{c_T} \right)^2 = (kh)^2 + (n\pi)^2 \quad (3)$$

where  $n = 0, 2, 4, \dots$  for symmetric modes, and  $n = 1, 3, 5, \dots$  for antisymmetric modes. Different with all modes of Lamb waves and the other modes SH waves, the fundamental mode (SH<sub>0</sub>) of SH wave is non-dispersive. The SH waves motion is normal to the direction of propagation,  $x_1$ . Its only displacement component  $v$  can be written as:

$$v = \begin{cases} B_n \cos\left(\frac{n\pi x_3}{h}\right) (\text{symmetric modes}) \\ A_n \sin\left(\frac{n\pi x_3}{h}\right) (\text{antisymmetric modes}) \end{cases} \quad (4)$$

where  $A_n$  and  $B_n$  are unknown coefficients. Note that the amplitude of SH modes are independent of frequency and wavenumber. Hence the wave structure of the SH mode does not vary along the entire dispersion curve. This is in contrast to Lamb wave behavior, where the wavefield is a function of the position on the dispersion curve. For SH wave excitation, the shear deformation is necessary for the sensor or actuator working in plane, thus traditional piezoelectric is unable to generate and sense it.

## 3. $d_{36}$ Type piezoelectric wafers

The most widely used piezoelectric material in most of the transducer and actuator applications is lead zirconate titanate (PbZr(Ti)O<sub>3</sub>, PZT), a ferroelectric ceramic. When this material is poled, it attains a permanently polarized state. For a small variation in the electric field, it behaves approximately linearly near that state, which can be expressed by the following linear piezoelectric equations:

$$\begin{aligned} \varepsilon_{ij} &= c'_{ijkl} \sigma_{kl} + d_{kij} E_k \\ D_i &= d_{ikl} \sigma_{kl} + \varepsilon_{ik}^T E_k \end{aligned} \quad (5)$$

where  $\varepsilon_{ij}$  and  $\sigma_{kl}$  are the strain and stress tensors,  $c'_{ijkl}$ ,  $d_{kij}$ , and  $\varepsilon_{ik}^T$  are the inverse elastic, piezoelectric, and dielectric constants, respectively.  $E_k$  and  $D_i$  are electric field and electric displacement component respectively. The poled piezoelectric ceramic is an orthotropic material with a plane of symmetry whose normal is in the poled direction, so the electric constants matrix has only three independent piezoelectric coefficients,  $d_{31}$ ,  $d_{33}$  and  $d_{15}$ . All other coefficients are zero, as following:

$$d_{ij} = \begin{bmatrix} 0 & 0 & 0 & 0 & d_{15} & 0 \\ 0 & 0 & 0 & d_{24} & 0 & 0 \\ d_{31} & d_{32} & d_{33} & 0 & 0 & 0 \end{bmatrix} \quad (6)$$

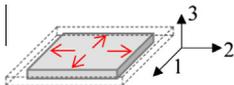
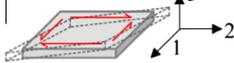
where  $d_{31} = d_{32}$ , and  $d_{24} = d_{15}$ . The piezoelectric wafer made of this piezoelectric ceramic is called  $d_{31}$  type piezoelectric wafer in this paper.

The proposed piezoelectric wafer in this paper has five independent coefficients, shown as following:

$$d'_{ij} = \begin{bmatrix} 0 & 0 & 0 & d_{14} & d_{15} & 0 \\ 0 & 0 & 0 & d_{24} & d_{25} & 0 \\ d_{31} & d_{32} & d_{33} & 0 & 0 & d_{36} \end{bmatrix} \quad (7)$$

**Table 1**

Comparison of deformations induced by  $d_{31}$  and  $d_{36}$  coefficients.

Piezoelectric coefficient	Deformation	Strain
$d_{31}/d_{32}$		$\varepsilon_1/\varepsilon_2$
$d_{36}$		$\gamma_{12}$

where  $d_{31} = d_{32}$ ,  $d_{14} = d_{25}$ , and  $d_{24} = d_{15}$ . For this wafer, when an external electric field is applied to the wafer in the  $z$  direction, the in-plane strain will include not only the normal strain, which is induced by  $d_{31}$  and  $d_{32}$ , but also the shear strain, which is induced by  $d_{36}$ . Table 1 lists the detailed difference between these two coefficients.

The proposed piezoelectric wafer may be obtained by means of cutting a type of single crystal,  $\text{Pb}(\text{Mg}_{1/3}\text{Nb}_{2/3})\text{O}_3\text{-PbTiO}_3$  (PMN-PT), in a special direction, shown as Fig. 1, in which  $[hkl]$  in crystallography denotes a direction and  $(hkl)$  a plane orthogonal to the direction  $[hkl]$ . The PMN-PT crystal first is orientated in a Cartesian coordinate system with crystal growth direction parallel to  $z$  axis. Then the PMN-PT crystal is poled in the  $[011]$  direction, and the wafer is cut in the plane  $(011)$ . Zhang et al. [12] gives an example of this piezoelectric constant matrix as following:

$$d'_{ij} = \begin{bmatrix} 0 & 0 & 0 & -696 & 935 & 0 \\ 0 & 0 & 0 & 935 & -696 & 0 \\ -364 & -364 & 890 & 0 & 0 & -1648 \end{bmatrix} \times 10^{-12} \text{C/N} \quad (8)$$

where  $d_{36}$  is more than four times of  $d_{31}$  and  $d_{32}$ . Generally, conventional  $d_{31}$  type piezoelectric wafer is able to excite Lamb waves (includes  $A_0$  and  $S_0$  modes) on the thin plate. Thanks to the shear deformation induced by  $d_{36}$ , the proposed piezoelectric wafer is expected to excite both Lamb waves and shear horizontal waves.

When the piezoelectric wafer is used routinely, the direction of external electric field  $E$  is usually along the  $z$ -axis and is perpendicular to the surface of the wafer, which is bonded on the surface of plate structure. For the moment, only  $d_{36}$  is considered, and then Eq. (5) can be simplified as:

$$\begin{aligned} \sigma_{12} &= \frac{E}{1+\nu} (\epsilon_{12} - d_{36}E_3) \\ D_3 &= d_{36}\sigma_{12} + \epsilon_3^T E_3 \end{aligned} \quad (9)$$

For a free  $d_{36}$  type piezoelectric wafer without in-plane external electric field and applied stress, the induced in-plane strain by applying a voltage across the  $z$ -direction can be expressed as:

$$\epsilon_{12} = d_{36}E_3 = \frac{d_{36}}{h_p} V_{in} \quad (10)$$

where  $h_p$  is the thickness of the wafer. Fig. 1(b) shows the deformation of a  $d_{36}$  type piezoelectric wafer placed on the surface of a plate. The deformation can be considered due to a pair of shear strains,  $\gamma_{xy}$  and  $\gamma_{yx}$ . Moreover, the wafer can be equivalent to a group of line force, which is along the four edges of the wafer, applied on the plate. The induced pure shear deformation along the edges results in the guided waves in the plate.

#### 4. Finite element analysis of $d_{36}$ type wafer as actuator and sensor

##### 4.1. Guided waves generation by $d_{36}$ type wafer

Through analyzing the displacement components, different guided wave modes can be understood accurately. Independent displacement components may be extracted naturally from finite element results. In this work, a metallic plate bonded with a  $d_{36}$  actuator was modeled with ANSYS to conduct the wave generation and propagation analyses. Fig. 2 shows the aluminum plate with 1.0 mm in thickness, 18 cm in both length and width. With this dimension, all guided waves will be able to prevail and the extensive computation is avoided. The square shaped actuator (4 mm  $\times$  4 mm), which is modeled by a Solid5 element in ANSYS, was placed on the location 70 mm away from two nearer edges of the plate. Seven sensing points are positioned on a circle with 60 mm radius around the excitation point with 15° increment. Due to the geometrical symmetry only a quadrant of the circumference is need to capture all guided wave phenomena from the square-shaped actuator located at the center. Displacement components of the sensing point were extracted from the simulation results to explore their variation with respect to different directions. To highlight wave generation using  $d_{36}$  type piezoelectric wafer, the relatively small values of piezoelectric coefficients related to the induced in-plane normal deformation are intentionally removed from the piezoelectric material matrix in Eq. (4).

The excitation from the piezoelectric wafer is a 5-cycle sinusoid tone-burst signal enclosed in a Hanning window with a center frequency  $f_c$  kept fixed at 160 kHz, since according to results from  $d_{31}$  type of PZT, both  $A_0$  and  $S_0$  modes of Lamb waves appear simultaneously under the excitation of this frequency. It can be seen through frequency domain analysis that the frequency components for this narrowband signal are mainly concentrated in a small region around the center frequency  $f_c$ , thus the dispersive effect can be significantly reduced. The amplitude spectrum of

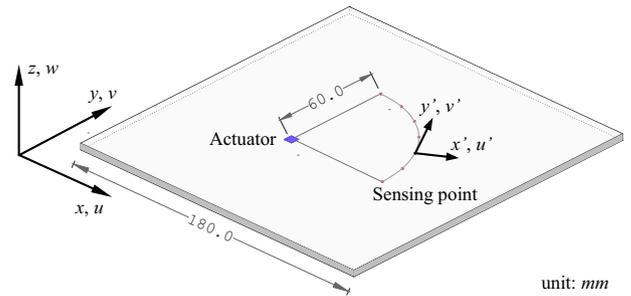


Fig. 2. Layout of the piezoelectric wafer and the sensing point.

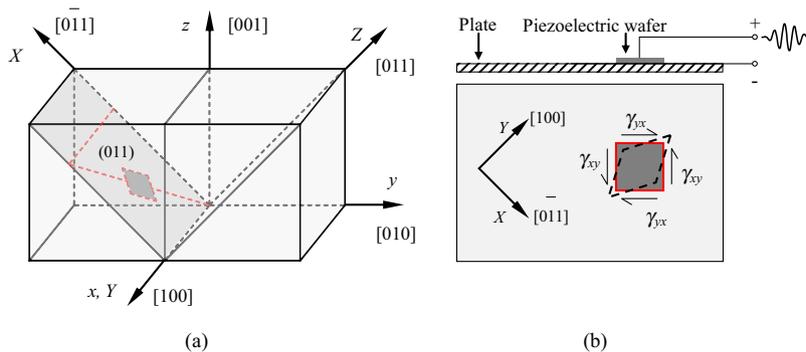


Fig. 1. Cutting direction of poled PMN-PT material and in-plane shear deformation of the PMN-PT mounted on a plate.

excitation signal was plotted together with group velocity dispersion curves corresponding to three wave modes in Fig. 3.

The size of the elements is determined according to the wavelength responding to the lowest phase velocity among  $A_0$ ,  $S_0$  and  $SH_0$  waves, i.e. the size of elements is less than  $1/20$  the shortest wavelength, while the time step during finite element analysis is chosen according to the center frequency  $f_c$  of excitation signal, i.e. the time step is less than  $1/20f_c$ .

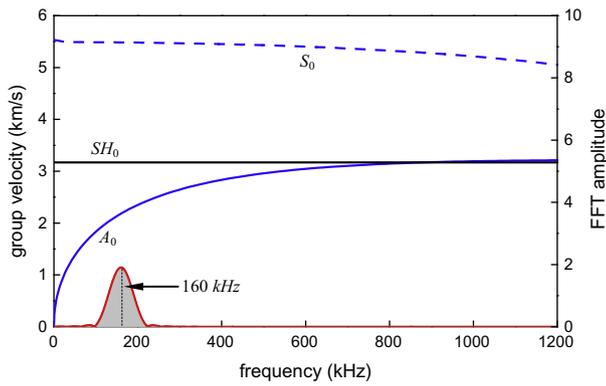


Fig. 3. Group velocity of Lamb wave and amplitude spectrum of excitation signal.

As shown in Fig. 2, for each sensing point, local coordinate ( $x'$  and  $y'$ ) is used, in which  $x'$  axis is along the radial direction; i.e., the direction of wave propagation. Then, in the local coordinate, all Lamb wave modes should have only two displacement components,  $u'$  and  $w'$ , based on the definition of Lamb wave described in Eq. (2); while the displacement component  $v'$  should indicate SH wave described in Eq. (4). Fig. 4 shows the snapshots of the displacement contours on the surface of the plate in three directions, and the total displacement contour, excited by the  $d_{36}$  actuator at the time  $28 \mu\text{s}$  respectively. For each figure, the displacement components are transformed to the local cylindrical coordinate, in which the displacement components  $u'$  for all nodes are in accordance with the radial direction, and  $v'$  the tangential direction. Lamb waves can be identified from Fig. 4(a) and (c) with significant values of  $u'$  and  $w'$ ; while SH wave also can be readily identified in Fig. 4(b). All values in Fig. 4 are extracted, and plotted in Figs. 5 and 6. Fig. 5 shows three displacement components that transform to the local coordinate at the sensing points from  $0^\circ$  to  $90^\circ$ . The range of  $y$  axis is kept fixed for illustrating their relative magnitudes. Fig. 5(a) shows three displacement components for sensing point on  $0^\circ$  and  $90^\circ$ . In this figure, both  $u$  and  $w$  vanish, and the displacement only occurs in  $v$ . This means that only SH wave is excited, and  $A_0$  and  $S_0$  Lamb waves do not exist in these two directions. According to the wave group velocity, the displacement  $v$  can be also identified as the direct  $SH_0$  waves, however their phases are

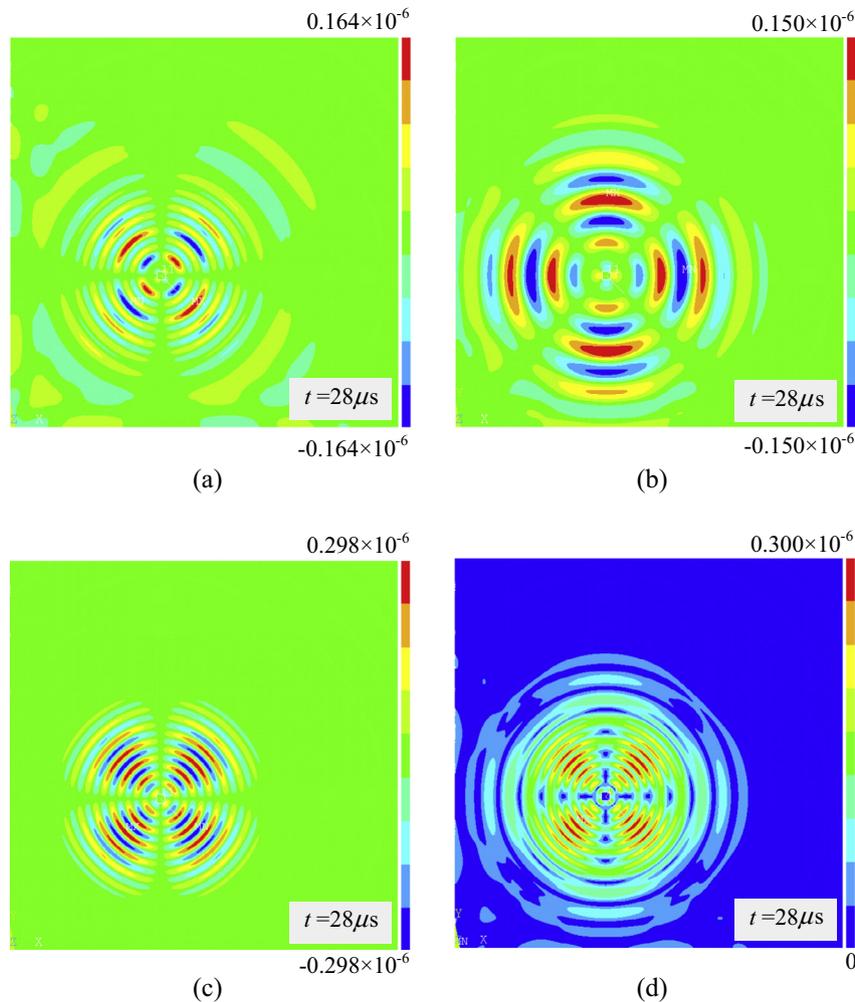


Fig. 4. Displacement contour in local coordinate under excitation of  $d_{36}$  type piezoelectric wafer: (a) displacement component  $u'$ , (b) displacement component  $v'$ , (c) displacement component  $w'$  and (d) total displacement.

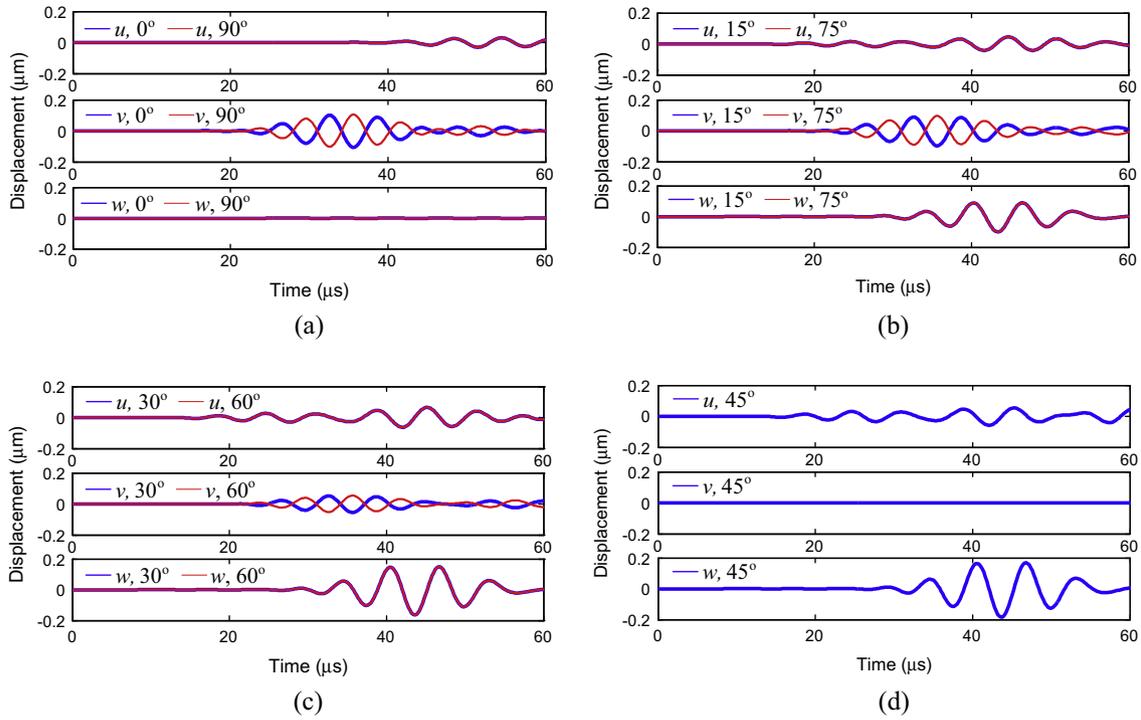


Fig. 5. Group velocity of Lamb wave and amplitude spectrum of excitation signal: (a) 0° and 90°, (b) 15° and 75°, (c) 30° and 60° and (d) 45°.

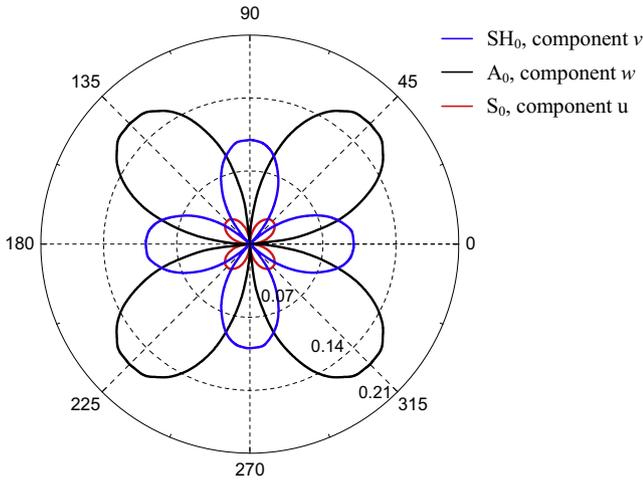


Fig. 6. Displacement amplitude versus direction of wave propagation for  $A_0$ ,  $S_0$ , and  $SH_0$  modes.

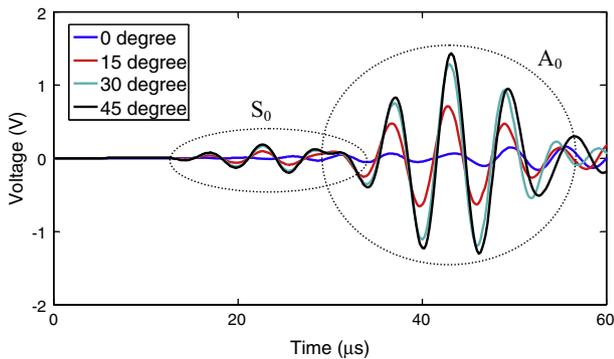


Fig. 7. Voltage responses of  $d_{36}$  type sensor to Lamb waves from different angles of wave propagation.

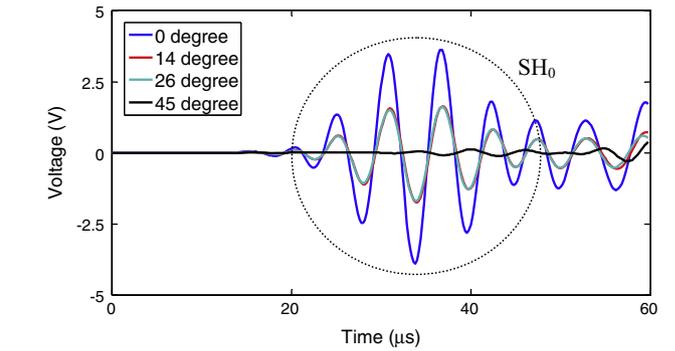


Fig. 8. Voltage responses of  $d_{36}$  type sensor to SH waves with different propagation directions.

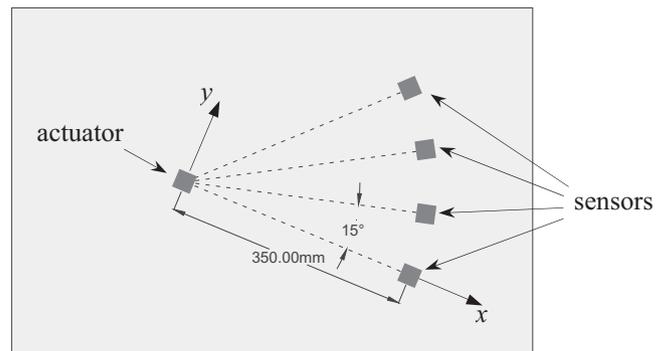
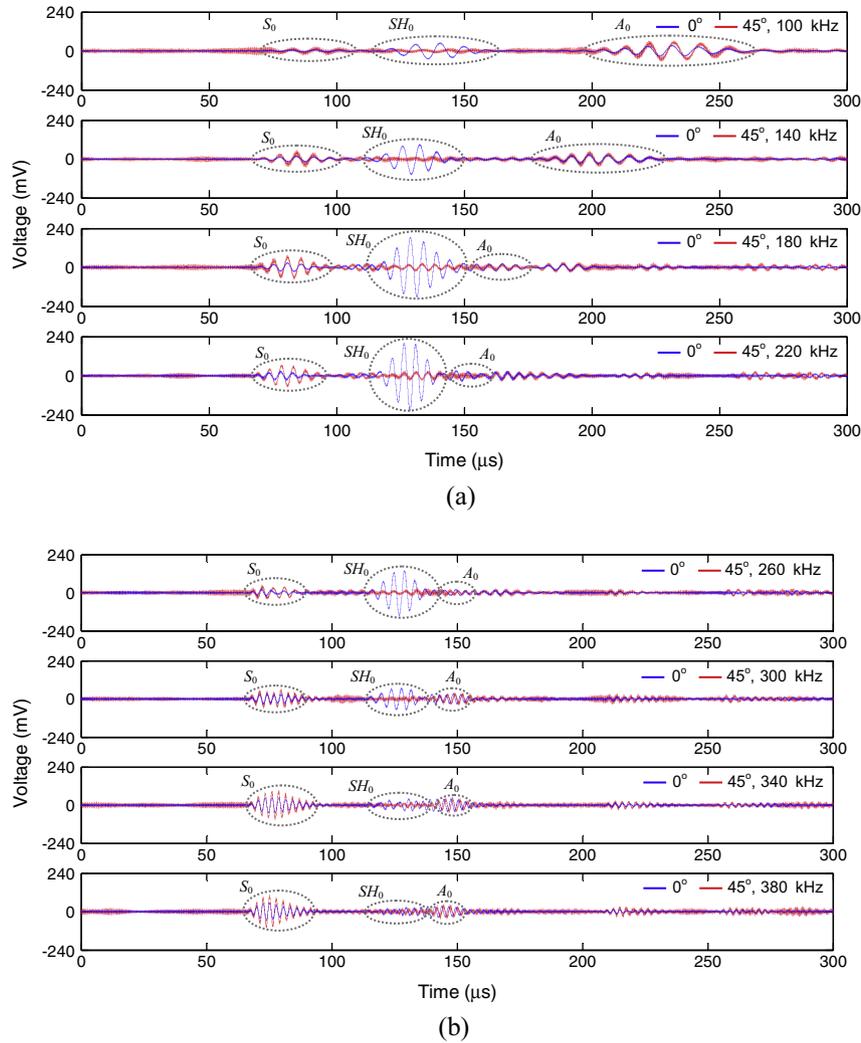


Fig. 9. Layout and location of piezoelectric wafers.



**Fig. 10.** Typical output voltage for different center frequencies at  $0^\circ$  and  $45^\circ$  in experiments: (a) 100–220 kHz and (b) 260–380 kHz.

opposite for  $0^\circ$  and  $90^\circ$ . In Fig. 5(b), the displacement  $v$  is smaller, and on the other two components, displacement responses are also obviously detected. Wave velocity indicates that they are  $A_0$  and  $S_0$  mode Lamb waves. For the sensing points on  $30^\circ$  and  $60^\circ$ , the amplitude of the SH wave becomes less and the  $A_0$  and  $S_0$  mode Lamb waves increase further. In Fig. 5(d), SH wave is unable to be identified, while the other two modes Lamb waves reach their crest value in the amplitude.

Fig. 6 shows absolute values of the maximum amplitude of three displacement components versus the angle of wave propagation from  $0^\circ$  to  $360^\circ$ . According to the guided wave structure,  $v$  is the only component of  $SH_0$  wave, and  $w$  and  $u$  are the dominant components of  $A_0$  and  $S_0$  wave, respectively. For  $SH_0$ ,  $A_0$  and  $S_0$  wave modes, the amplitude curve is extracted at the time  $26.8 \mu\text{s}$ ,  $31.2 \mu\text{s}$  and  $15.2 \mu\text{s}$  respectively (they do not occur simultaneously). Due to the symmetry, the curves in the other three quadrants have the same shapes. It should be noted that the relative values among these three mode waves in this figure are only available at 160 kHz. With tuning frequencies, their relative values will change.

#### 4.2. Guided waves detection using $d_{36}$ type wafer

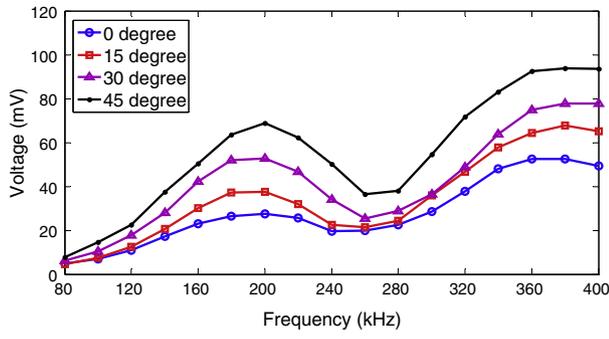
According to reciprocity principle, the voltage responses of  $d_{36}$  type sensor excited by the wave motion will also vary from different angles of wave propagation, since the  $d_{36}$  type

transducer shows apparent directionality when it is used as actuator. Moreover,  $d_{36}$  type actuator is able to generate both Lamb and SH waves, thus the  $d_{36}$  type sensor is expected to detect above two types of guided waves. In this work, two simulations are set up to explore the voltage response of the  $d_{36}$  type sensor received from the waves coming from different angles.

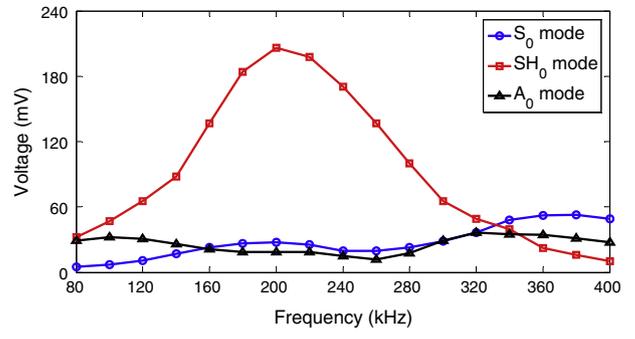
For Lamb waves, a square  $d_{31}$  type piezoelectric wafer is used as actuator, four  $d_{36}$  type sensors are placed around it away the same distance, and their relative angle are  $0^\circ$ ,  $15^\circ$ ,  $30^\circ$  and  $45^\circ$ , respectively. Since the  $d_{31}$  type piezoelectric actuator is omnidirectional, above four sensors will incur the same wave motion. Fig. 7 shows the voltage responses for all sensors.  $d_{36}$  type sensor has maximum sensitivity on  $45^\circ$ , but minimum sensitivity on  $0^\circ$  for both  $A_0$  and  $S_0$  mode Lamb waves.

For investigating the response of  $d_{36}$  sensor to  $SH_0$  wave, a square  $d_{36}$  type actuator is used to excite this non-dispersive wave. Four  $d_{36}$  type sensors are placed on  $0^\circ$ ,  $90^\circ$ ,  $180^\circ$  and  $270^\circ$  of the actuator to detect SH wave, but their relative angle to the actuator are  $0^\circ$ ,  $14^\circ$ ,  $26^\circ$  and  $45^\circ$ , respectively. These angles are selected due to the orthotropic grid in the finite element model.

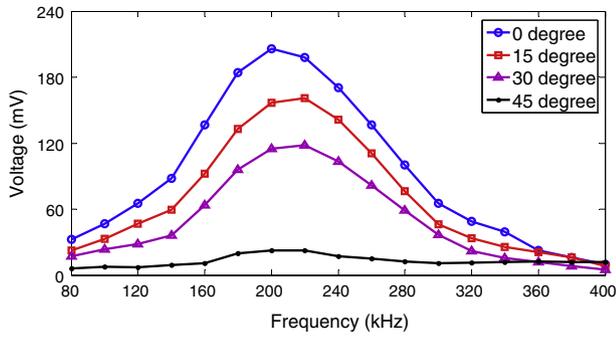
Fig. 8 shows the results from this simulation. The maximum voltage output occurs at the  $0^\circ$ , and the minimum voltage output at the  $45^\circ$ . Since the square wafer covers a small area but not a point, the response at the  $45^\circ$  is not absolute zero.



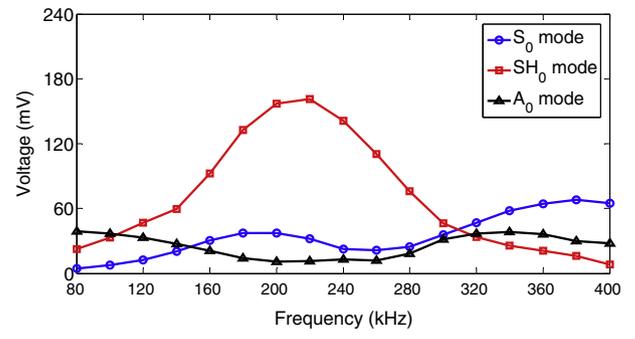
(a)



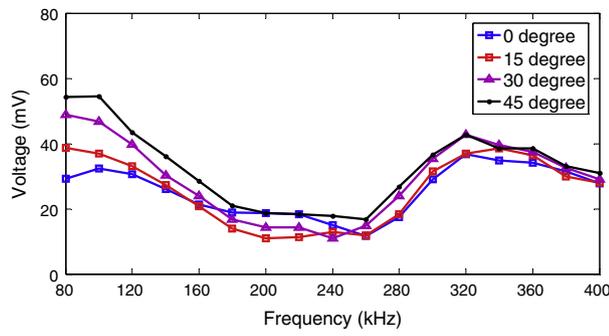
(a)



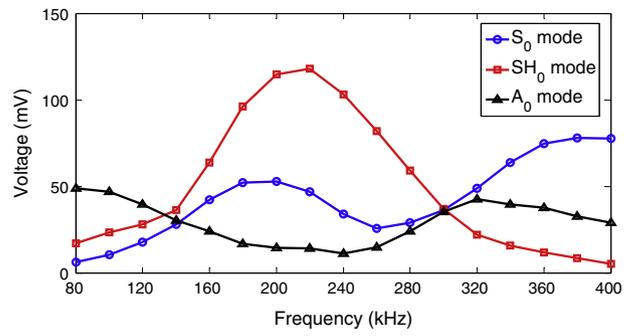
(b)



(b)



(c)



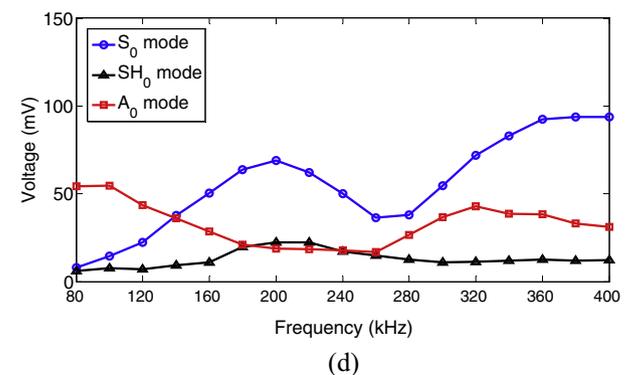
(c)

Fig. 11. Wave amplitude versus frequencies: (a)  $S_0$  mode, (b)  $SH_0$  mode and (c)  $A_0$  mode.

## 5. Experimental investigations

### 5.1. Experimental setup

Experiments were designed to investigate not only the directionality of  $d_{36}$  type actuator and sensor, but also the amplitude sensitivity to frequency. Different from the simulations, in the experiments effects of  $d_{31}$  and  $d_{32}$  piezoelectric coefficients also can be considered. According to the results from FEM, three modes guided waves will be excited and propagate with their respective group velocities, thus the waves with higher velocity will reflect and overlap the waves with slower waves, which results in difficulty in amplitude analysis. Therefore, a large aluminum plate (1200 mm  $\times$  2400 mm  $\times$  1.0 mm) is used in the tests to extend the paths of wave reflections. Five  $d_{36}$  type wafers are mounted on the top surface of an aluminum plate. One of them is used as actuator, and the other four used as sensors are arranged around the actuator at 0°, 15°, 30° and 45° respectively. Their layout and location are shown in Fig. 9, in which the plate and wafers are



(d)

Fig. 12. Wave amplitude versus frequencies: (a) 0°, (b) 15°, (c) 30° and (d) 45°.

not actual size. The distance between all wafers is 35 cm. The size of the wafers as both actuator and sensor is 7 mm  $\times$  7 mm  $\times$  0.5 mm. This thin wafer can reduce the effect of induced bending moment applied to the plate from the wafer. The super glue, which may form very thin and stiff adhesive layer, was used to attach these wafers on the plate. The actuator was driven by a function

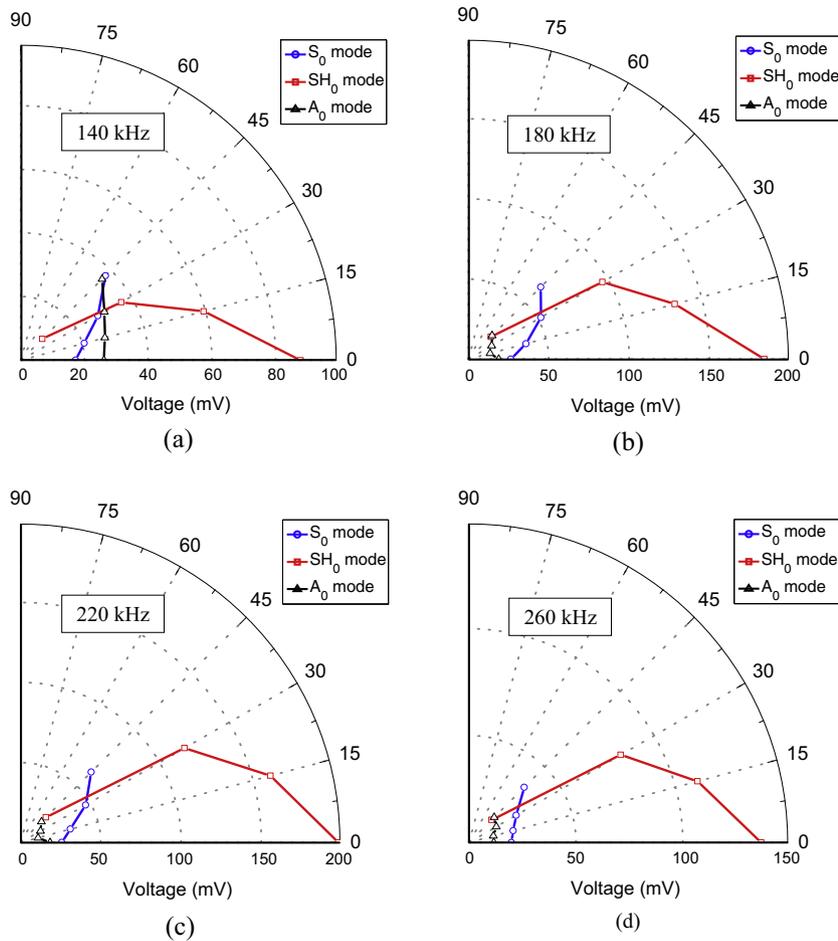


Fig. 13. Wave amplitude versus angle of wave propagation: (a) 140 kHz, (b) 180 kHz, (c) 220 kHz and (d) 260 kHz.

generator (AFG3252C, Tektronix, Inc.), into which a five-peaked wave signal was preprogrammed. Its amplitude and center frequency can be changed during the test. A wideband (DC–2 MHz) power amplifier (TEGAM 2350) was used to amplify excitation voltage. The signal induced from the sensor was collected directly with 2GS/s sampling rate by a digital oscilloscope (DPO 2024, Tektronix, Inc.). The center frequencies of excitation signal  $f_c$  is varied from 80 kHz to 400 kHz in increments of 20 kHz for all angles. According to the group velocities from Fig. 3, among these frequencies, overlapping hardly occurs for all three modes waves, so the reliable amplitudes are able to be obtained from the experimental results.

### 5.2. Amplitude versus frequency for all modes

The voltage responses of piezoelectric sensor vary as the exciting frequency sweeps. Since the experiments obtained a large number of data, only typical output voltage of sensors for different center frequencies at  $0^\circ$  and  $45^\circ$  is selected to be shown in Fig. 10, in which the amplitudes of excitation signals were kept fixed on  $V_{pp} = 2$  V. The  $A_0$ ,  $S_0$  and  $SH_0$  wave modes can be identified from experimental data according to group velocity. These curves show the varying amplitudes and effects of dispersion to  $A_0$  and  $S_0$  modes from lower to higher frequency. The waveforms obviously disperse due to wave dispersion characteristics.

All wave amplitudes with respect to frequencies for all wave modes are shown in Fig. 11. For comparison purpose, all data is

also replotted according to angle of wave propagation in Fig. 12. It can be found that for different angles of wave propagation, the  $S_0$ ,  $SH_0$  and  $A_0$  wave modes have the same tuning frequency characteristics, respectively, but for the same angle, the tuning frequency characteristics are different for these three modes. For  $S_0$  mode, it reaches the first peak at 200 kHz, then decrease to the minimum around 260 kHz and then increase. For  $SH_0$  mode, it reaches maximum value at around 200 kHz, then decrease monotonically. Additional tests indicate that the wave amplitude of the  $SH_0$  mode does not increase again until 1020 kHz, therefore no more data is shown in this work. For  $A_0$  mode, the wave amplitudes also vary with frequency. For the data from  $0^\circ$  direction, the  $SH_0$  mode is dominant from about 160 kHz to 280 kHz. It can be expected that the  $SH_0$  wave mode is easy to be distinguished in this frequency range. In addition, it seems that the data in Fig. 11(c) was affected by errors, which might be induced during amplitude identification for the dispersed wave forms.

### 5.3. Amplitude versus angle of wave propagation for all wave modes

Four excitation frequencies where the  $SH_0$  mode is dominant are selected to show the amplitude variation with angles of wave propagation, shown in Fig. 13. For  $A_0$  and  $S_0$  modes, they still appear due to the existence of  $d_{31}$  and  $d_{32}$  piezoelectric coefficients. In addition, according to Fig. 13 the rates of amplitude change from  $0^\circ$  to  $45^\circ$  for  $SH_0$  mode are various for different frequencies, but the

angles corresponding to maximum and minimum wave amplitudes are the same as those from simulation.

## 6. Conclusions

This paper examines the fundamental characteristics of a new  $d_{36}$  type piezoelectric transducer, made from a PMN–PT crystal. Specifically the tuning frequency characteristics and directionality for the piezoelectric wafer are investigated in application of guided waves generation and sensing. According to the guided wave theory, the  $SH_0$  wave mode is a non-dispersive, which has many advantages for structural damage detection. The PMN–PT wafer has been demonstrated that it is capable of generating and detecting  $SH_0$  wave mode. Due to its materials anisotropy, the  $SH_0$  wave mode exhibits wave directionally. In the preliminarily FEM analyses, it proves above points based on the displacement components generated by the  $d_{36}$  type piezoelectric wafer, and then voltage responses show further that  $SH_0$  wave mode is dominant at  $0^\circ$ , but negligible small at  $45^\circ$ ; while the  $A_0$  and  $S_0$  wave mode show the reverse trend.

Experiments were conducted to investigate both the tuning frequency characteristics and wave mode directionality for the proposed wafer. Based on wave amplitudes obtained, the following conclusions are drawn:

- (1) The  $A_0$  and  $S_0$  wave modes are affected greatly by dispersion; while the  $SH_0$  wave mode maintains original waveform and constant group velocity with respect to varying frequencies.
- (2) In the direction along  $0^\circ$ ,  $SH_0$  wave mode is dominant between about 160 kHz and 280 kHz, and then decreases gradually. Similar phenomena are observed for the other three directions. In the direction along  $45^\circ$ ,  $SH_0$  wave mode is hardly detected, whereas the  $A_0$  and  $S_0$  wave modes are mostly dominant. It indicates that among the specific frequency range, the proposed wafer can be used to generate and detect  $SH_0$  wave in specific directions. With this property, the  $d_{36}$  type piezoelectric wafer is suited to be used in long structure for damage detection, because the other modes wave will become decrease in magnitude through long travel.

- (3) Wave amplitudes of  $A_0$  and  $S_0$  wave modes in different directions have similar frequency tuning characteristic, but have different turning points from the PZT reported in the literature [10]. This is due to inherently different mechanical properties of piezoelectric wafer.

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